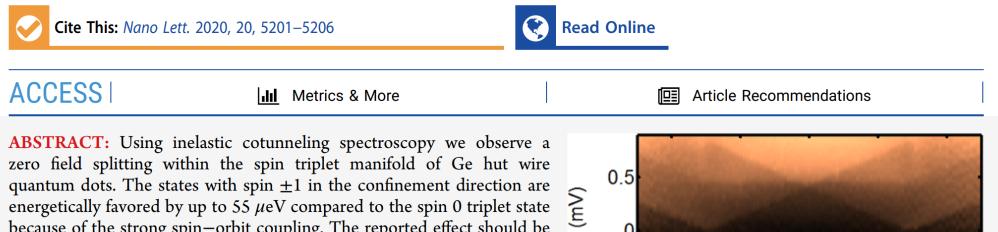
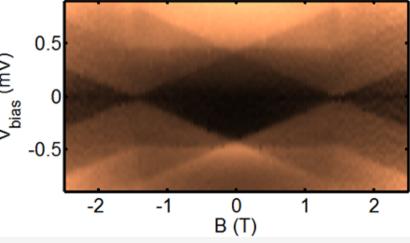
Zero Field Splitting of Heavy-Hole States in Quantum Dots

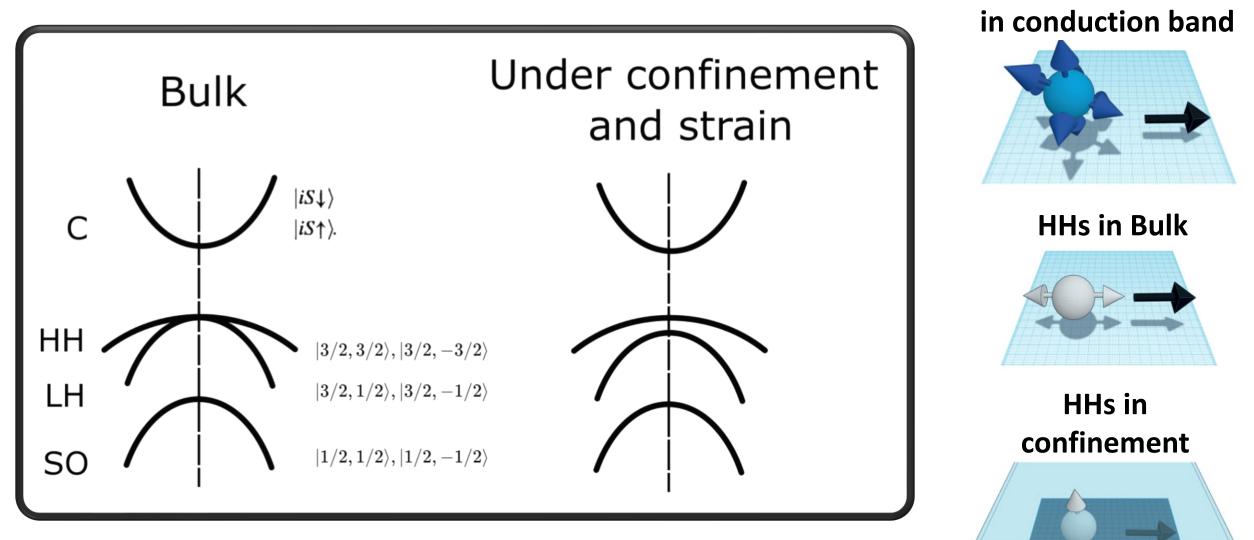
Georgios Katsaros,* Josip Kukučka, Lada Vukušić, Hannes Watzinger, Fei Gao, Ting Wang, Jian-Jun Zhang, and Karsten Held



quantum dots. The states with spin ± 1 in the confinement direction are energetically favored by up to 55 μ eV compared to the spin 0 triplet state because of the strong spin-orbit coupling. The reported effect should be observable in a broad class of strongly confined hole quantum-dot systems and might need to be considered when operating hole spin qubits.

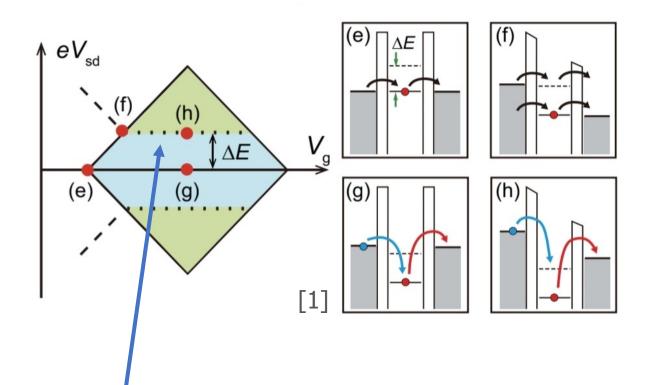


Heavy hole states



Electrons

Elastic/Inelastic Cotunneling



Onset of inelastic cotunneling

No VG dependence

[1] Phys. Rev. B **87**, 041302

First order tunneling Processes

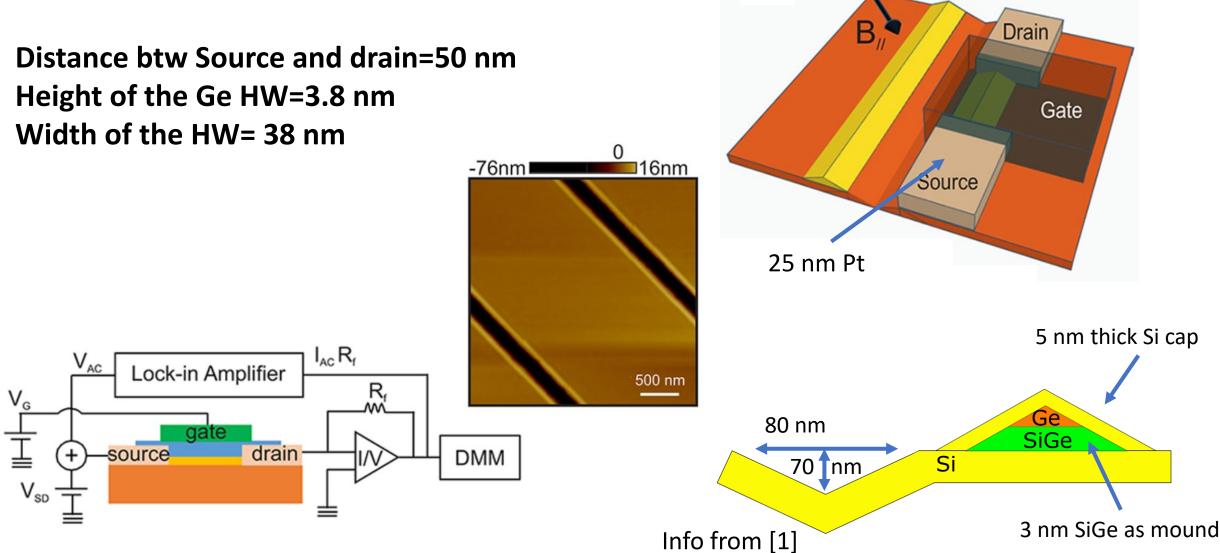
(e), and (f) Sequential tunneling transport

Blocked inside the diamond

Higher order tunneling processes

- (g) Elastic Cotunneling
- (h) Inelastic Cotunneling $eV_{sd} \ge \Delta E$

Setup and device

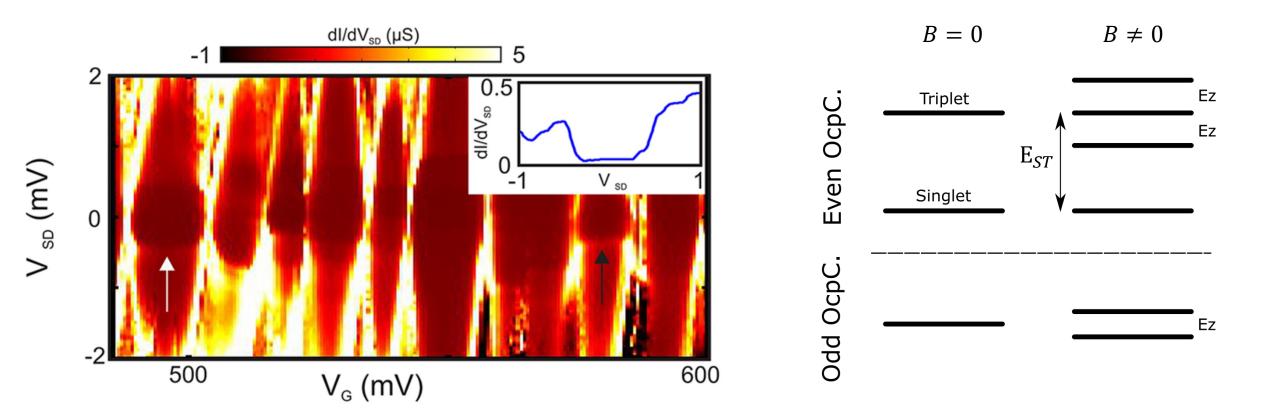


B⊥

[1] F. Gao, et al. Advanced Materials 32.16 (2020): 1906523

Stability diagram

Odd/even occupancy



By comparing E_{ST} and E_{ORB}, one can obtain information about the strength of hole-hole interactions

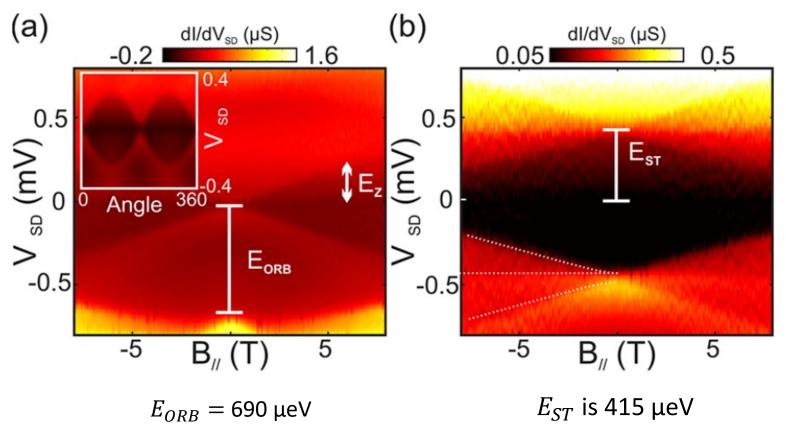
Effective electron temperature of 100 mK

Odd no. of holes \rightarrow splitting of the doublet Even no. of holes \rightarrow splitting of the triplet

(a) doublet, VG= 510.5 mV

Inset: B = 1 T,

- \rightarrow g-factor anisotropy of about 7.5
- $(g \parallel = 0.56 \pm 0.06 \text{ and } g \perp = 4.17 \pm 0.22)$
- This anisotropy is due to the HH character of the confined states [1]



b) singlet VG = 528.3 mV \rightarrow g|| = 0.57 ± 0.01 and g \perp =4.56 ± 0.16

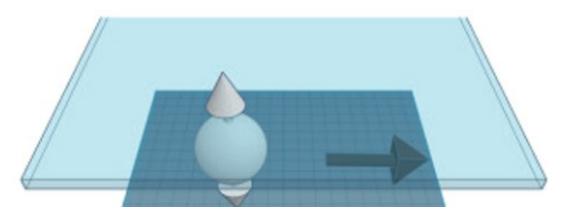
- triplets are nonequally spaced
- ZFS of 55 μeV

 $E_{ORB} - E_{ST} = 275 \,\mu\text{eV}$

Coulomb interaction energy

Model

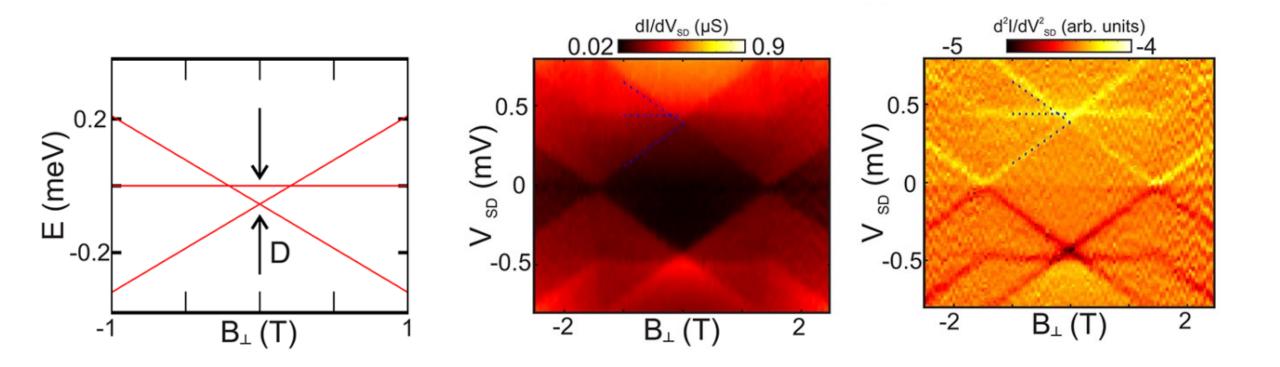
Ge is known to have a very strong valence band SOC which leads to the HH spin pointing in the perpendicular direction. Assuming the triplet state as being made up of two such HH spins.



$$H = -J/2SS + g_{\perp}\mu_{B}S_{\perp}B_{\perp} + g_{\parallel}\mu_{B}S_{\parallel}B_{\parallel} - DS_{\perp}^{2}$$

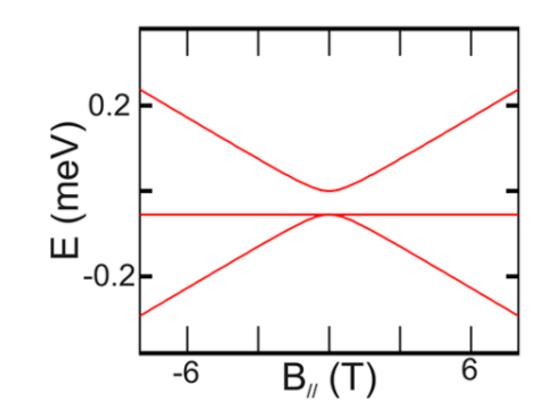
Exchange term J differentiates
singlet and triplet
Zeeman terms
Zeeman terms
The singlet and triplet terms
The singlet and triplet terms
The singlet and triplet terms terms
The singlet and triplet terms terms

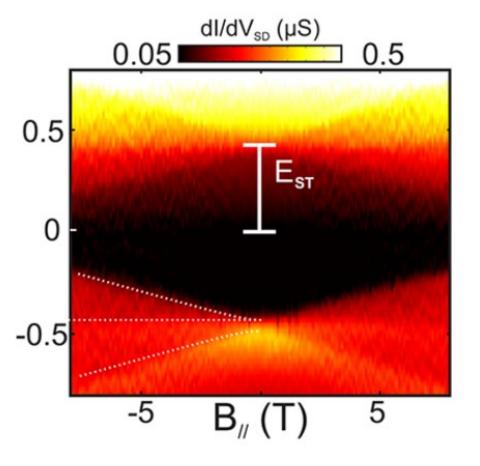
Simulation and the measurement for B_{\perp}



 g_{\parallel} = 0.57, g_{\perp} = 4.56, and D = 55 µeV, as extracted from the

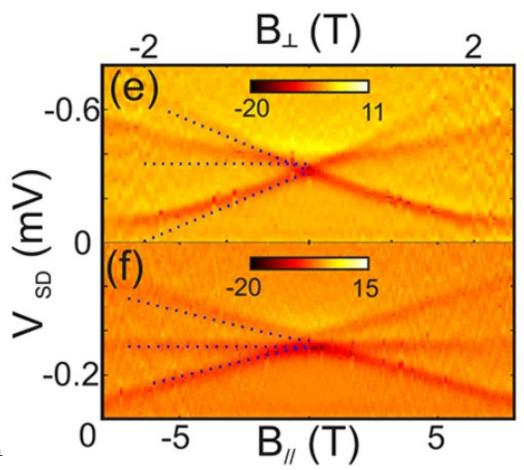
Simulation and the measurement for $B_{||}$





For large $g_{\parallel}\mu B_{\parallel} \gg D$, the usual Zeeman splitting of the triplet states into $S_{\parallel} = \pm 1$, 0 is recovered as the HH pseudo spins now reorient along B_{\parallel} .

Measurements on a 2nd device



In this case **orbital effects** \rightarrow slight bending of the states for B₁ ZFS=35 µeV g|| = 0.52 ± 0.13, g1 = 2.78 ± 0.06

Summary

- They investigated spin anisotropy of HH states
- They measured the ZFS for HH states in a 2D quantum dot
- They extracted the hole-hole interaction

Thanks for your attention

