

**Transport Signatures of Quasiparticle Poisoning in a Majorana Island**

S. M. Albrecht,¹ E. B. Hansen,¹ A. P. Higginbotham,^{1,2} F. Kuemmeth,¹ T. S. Jespersen,¹ J. Nygård,¹ P. Krogstrup,¹
J. Danon,^{1,3} K. Flensberg,¹ and C. M. Marcus¹

¹*Center for Quantum Devices and Station Q Copenhagen, Niels Bohr Institute, University of Copenhagen, Copenhagen 2100, Denmark*

²*JILA, University of Colorado and NIST, Boulder, Colorado 80309, USA*

³*Department of Physics, NTNU, Norwegian University of Science and Technology, 7491 Trondheim, Norway*

(Received 17 December 2016; published 27 March 2017)

We investigate effects of quasiparticle poisoning in a Majorana island with strong tunnel coupling to normal-metal leads. In addition to the main Coulomb blockade diamonds, “shadow” diamonds appear, shifted by $1e$ in gate voltage, consistent with transport through an excited (poisoned) state of the island. Comparison to a simple model yields an estimate of parity lifetime for the strongly coupled island ($\sim 1 \mu\text{s}$) and sets a bound for a weakly coupled island ($> 10 \mu\text{s}$). Fluctuations in the gate-voltage spacing of Coulomb peaks at high field, reflecting Majorana hybridization, are enhanced by the reduced lever arm at strong coupling. When converted from gate voltage to energy units, fluctuations are consistent with previous measurements.

LETTER

doi:10.1038/nature17162

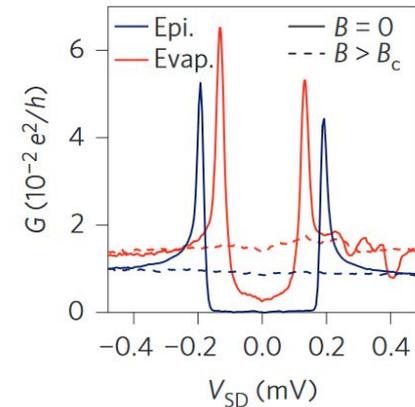
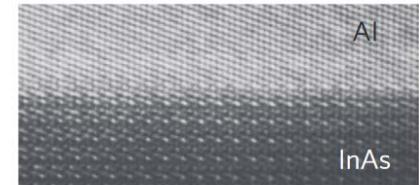
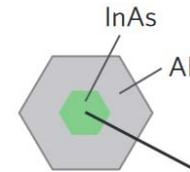
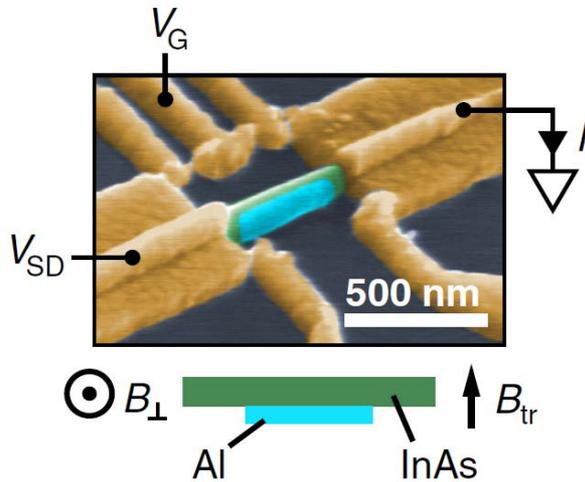
206 | NATURE | VOL 531 | 10 MARCH 2016

Exponential protection of zero modes in Majorana islands

S. M. Albrecht^{1*}, A. P. Higginbotham^{1,2*}, M. Madsen¹, F. Kuemmeth¹, T. S. Jespersen¹, J. Nygård¹, P. Krogstrup¹ & C. M. Marcus¹

Device Layout

- Sample:
- [0001] wurtzite InAs nanowire
 - MBE Al on 2 of 6 faces (hard gap)
 - Chemical etch to remove Al shell at ends
→ L=400nm Al island
 - Ti/Au contacts
 - Gates near 50nm exposed segment to form dot



- B-field:
- perpendicular to wire
 - Thin film Al (10nm) → high critical field:
 $B_{\text{crit}} = 0.7 \text{ T}$ (perp. substrate)
 $B_{\text{crit}} = 0.2 \text{ T}$ (in plane)

Measurements: Lockin 5 μ V, 314Hz, 50mK (fridge)

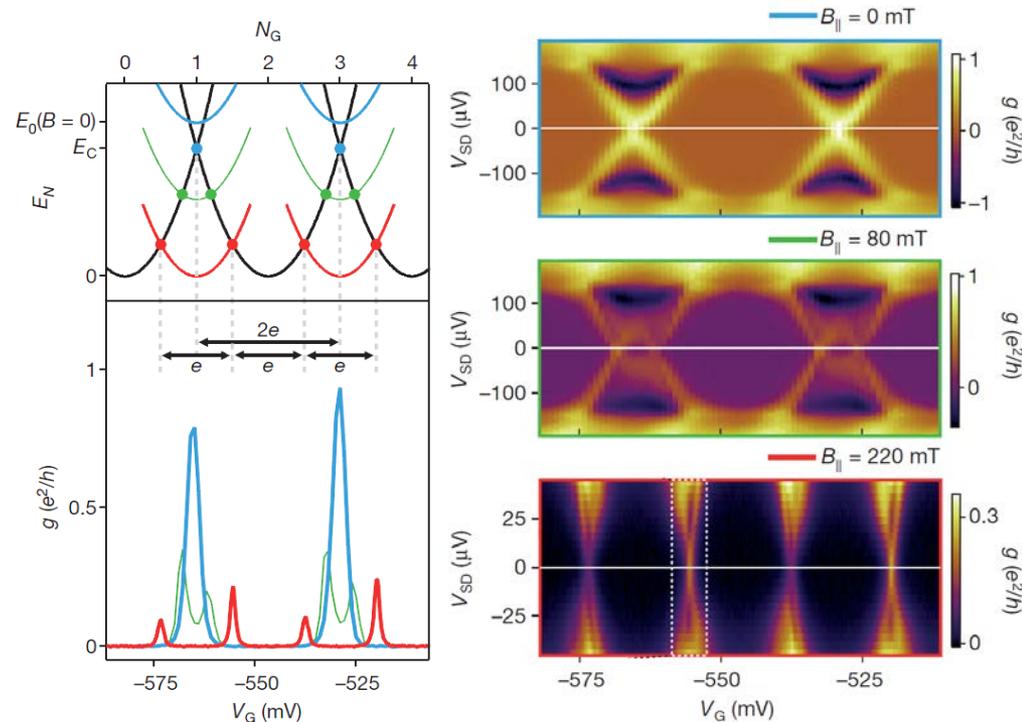
W. Chang et al., Nat. Nanotech. **10**, 232 (2015)

Charge states of island (no poisoning)

Number of charges on island: $N = 2N_{cp} + N_0$
 N_{cp} : Cooper pairs on island
 N_0 : (0/1) occupation of subgap state

State energies $E(N_{cp}, N_{\Delta}, N_0) = E_C/2 * (N_G - N)^2 + N_0 * E_0$
 Gate induced charge $N_G = C_G V_G / e$
 Capacitance C_G (island – side gate)
 Charging energy $E_C = e^2 / C_{tot}$

Even N: $(N_{cp}, 0)$ lowest state: condensate
 Odd N: $(N_{cp}, 1)$ 1 particle in subgap state



S.M. Albrecht et al., Nature **531**, 206 (2016)

Charge transport @ degeneracies

- $E_0 > E_C$ Transport through cooper pairs => $2e$ spacing
- $E_0 < E_C$ Alternating peak spacing S_o, S_e (E_0 lowered by Zeeman energy)
- $E_0 = 0$ Transport through Majorana zero modes => $1e$ spacing

Charge states of island

Number of charges on island: $N = 2N_{cp} + N_{\Delta} + N_0$
 N_{cp} : Cooper pairs on island
 N_{Δ} : Quasiparticles in BCS continuum
 N_0 : (0/1) occupation of subgap state

State energies $E(N_{cp}, N_{\Delta}, N_0) = E_C/2 * (N_G - N)^2 + N_{\Delta} * \Delta + N_0 * E_0$
 Gate induced charge $N_G = C_G V_G / e$
 Capacitance C_G (island – side gate)
 Charging energy $E_C = e^2 / C_{tot}$

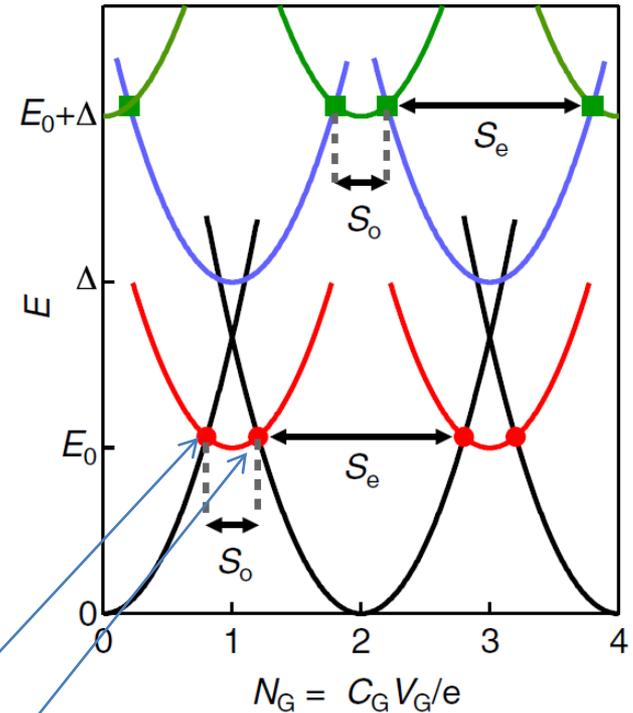
Even N: $(N_{cp}, 0, 0)$ lowest state is pure condensate, **black**
 $(N_{cp}, 1, 1)$ 1st excited state, **green**
 Odd N: $(N_{cp}, 0, 1)$ 1 particle in subgap state, **red**
 $(N_{cp}, 1, 0)$ 1st excited state, **blue**

Charge transport @T=0 (no poisoning):

- Island in ground state
- Transport @ degeneracies

$$(N_{cp}, 0, 0) \rightleftharpoons (N_{cp}, 0, 1)$$

$$(N_{cp}, 0, 1) \rightleftharpoons (N_{cp} + 1, 0, 0)$$



Transport with poisoning

Poisoning: Presence of quasiparticles (relax in energy down to BCS peak @ Δ)

Transport at degeneracies with $N_{\Delta}=1$ (green squares)

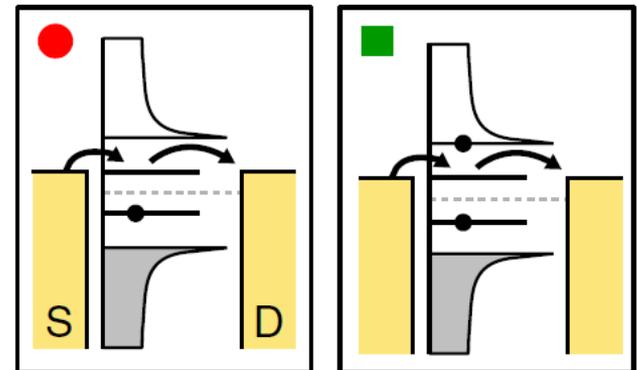
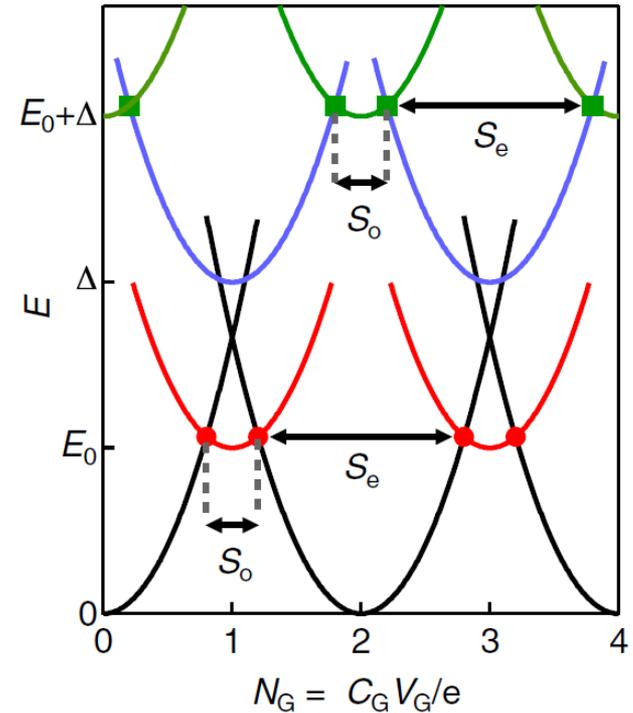
Returning to unpoisoned state:

- Cooper pair recombination ($(N_{cp}, 1, 1) \rightarrow (N_{cp}+1, 0, 0)$)
(possible due to lack of transl. invariance)
- Quasiparticle relaxation to subgap state:
($(N_{cp}, 1, 0) \rightarrow (N_{cp}, 0, 1)$)
- Quasiparticle tunneling to lead
($(N_{cp}, 1, N_0) \rightarrow (N_{cp}, 0, N_0)$)

Visibility of “excited state transport”:

- Relaxation rate
- Poisoning rate
- Γ_{SD}

Appearance: Same as ground state, but shifted by $1e$ in gate voltage



Coulomb blockade diamonds

Without subgap state:

1 set of diamonds with $2e$ peak spacing

Subgap state:

Alternating peak spacing S_o, S_e

2 sets of diamonds

- Small strong
- Faint large

Nearly vanishing odd diamonds

=> E_0 only slightly smaller than $E_C/2$

- Weak lines within large diamond
→ quasiparticle transport

Extracted parameters:

$E_C=210\mu\text{eV}$, $E_0=75\mu\text{eV}$, $\eta=6\text{meV/V}$,

$\Delta=140\mu\text{eV}$ (induced gap, onset negative diff. cond.)

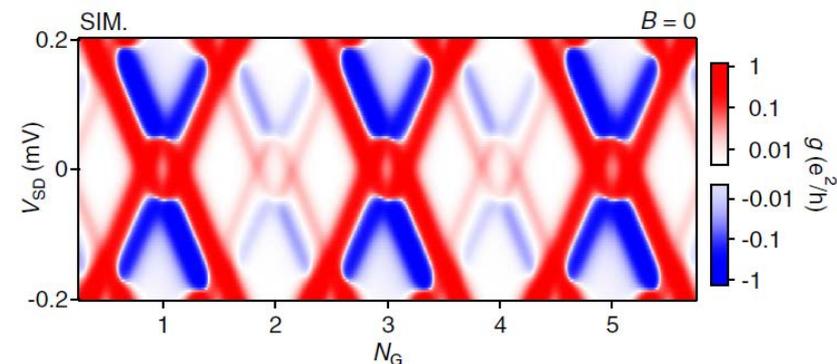
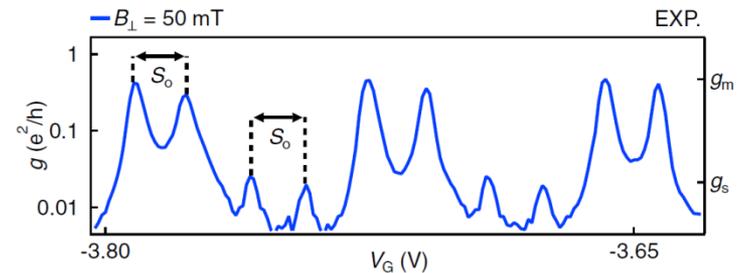
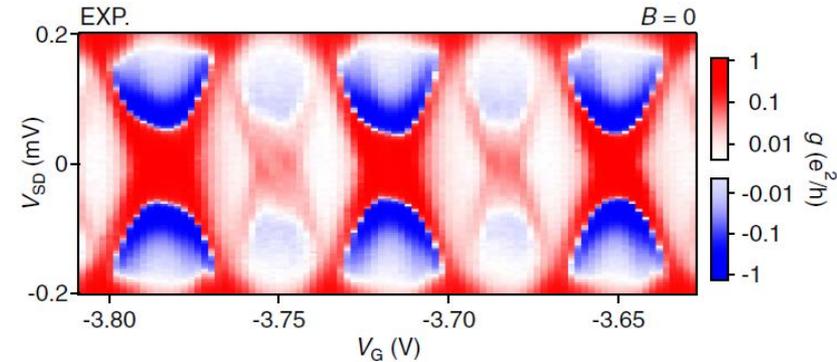
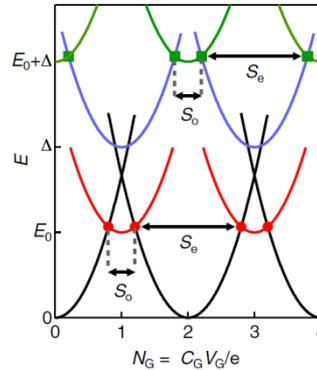
$\Gamma_S=1\text{GHz}$, $\Gamma_D=6\text{GHz}$ (peak height, width), previously $\approx 0.5\text{GHz}$

$\tau_{qp}=0.1\mu\text{s}$ (relaxation quasiparticle to subgap state, previously meas. in similar devices)

Numerical simulations (using above parameters):

Best match using $\tau_p=1.2\mu\text{s}$ (poisoning time)

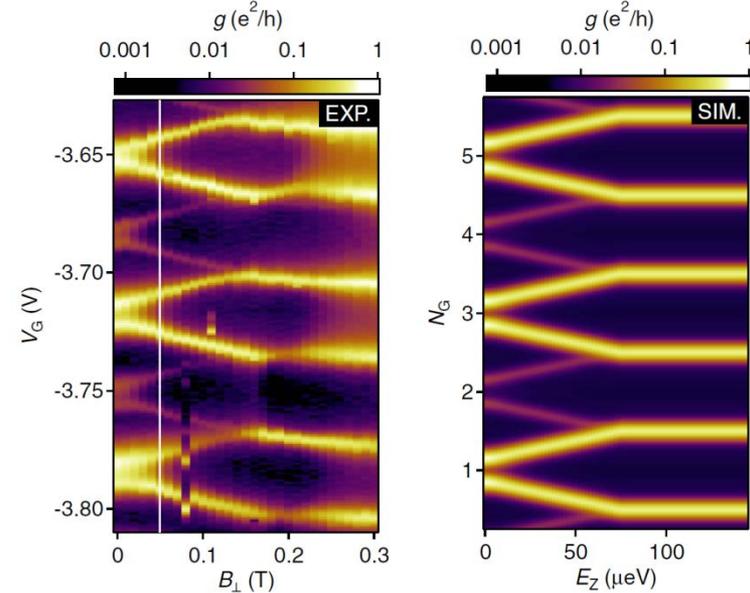
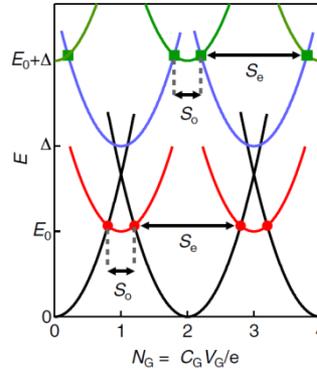
- Electron like $(N_{cp}, 0, N_0) \rightarrow (N_{cp}, 1, N_0)$
- Hole like: $(N_{cp}, 0, N_0) \rightarrow (N_{cp}-1, 1, N_0)$



Evolution to Majorana modes

Emergence of zero energy mode:

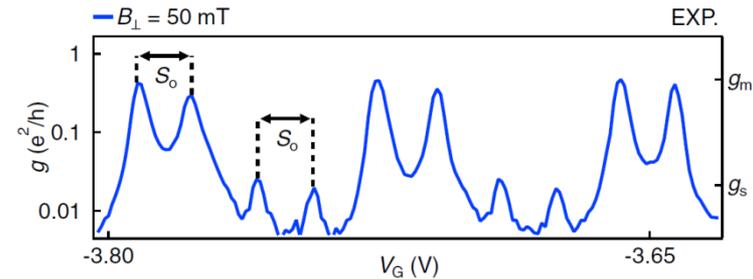
- Increasing B_{perp} lowers state energy of E_0 :
 $E_0 = 75 \mu\text{eV} - E_Z$
- Uniform spacing $S_0 = S_e$ for $B_{\text{perp}} > 0.16\text{T}$
 \Rightarrow zero energy state
- For $B > 0.16\text{T}$: state remains at zero with small oscillations
 \Rightarrow hybridized Majorana modes



Simulations use same parameters as for CBP

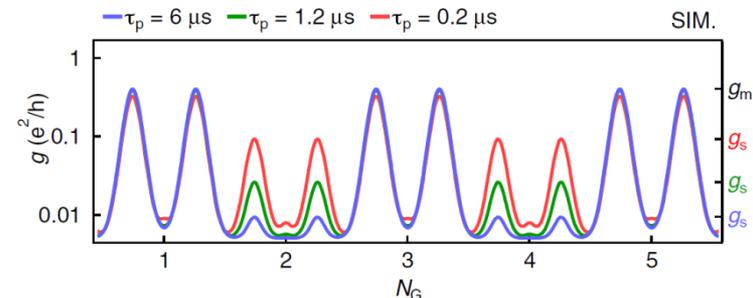
Shadow diamond & poisoning rate:

- Clearly split Coulomb peaks at finite field
- Strongly suppressed shadow peaks
- Peak height depends strongly on τ_p
 \Rightarrow Poisoning rate $\tau_p = 1.2 \mu\text{s}$



Parity lifetime of $\tau_p = 1.2 \mu\text{s}$ (obtained from ration g_m/g_s)

- Here, not ideal tuning for long lifetimes (large source drain coupling)
- Estimate g_m/g_s for devices without shadow diamond ($g_s =$ noise floor) \Rightarrow conservative estimate $\tau_p > 10 \mu\text{s}$



Hybridized Majorana modes

B-field evolution of subgap state:

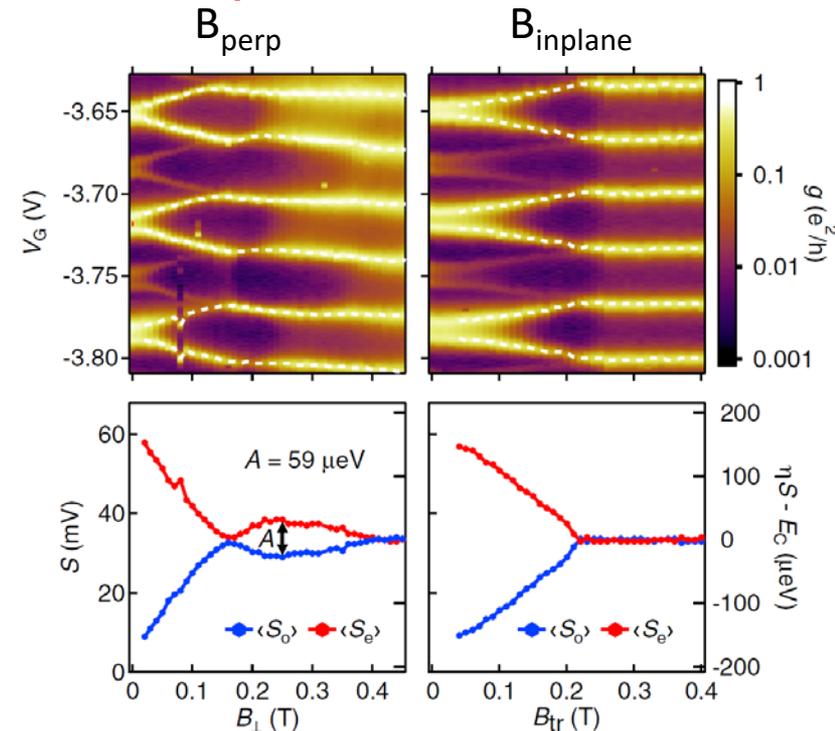
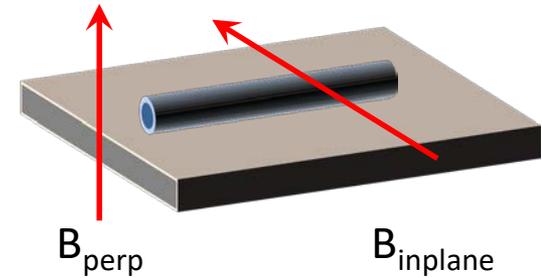
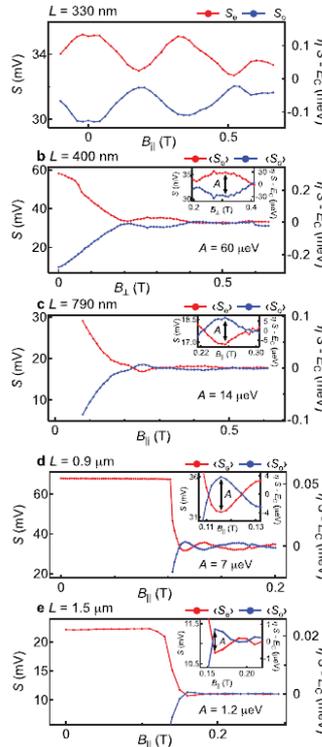
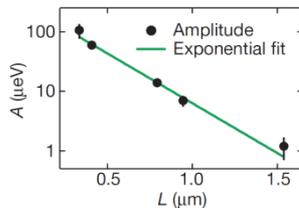
- spacing $S_0(S_e)$ increases (decreases) with growing field, both B_{perp} , B_{inplane}
- Strong and faint peaks merge
 - \Rightarrow Uniform peak spacing
 - \Rightarrow zero energy state
- Small oscillations for B_{perp} consistent with hybridizing Majorana modes in finite length device:

$$A = A_0 * e^{-L/\xi} = 64 \mu\text{eV}$$

$$L = 400\text{nm}$$

$$A_0 = 300 \mu\text{eV}$$

$$\xi = 260\text{nm}$$



No oscillations for B_{inplane} ?

- Merging happens at higher field, close to gap closing
- \Rightarrow expect "trivial" $1e$ period, no Majorana modes

Summary

- Observation of merging (oscillating) subgap states
- consistent with hybridizing Majorana modes in finite length samples
- Observation of shadow Coulomb diamond => quasiparticles
- Intensity of shadow diamond => quasiparticle poisoning rates
- Majorana parity lifetime of $10\mu\text{s}$ ($1.2\mu\text{s}$) in weakly (strongly) coupled devices