

# Coherent Control of a Single Electron Spin with Electric Fields

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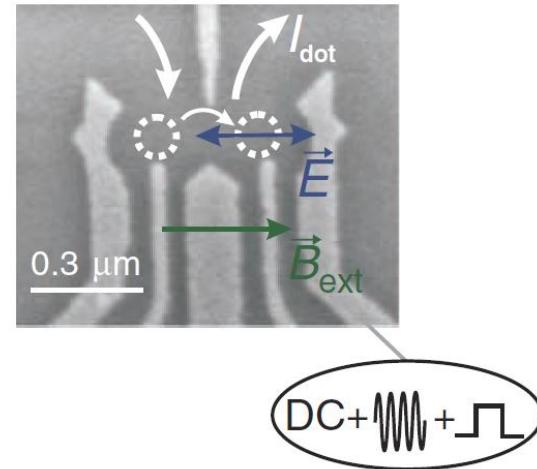
Manipulation of single spins is essential for spin-based quantum information processing. Electrical control instead of magnetic control is particularly appealing for this purpose, because electric fields are easy to generate locally on-chip. We experimentally realized coherent control of a single-electron spin in a quantum dot using an oscillating electric field generated by a local gate. The electric field induced coherent transitions (Rabi oscillations) between spin-up and spin-down with 90° rotations as fast as ~55 nanoseconds. Our analysis indicated that the electrically induced spin transitions were mediated by the spin-orbit interaction. Taken together with the recently demonstrated coherent exchange of two neighboring spins, our results establish the feasibility of fully electrical manipulation of spin qubits.

*Science* **318**, 1430 (Nov 30, 2007)

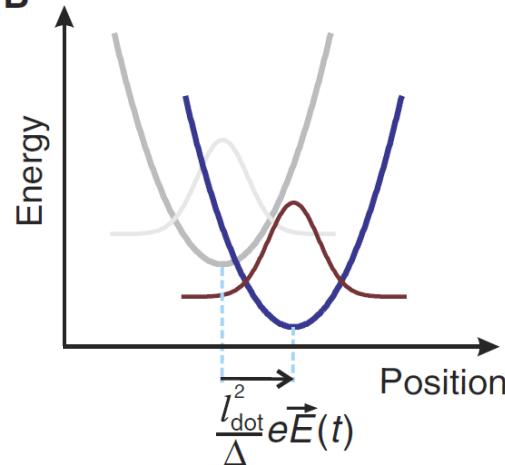
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<sup>†</sup>These authors contributed equally to this work.

A

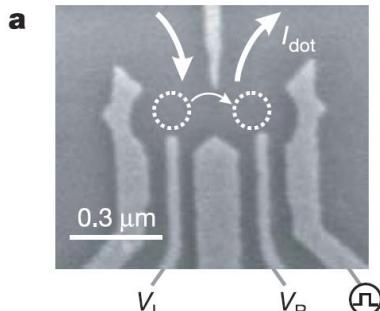


B

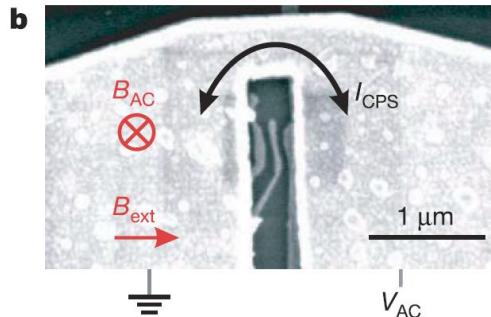


# Electron Spin Resonance (ESR)

Strip line produces  $B_{AC}$



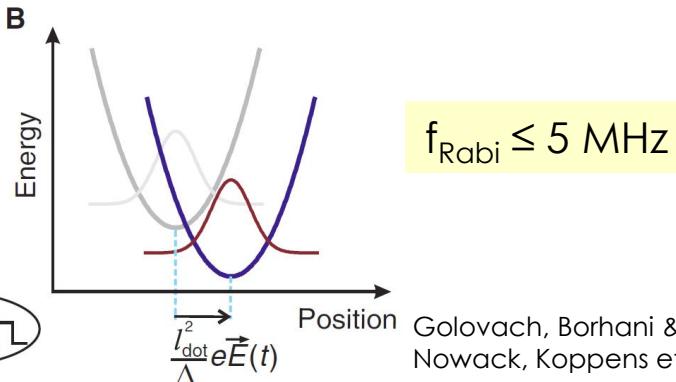
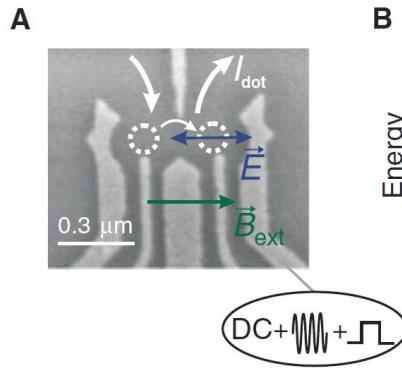
Koppens et al., Nature 2006



$$f_{\text{Rabi}} \leq 10 \text{ MHz}$$

comparison: SWAP > 1 GHz

Electric dipole spin resonance: spin-orbit coupling



Golovach, Borhani & Loss, PRB 2006  
Nowack, Koppens et al., Science 2007

other methods  
gradients

- nuclear
- micro magnet
- exchange
- g - factor

# Exploiting semiconductor spin-orbit coupling

$$|\mathbf{B}_{\text{eff}}(t)| = 2 |\mathbf{B}_{\text{ext}}| \frac{l_{\text{dot}} e |\mathbf{E}(t)|}{l_{\text{SO}}} \frac{l_{\text{dot}}}{\Delta}$$

dot orbitals       $\Delta, l_{\text{dot}}$   
spin-orbit length       $l_{\text{SO}}$   
electric field       $E(t)$

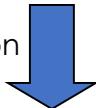
Golovach, Borhani, Loss, PRB 2006

semiconductor spin-orbit lengths  
dot sizes

$l_{\text{SO}} \sim 50 \text{ nm to few microns}$   
 $l_{\text{dot}} \sim 10 \text{ nm to } 100 \text{ nm (single charge dot)}$

$$\lambda_{\pm} = \hbar^2 / (2m^* |\alpha \pm \beta|)$$

$$H_{SO} = \underline{\alpha(p_x\sigma_y - p_y\sigma_x)} + \underline{\beta(p_x\sigma_y + p_y\sigma_x)}$$

45° rotation  Rashba term  
structural asymmetry

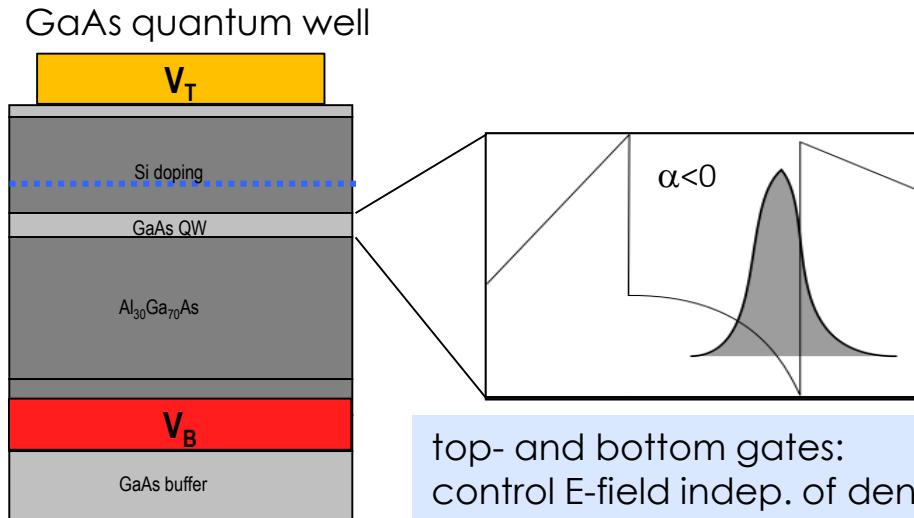
Dresselhaus term (linear)  
crystal asymmetry

$$H_{SO} = (\beta - \alpha) p_y \sigma_x + (\beta + \alpha) p_x \sigma_y$$

Typical wafers:  $\alpha \sim \beta$   
“weak” direction  $\alpha - \beta$  : slow EDSR < 10 MHz  
“strong” direction  $\alpha + \beta$  : fast EDSR > 100 MHz

# Rashba term

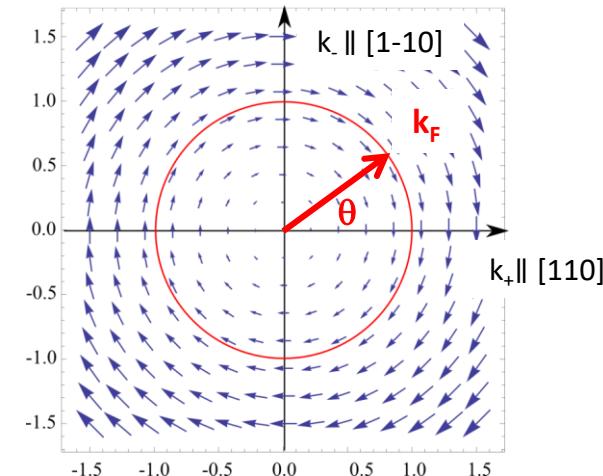
- Structure inversion asymmetry:
  - electric field, QW asymmetry
  - doping (effective E-field)
  - tilting of the QW with gate



2D Rashba Hamiltonian

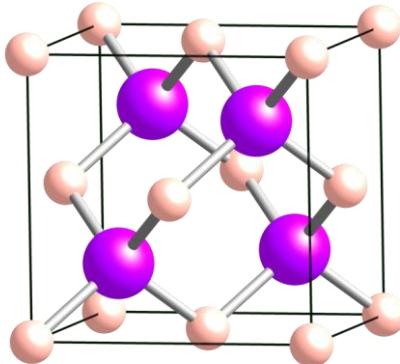
$$\mathcal{H}_R = \alpha(\sigma_x k_y - \sigma_y k_x)$$

$$\alpha \sim E_{\text{eff}}$$



first harmonics

# Dresselhaus term



GaAs,InAs,InSb,.... (III-V compounds)

- zinc-blende structure
- 2 atomic basis, polar semiconductor
- No center of inversion for electric fields → Bulk inversion asymmetry (BIA)

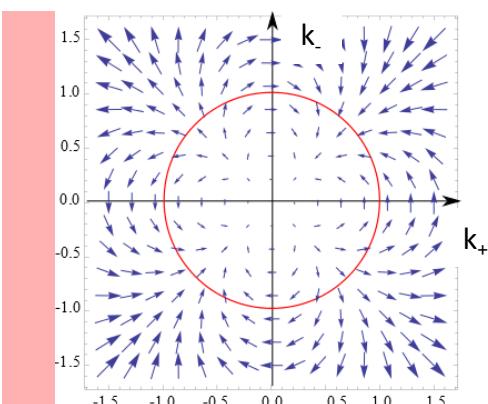
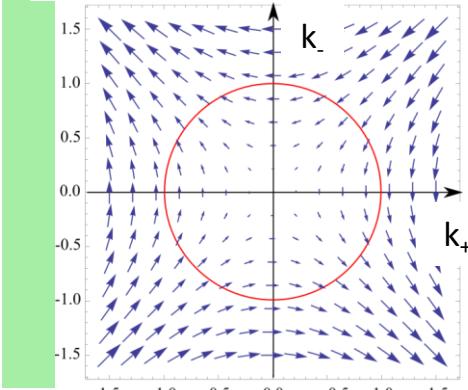
Dresselhaus Hamiltonian in 3D:

$$\mathcal{H}_D = \gamma(\sigma_x k_x k_y^2 - \sigma_y k_y k_x^2) + \text{cycl. perm.}$$

g: Dresselhaus coefficient: material constant

# Dresselhaus term in a quantum well

$$H_D = \beta(p_y\sigma_y - p_x\sigma_x)$$



- Coordinate transformation [110] main axes
- quantum well: quantize z-direction( $\langle k_z \rangle = 0$ )
- Rewrite in polar coordinates

$$\mathcal{H}_D = k_F [(\beta \sin(\theta) + \beta_3 \sin(3\theta))\sigma_+ \\ - (\beta \cos(\theta) + \beta_3 \cos(3\theta))\sigma_-]$$

density n  
gate tunable

$$\beta = \gamma \langle k_z^2 \rangle - \gamma \frac{\pi n}{2}$$

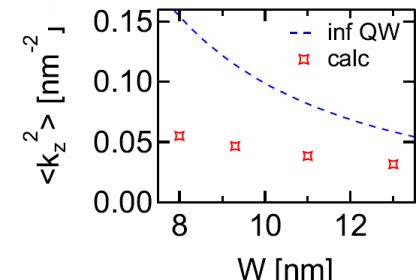
well width

$$\beta_3 = \gamma \frac{\pi n}{2}$$

→  $\beta$  and  $\beta_3$  are density dependent!  
→ **Dresselhaus SOI can be gate-tuned**

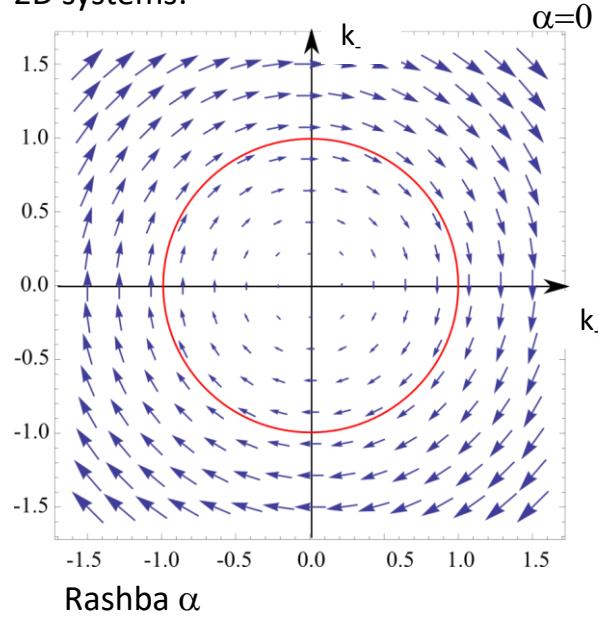
first harmonics

third harmonics

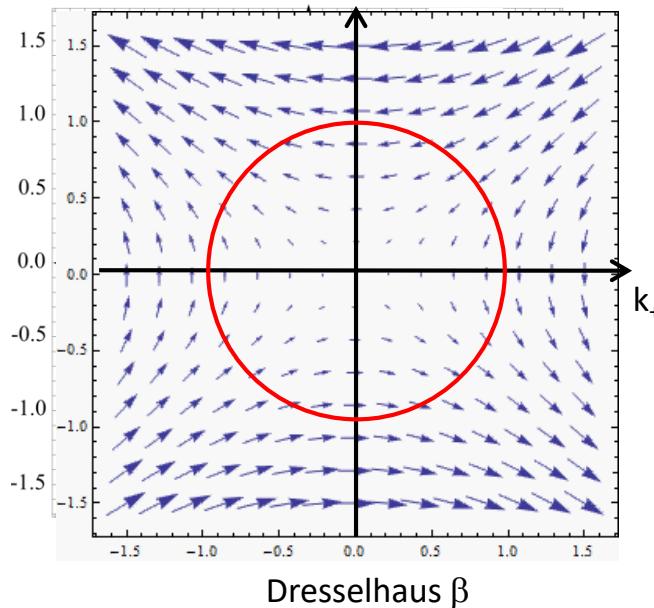


# The Persistent Spin Helix

2D systems:



+



Rashba  $\alpha$

first harmonics

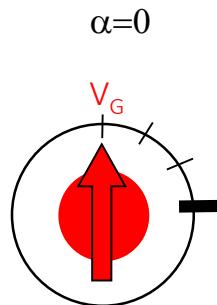
third harmonics

J. Schliemann, J.C. Egues and D. Loss, PRL 2003  
A. Bernevig, J. Orenstein, S. Zhang, PRL 2006

# The Persistent Spin Helix

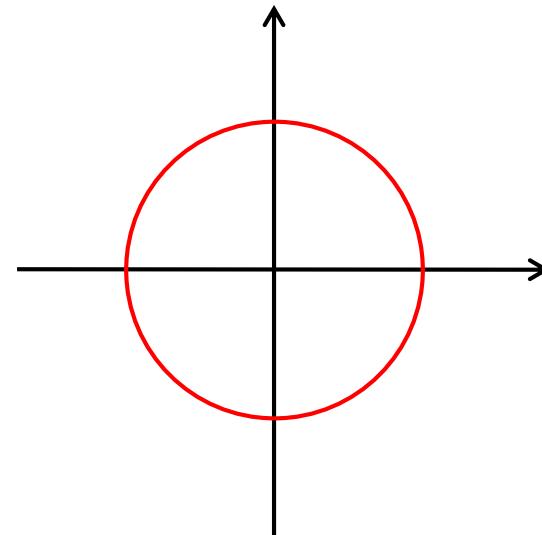
2D systems:

$$\alpha(V_G) = \beta(V_G)$$



- Persistent spin helix (PSH) regime at  $\alpha=\beta$
- UNIAXIAL spin orbit field

Rashba  $\alpha$



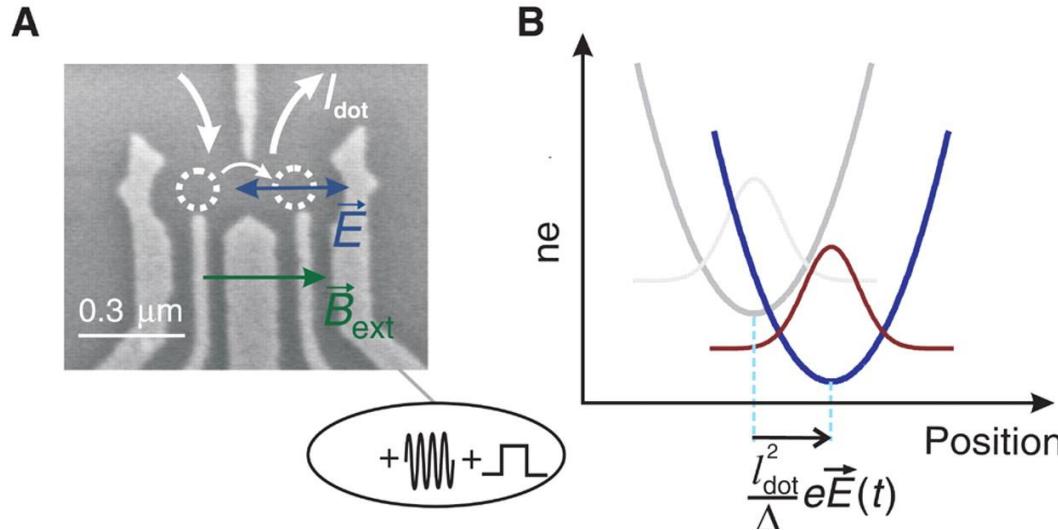
Dresselhaus  $\beta$

first harmonics

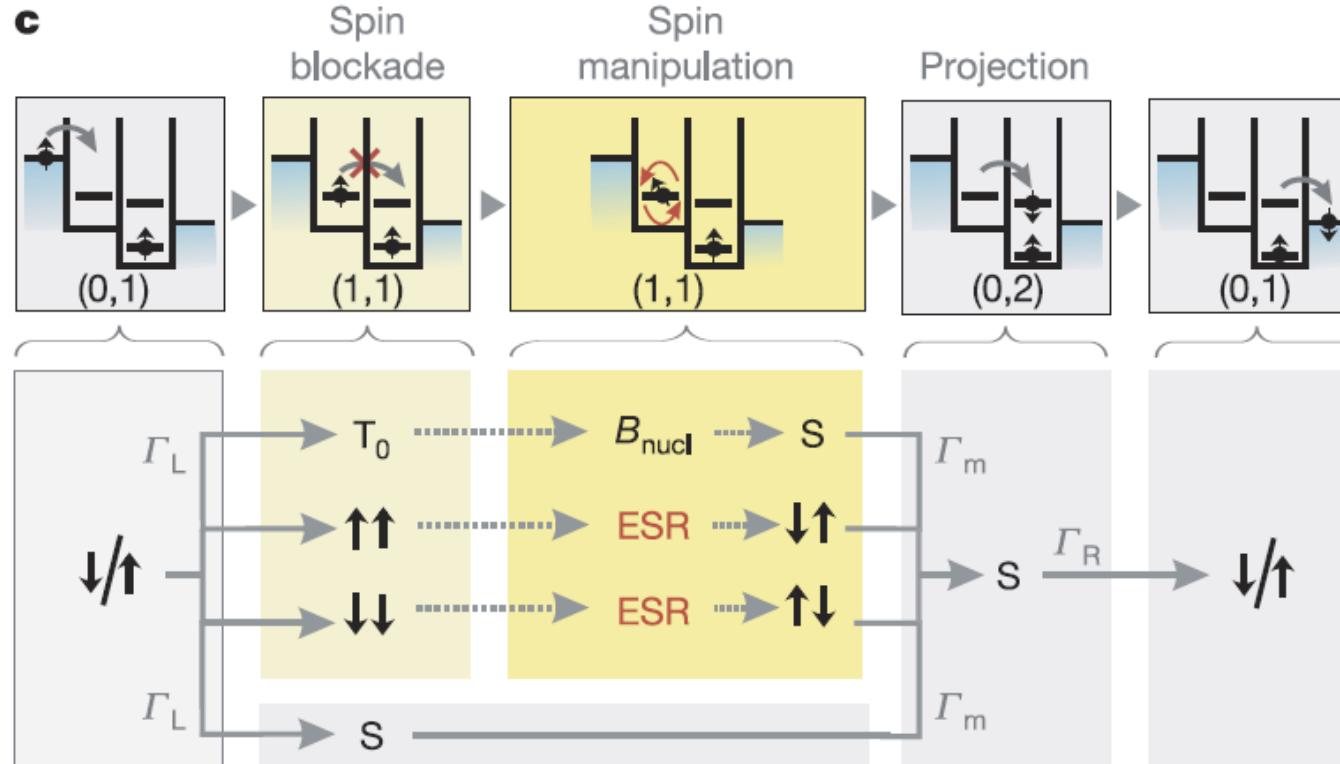
third harmonics

# Device and Pulse Scheme

Fig. 1. (A) Scanning electron micrograph of a device with the same gate structure as the one used in this experiment.

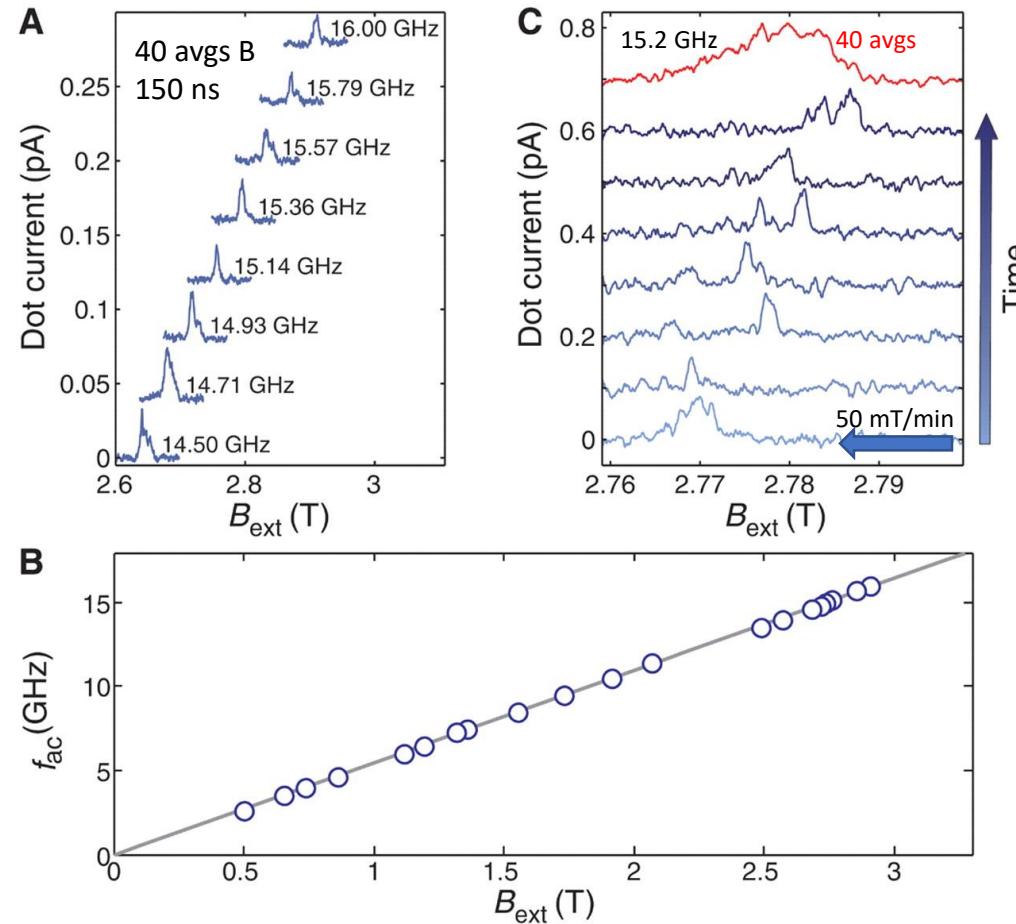


# Using Pauli Spin Blockade for Spin Readout

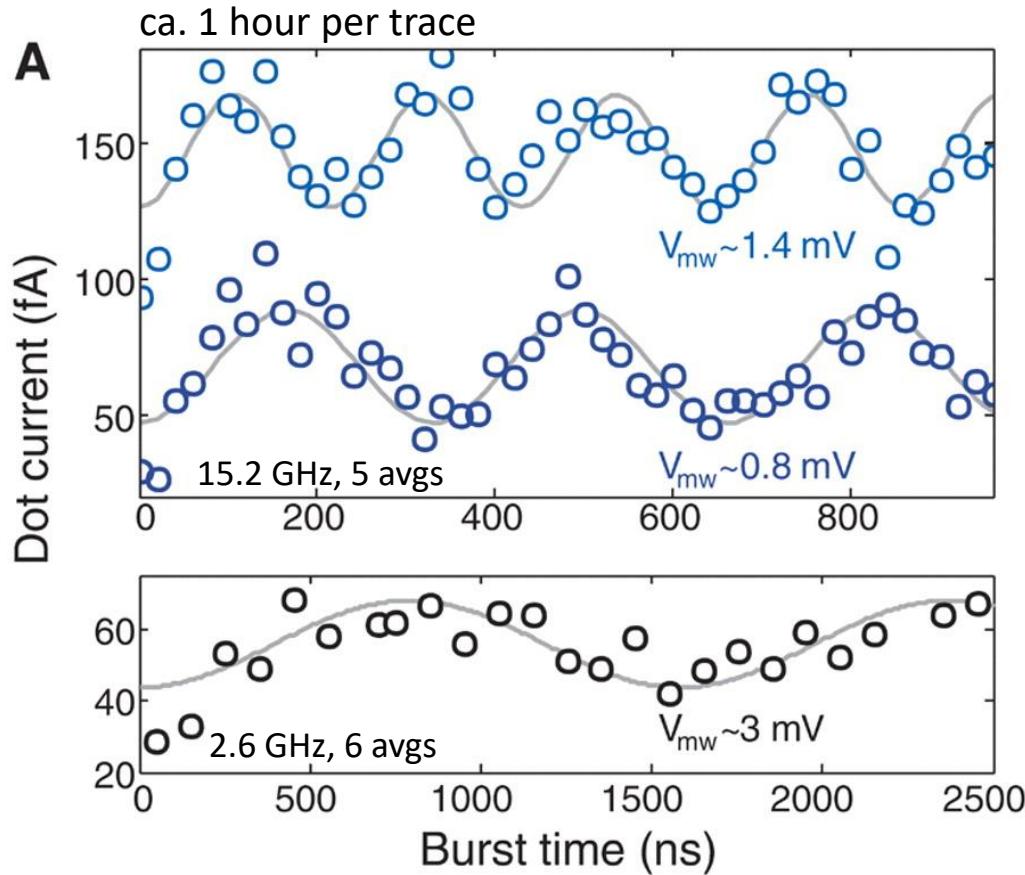


# ESR signature and B-field noise

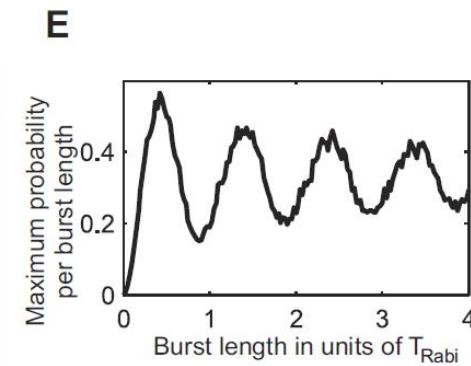
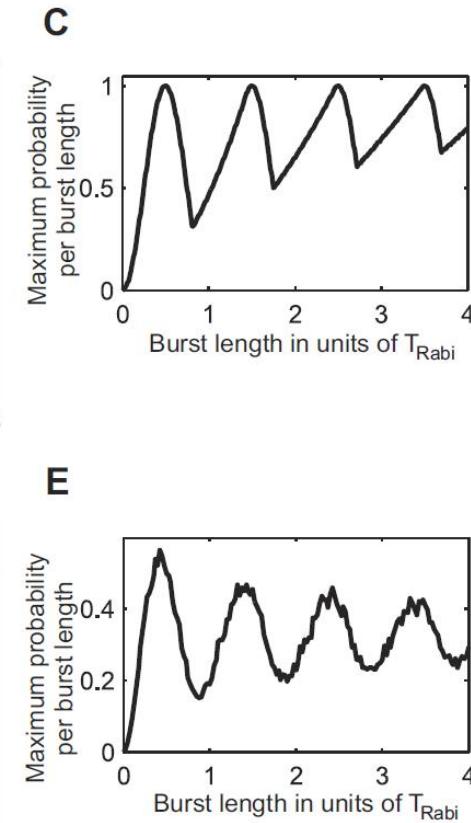
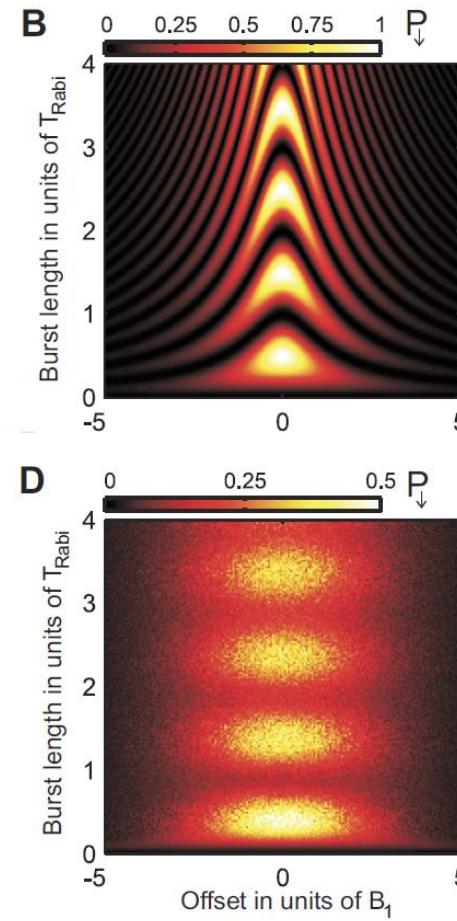
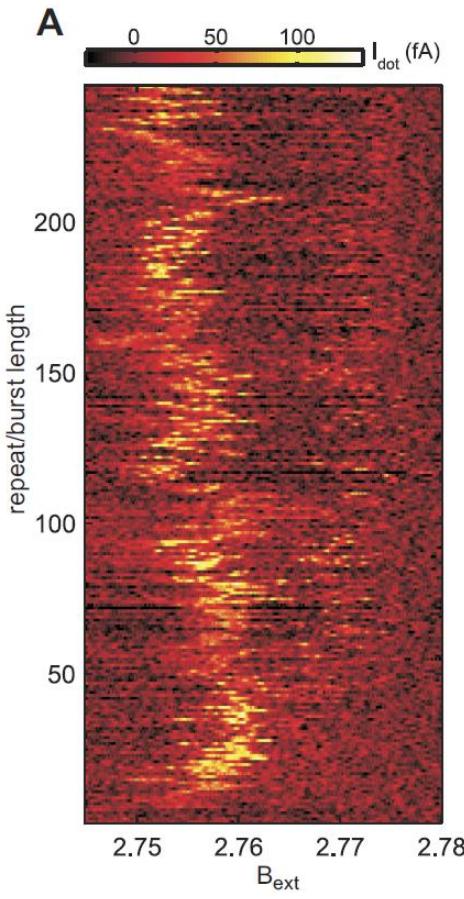
**Fig. 2. (A)** The current averaged over 40 magnetic field sweeps is given for eight different excitation frequencies, with a microwave burst length of 150 ns.



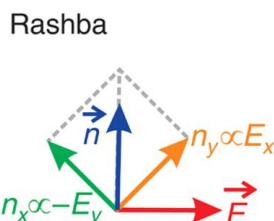
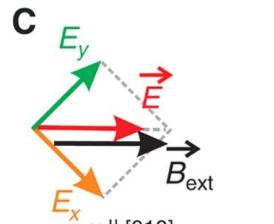
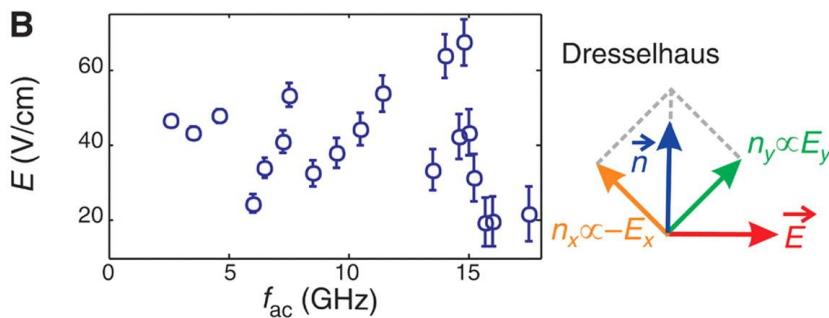
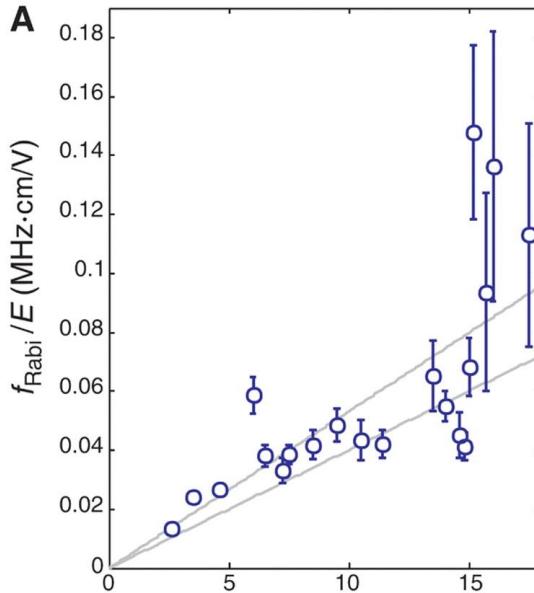
# Rabi Oscillations



# Chevrons with B-field noise



# Determining the EDSR Mechanism

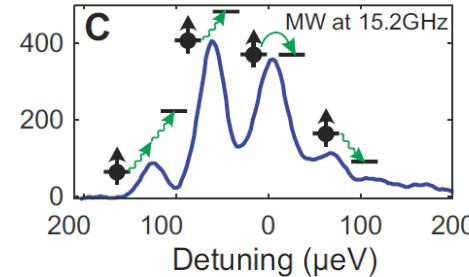
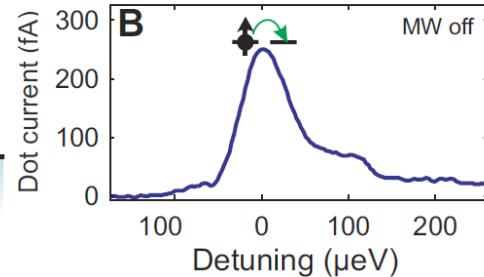
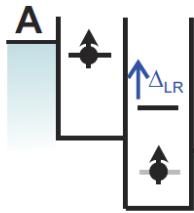


$$|\mathbf{B}_{\text{eff}}(t)| = 2 |\mathbf{B}_{\text{ext}}| \frac{l_{\text{dot}} e |\mathbf{E}(t)| l_{\text{dot}}}{l_{\text{SO}}} \frac{1}{\Delta}$$

dot orbitals     $\Delta, l_{\text{dot}}$   
spin-orbit     $\ell_{\text{SO}}$   
electric field     $E(t)$

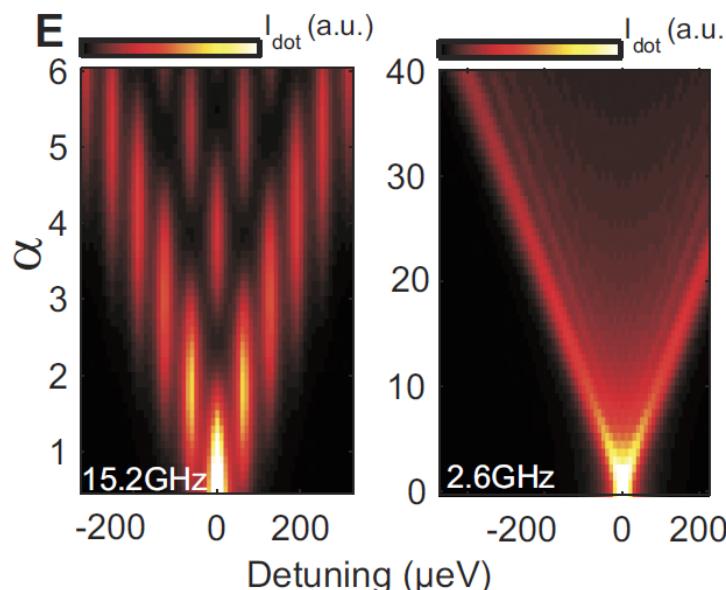
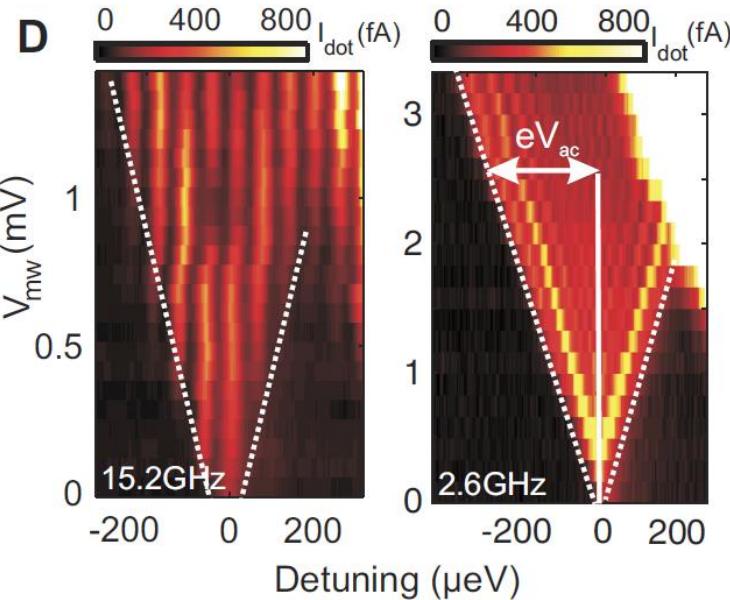
challenge: E-field vs. freq.

# Calibrating the E-field: Photon Assisted Tunneling


$$\alpha = (eV_{\text{ac}})/hf_{\text{ac}}$$
$$\tilde{\Gamma}(E) = \sum_{n=-\infty}^{+\infty} J_n^2(\alpha) \Gamma(E + nhf_{\text{ac}})$$

n-th order Bessel funct.

envelope ca.  $\alpha \sim n$



# Conclusions

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Demonstration of electric dipole spin resonance

- Rabi frequencies up to 5 MHz / 200 ns
- Consistent with B-scaling as expected for spin-orbit
- Nuclear spins clearly visible

Improvements, outlook:

- suppressing PAT
- higher B-fields
- optimized gate layout
- materials with stronger spin-orbit coupling
- individual addressing of spin in array
- single shot read out