

Clean ballistic quantum point contact in SrTiO₃

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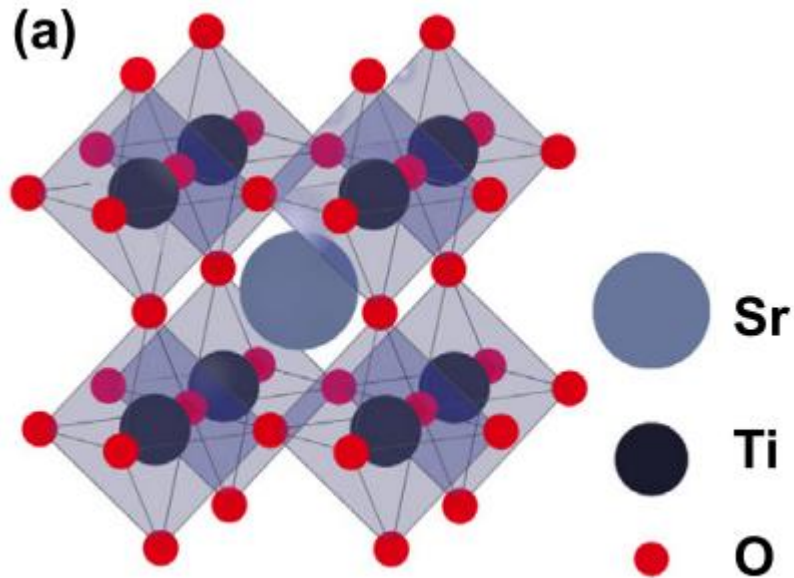
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Outline

- Why you should care about SrTiO_3
- Device design
- Data, results
- Conclusion/discussion

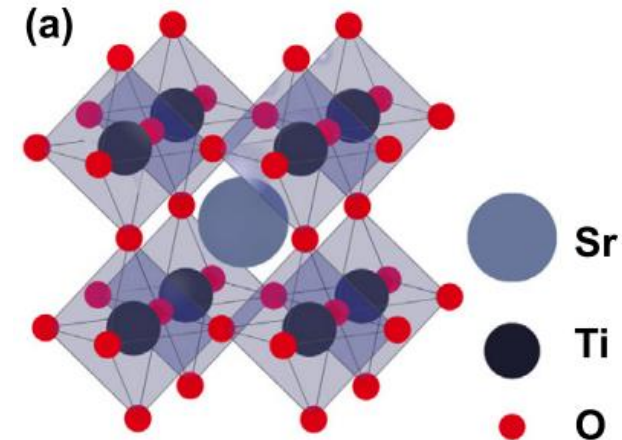


Strontium Titanate

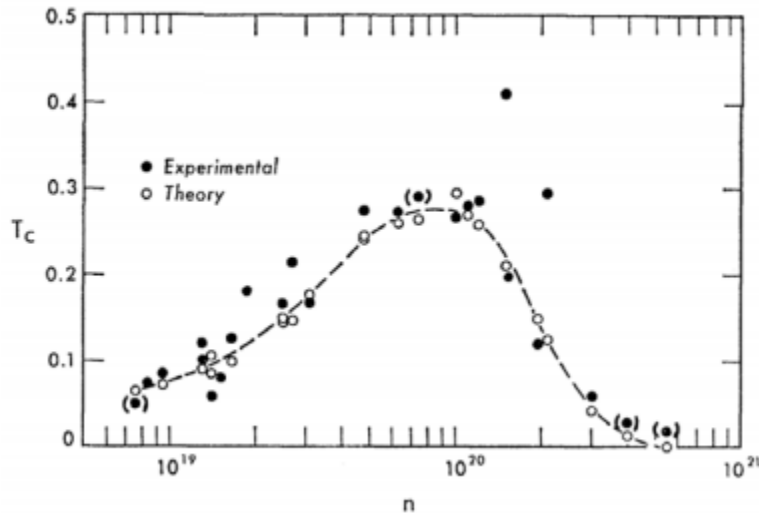


- Centrosymmetric (has inversion center)
- no DSOI
- piezoelectric effect (all centrosymmetric materials)
- Paraelectric
- Indirect band gap 3.25 eV
- Direct band gap 3.75 eV
- $\epsilon_r \sim 300$ at room temp.
- Insulator-to-metal transition in LaAlO₃/SrTiO₃ interfaces

Low Temp

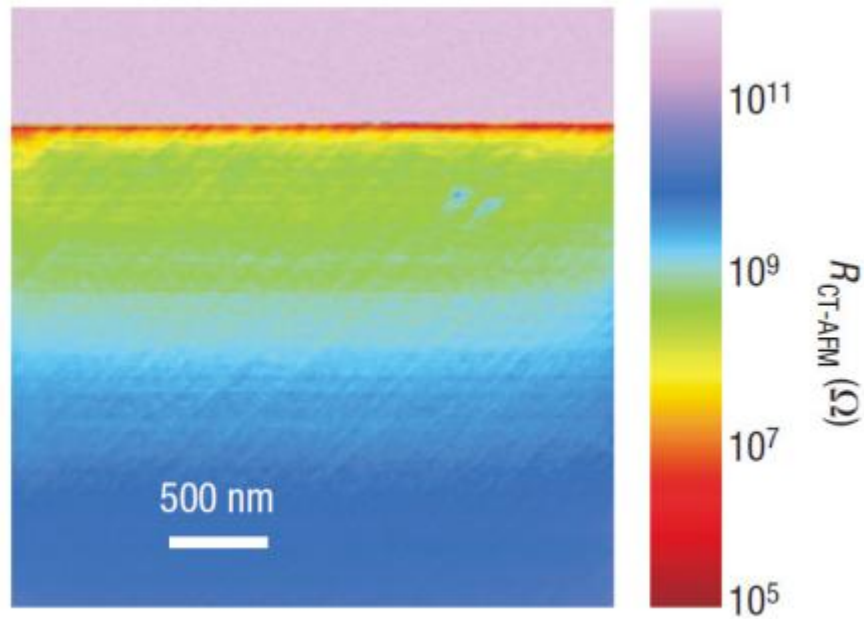


- Dielectric constant increases as temp drops
→ saturates around 4 K near $\epsilon_r = 10^5$
- Density dependent superconductivity below 450 mK
- Approaches ferroelectric phase transition but remains paraelectric due to quantum fluctuations
- Jahn-Teller effect manifests at low temp
→ molecules distort geometrically to remove degenerate ground-state, thus lowering overall energy



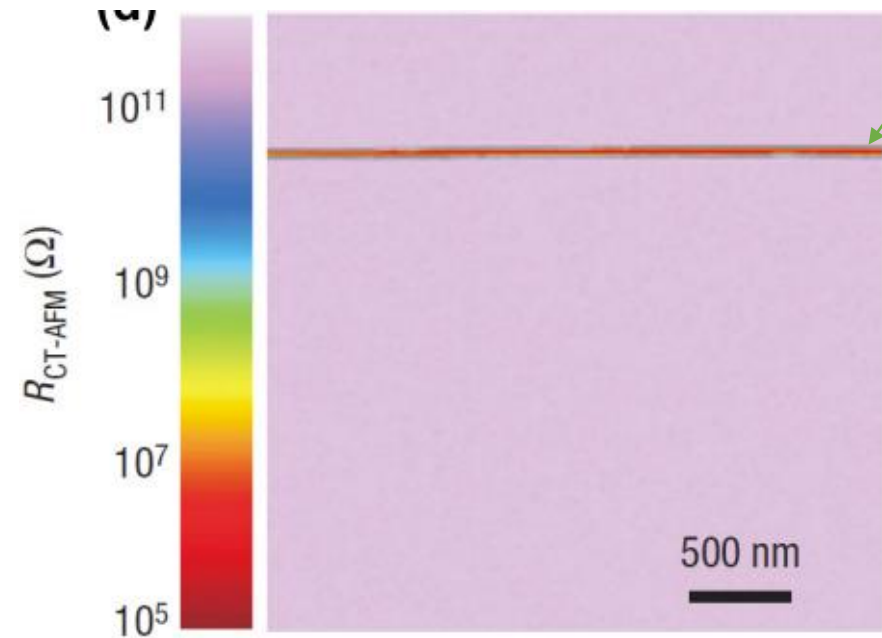
Will Anneal for 2DEGs

Take bulk, standard stock sample, anneal in high oxygen partial pressure
(here $\text{LaAlO}_3/\text{SrTiO}_3$)



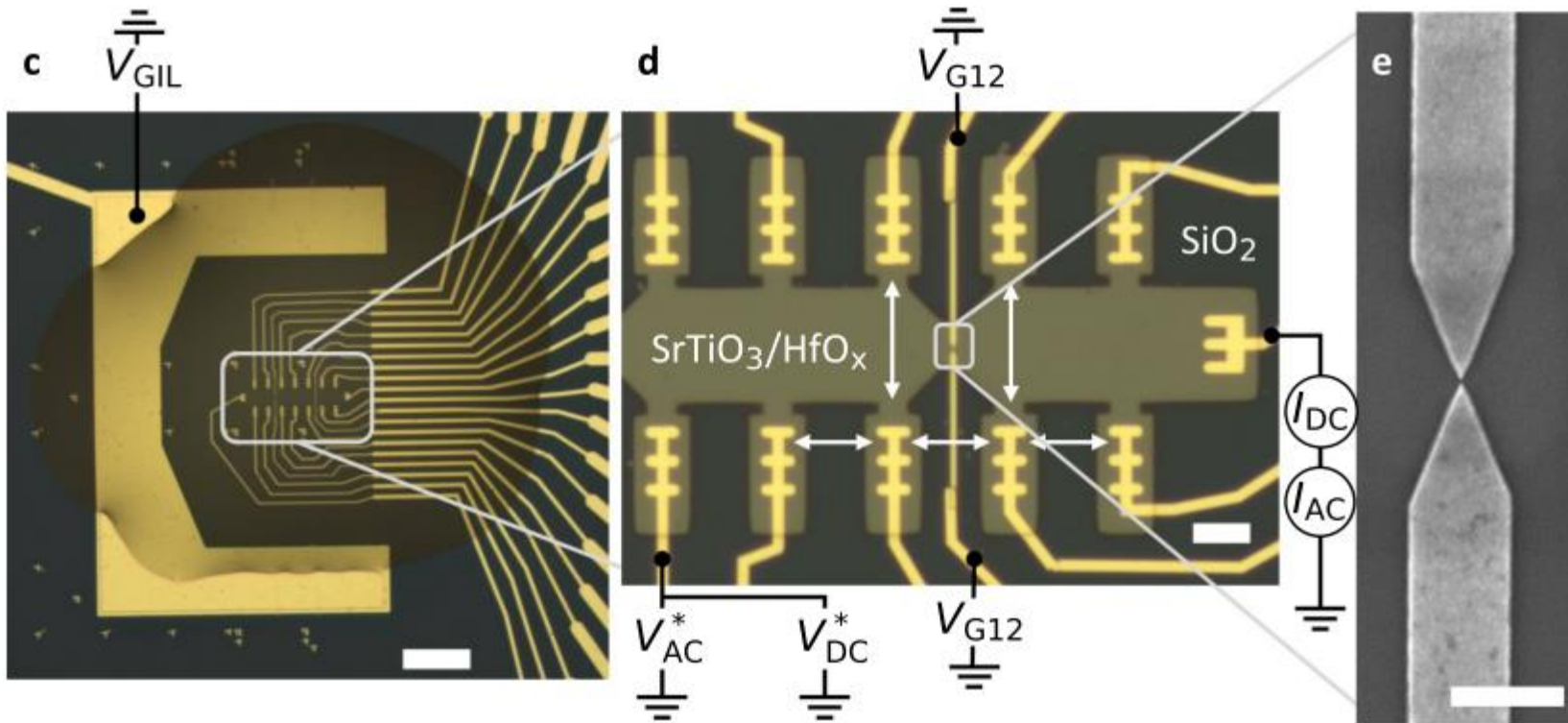
unannealed

(conductive-tip AFM probing electron profile)

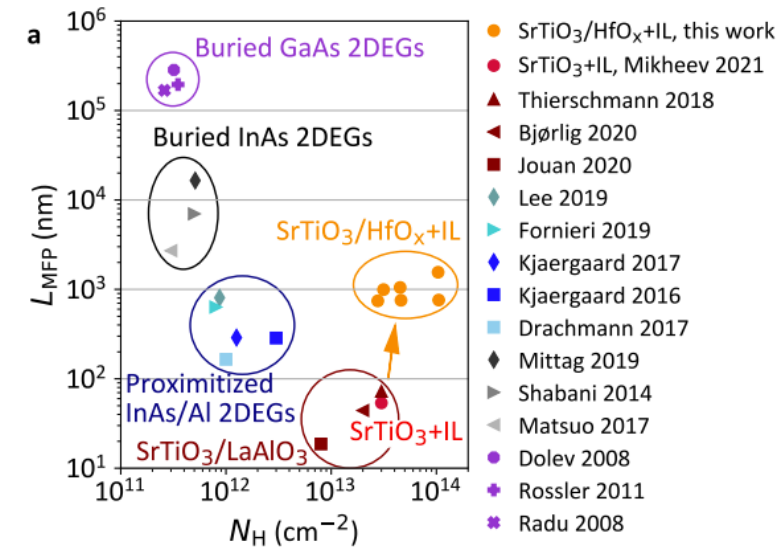


annealed

Device

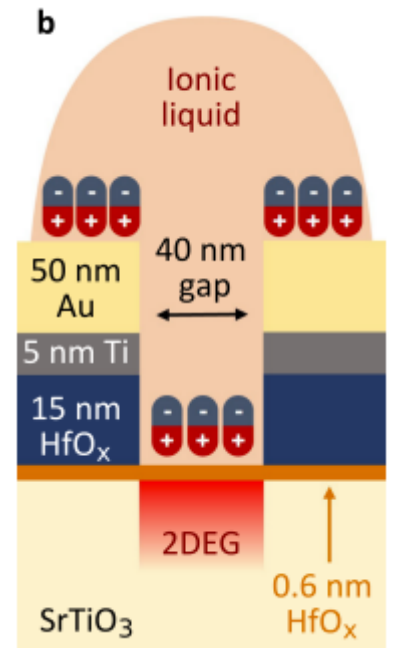
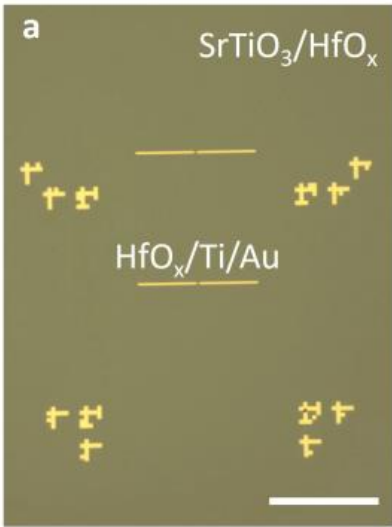


Scale bars:
c: 100 μm | d: 10 μm | e: 1 μm

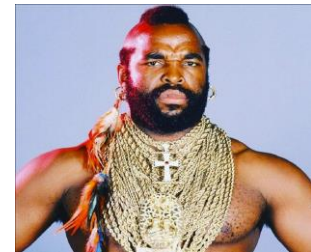


Fabrication

Fabricated on commercial (001) SrTiO₃ single crystal substrates



- Soak in DI water for 20 min, then annealed at 1000 C for 2 hours in Ar/O₂
- HfO_x (4 layers) deposited (85 C)
- (a) Selective patterning for HfO_x (15 nm) and Ti/Au (5/50 nm) for QPC tip
- (b) Ebeam + Ti/Au (40/100) for remaining patterning
- (c) Ebeam, Ar milling, Ti/Au (10/80)
- (d) Ebeam + sputtering of 80 nm SiO₂
- (e) ionic liquid (IL) drop (DEME-TFSI)
- Use IL to tune density before cooling below its freezing point (200 C) and doing measurements

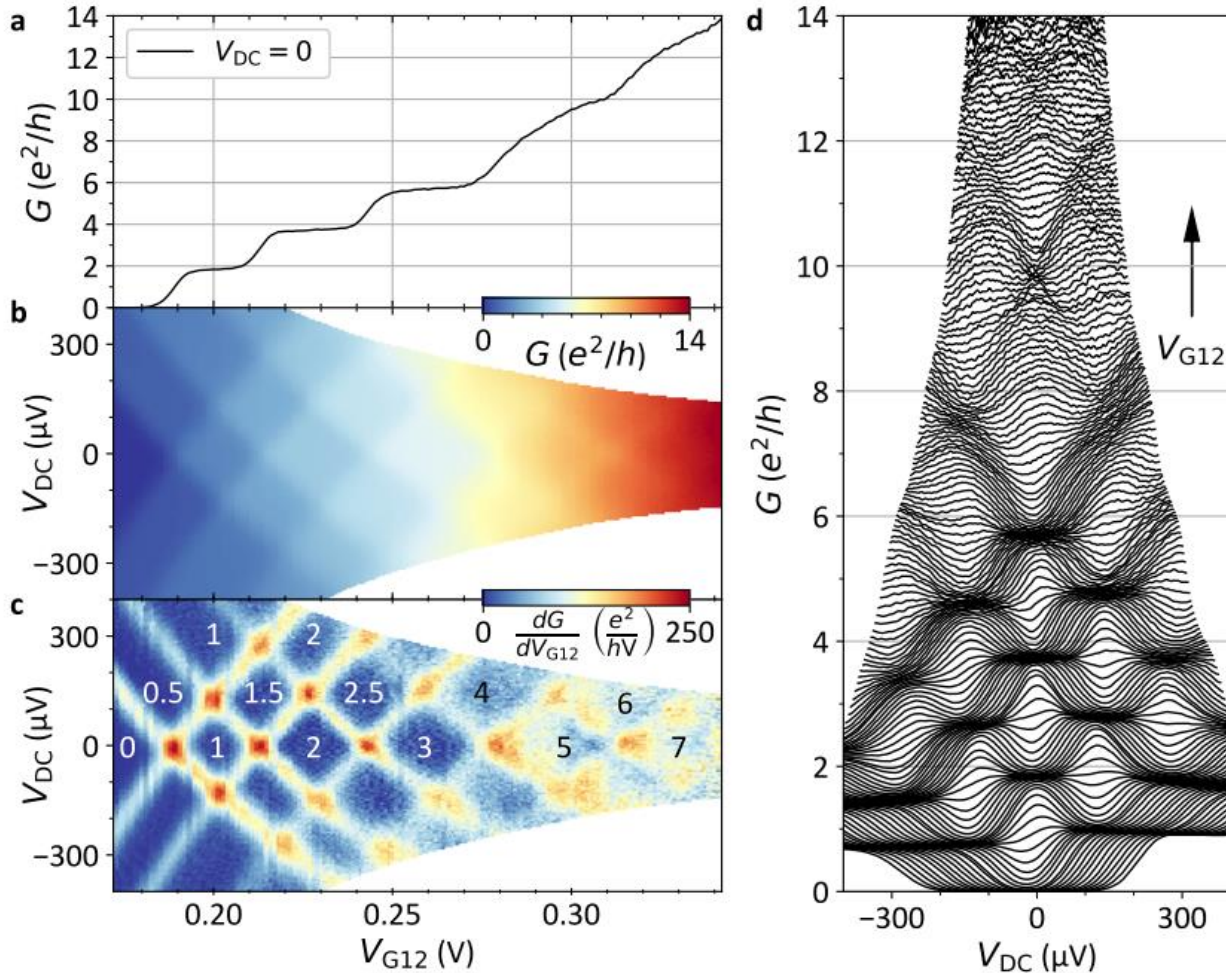


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Textbook* quantized conduction

4pt measurements, no series resistance subtracted

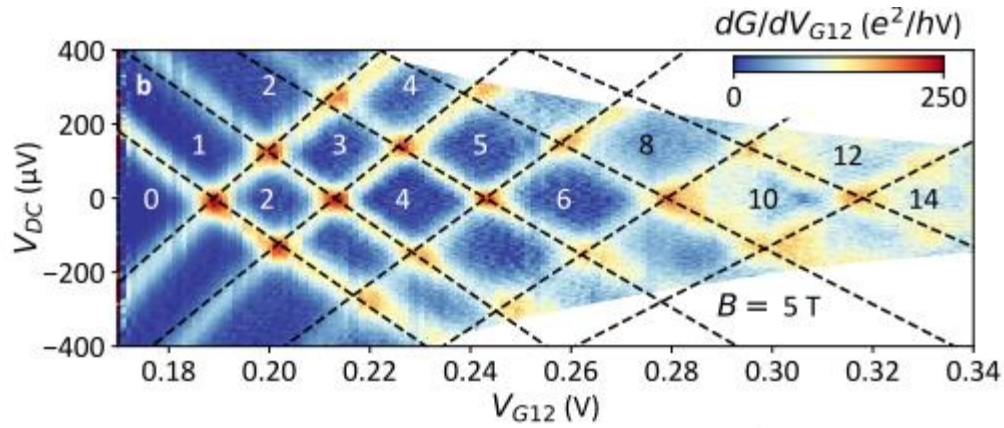


- Doubly degenerate subbands
- Some plateaus skipped

32 mK *B = 5 T (!!)

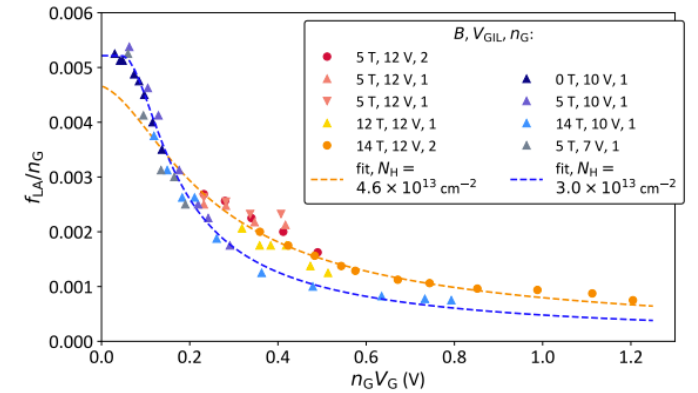
Lever Arm Non-Uniformity

$f_{LA} = dV_{DC}/2dV_G$ at dG/dV_G peaks

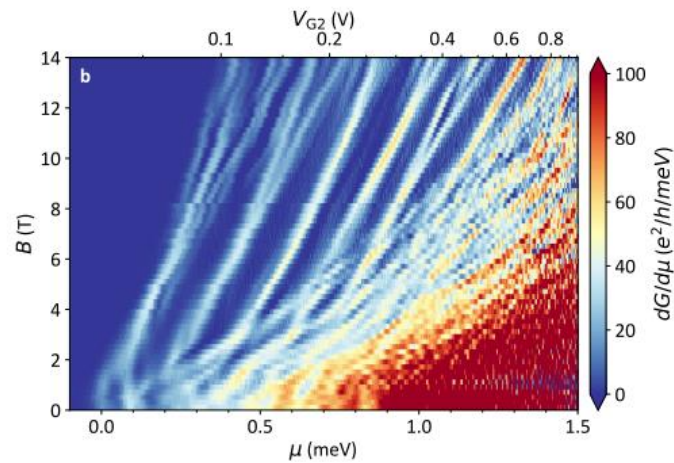
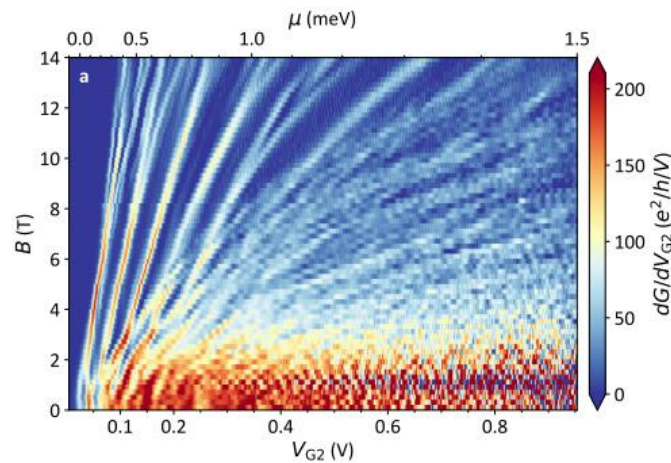


$$f_{LA}(V_G) = \frac{f_{LA}(0) \cdot V_{NL}}{\sqrt{V_{NL}^2 + (V_G - V_{G0})^2}}$$

Sharpness

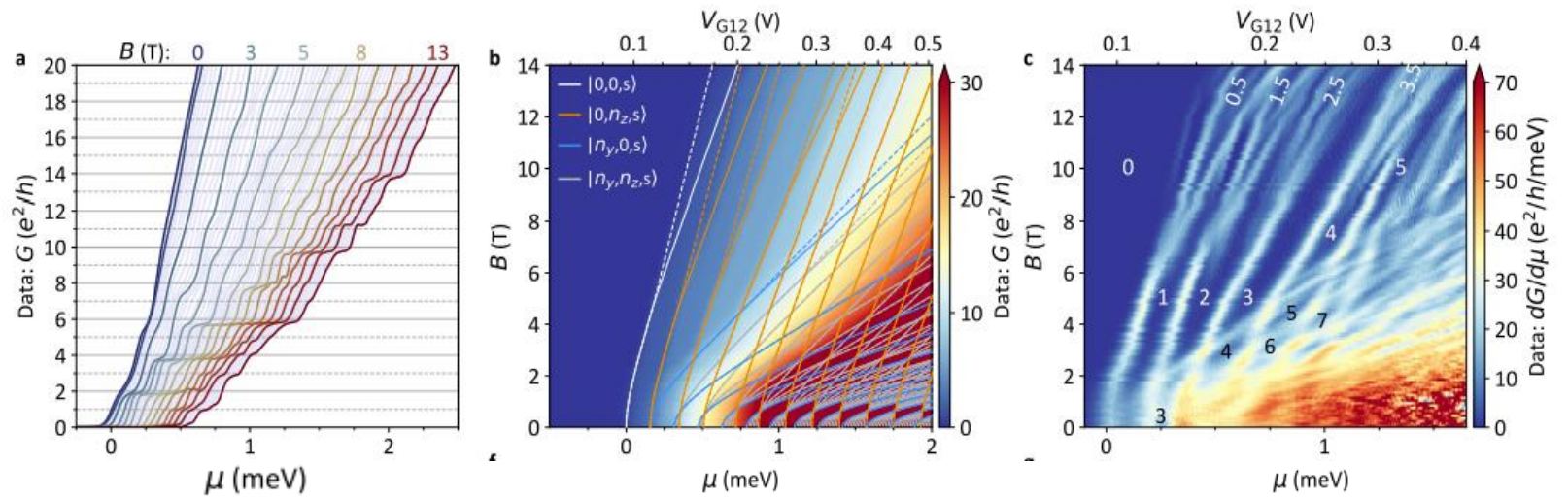


$$\mu(V_G) = \int_{V_{G0}}^{V_G} f_{LA}(V_G) \cdot (V_G - V_{G0}) dV_G = f_{LA}(0) \cdot V_{NL} \arctan \left(\frac{V_G - V_{G0}}{\sqrt{V_{NL}^2 + (V_G - V_{G0})^2}} \right)$$



Subband Evolution

- 2-fold degeneracy persists til ~ 7 T
- Fast and slow moving subbands
 - Cross resulting in intermittent 4-fold degeneracies



3D QPC Saddle Potential

- Extend classic 2D model to 3D
- Quadratic in x, y, z with polarity $P_x = -1, P_{y,z} = 1$

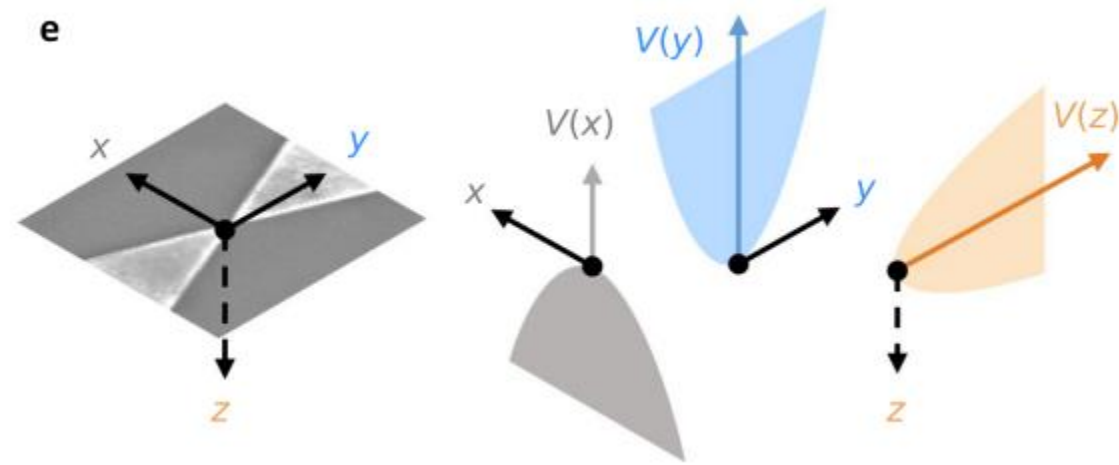
$$\mathcal{H} = \sum_{u=x,y,z} \left(-\frac{\hbar^2}{2m_u^*} \cdot \frac{\partial^2}{\partial u^2} + P_u \frac{m_u^* \epsilon_u^2 u^2}{2\hbar^2} \right) + E_Z \sigma_z$$

- At $B = 0$, $\epsilon_u(B = 0) = \frac{\hbar^2}{(m_u^* l_u^2)}$
- At $B \neq 0$, cyclotron frequency renormalizes x-y plane confinement:

$$\epsilon_x = \hbar \omega_x / (1 + \omega_c^2 / \omega_y^2)^{1/2}, \quad \epsilon_y = \hbar (\omega_y^2 + \omega_c^2)^{1/2},$$

ϵ_z unaffected

- Can separate Hamiltonian to y-z subbands discretized as $|n_y, n_z, s\rangle$ + x component which broadens them



Results

- Integers $n_{x,y} \geq 0$ and $s = \pm 1/2$ yield subband spectrum:

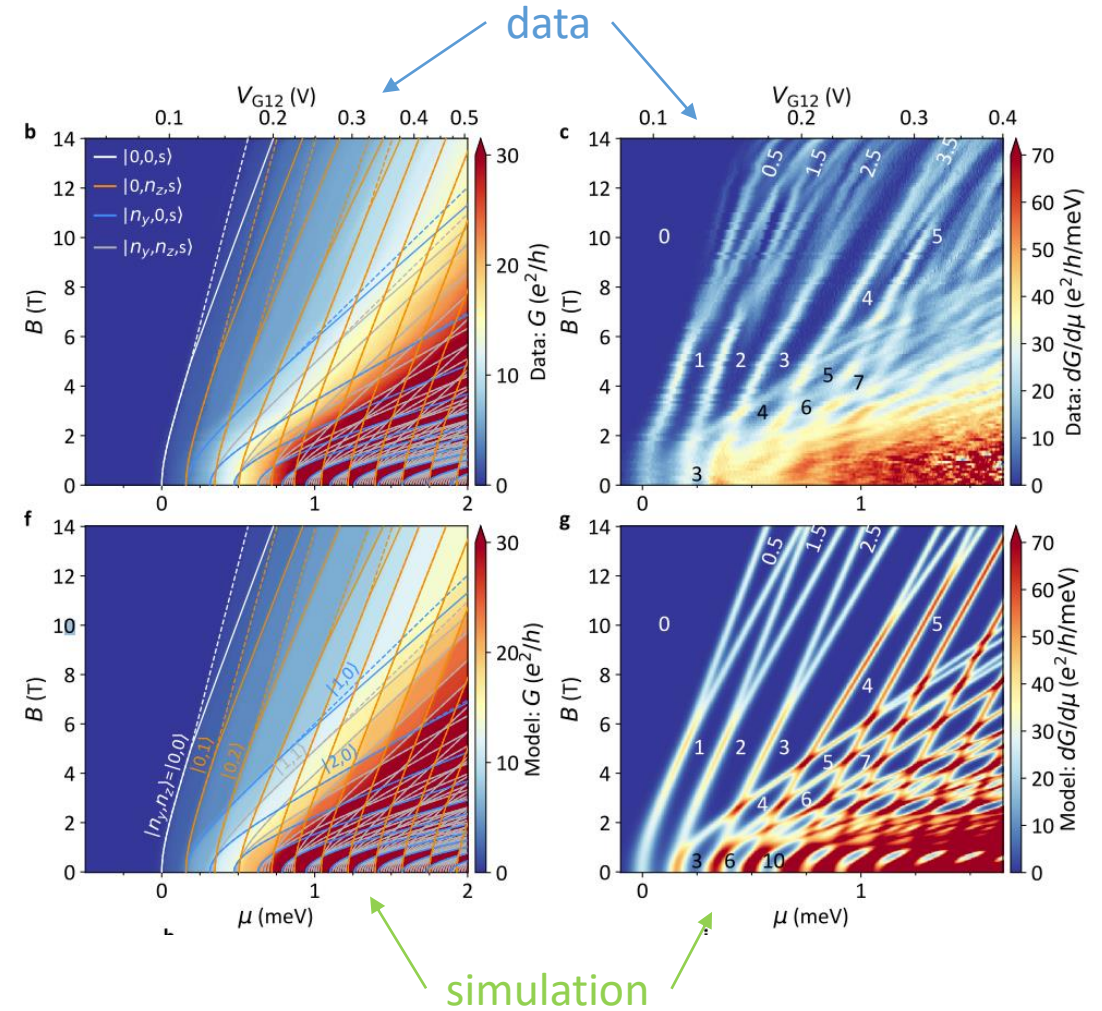
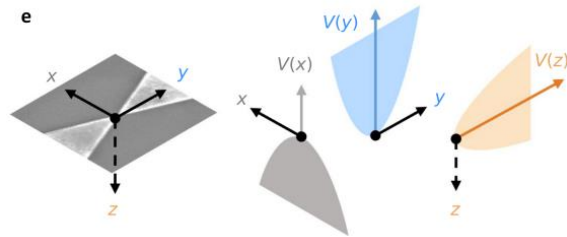
$$\epsilon_{yz} = \epsilon_y \left(n_y + \frac{1}{2} \right) + \epsilon_z \left(n_z + \frac{1}{2} \right) + E_Z(B, s)$$

- Empirically modify

$$E_Z(B, s) = g\mu_B s(B - B_P) \text{ when } B \geq B_P,$$

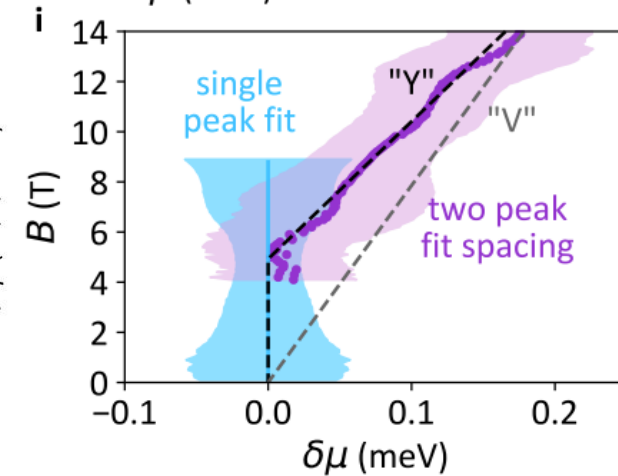
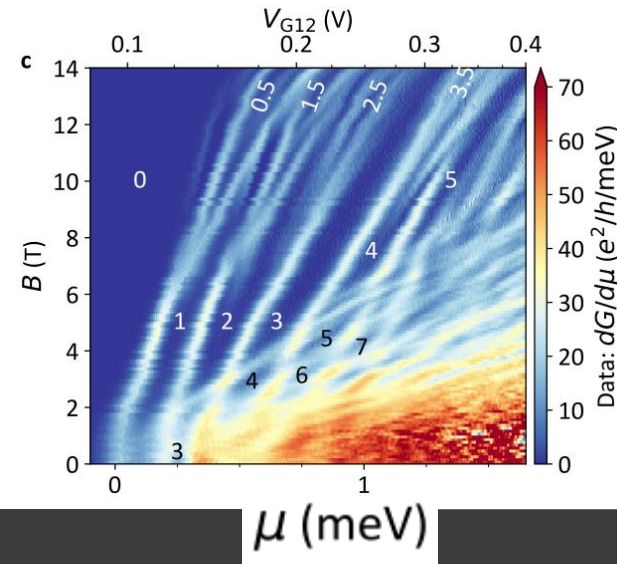
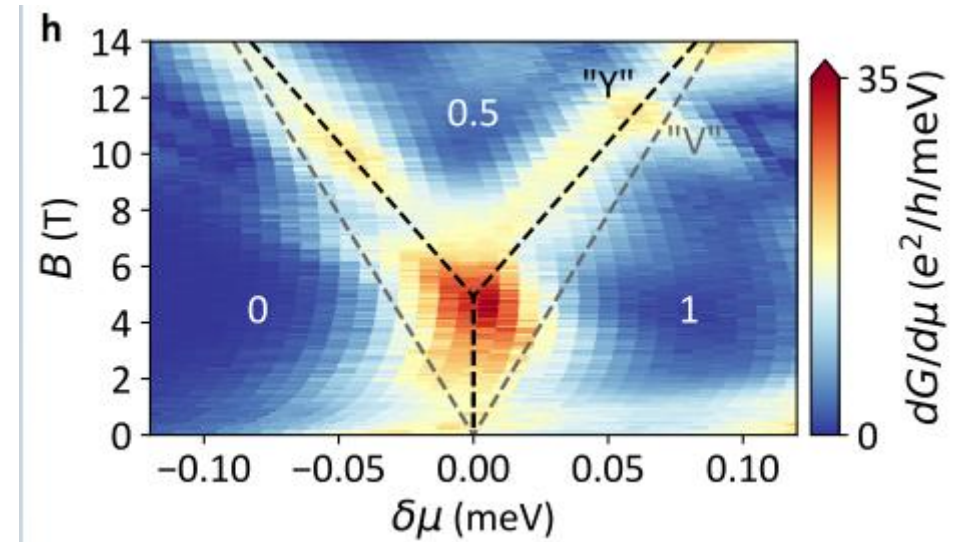
$$E_Z(B, s) = 0 \text{ when } B < B_P$$

- Slow moving subbands have $|n_y = 0, n_z, s = \pm 1/2\rangle$
- Fast subbands have $n_y > 0$, due to renormalization of ϵ_y by cyclotron frequency

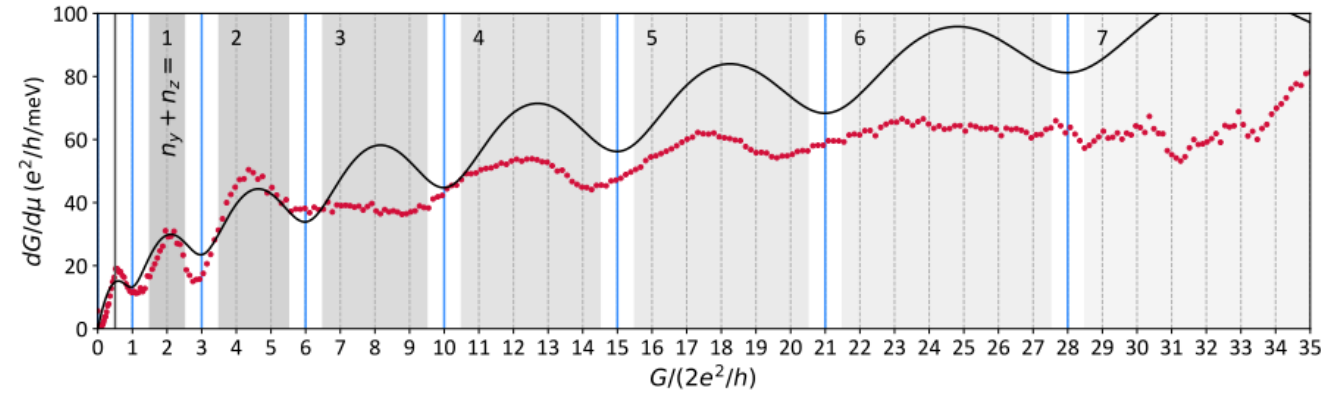
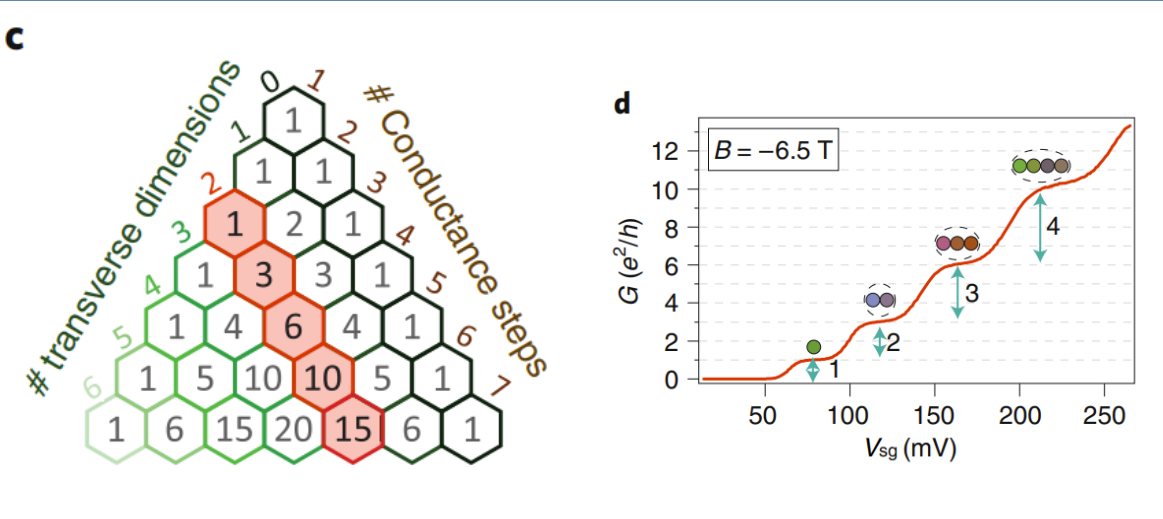


Unusual Spin-Splitting

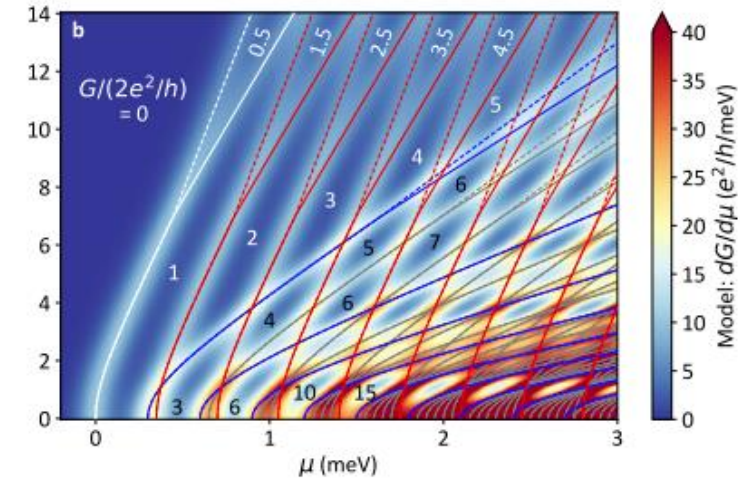
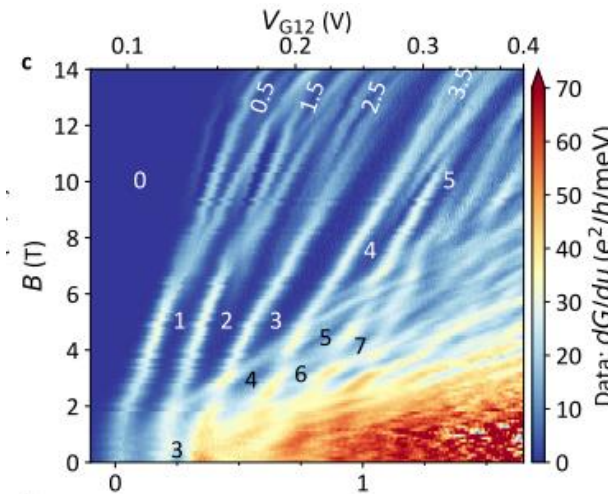
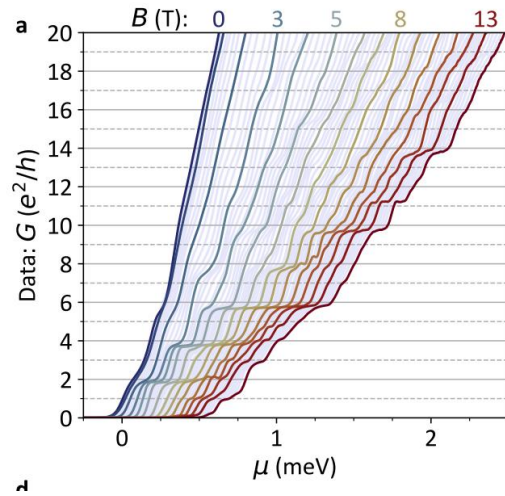
- Modified E_Z results in subband splitting into 'Y' shape
 - Y: $B_P = 4.9$ T, $g = 0.32$
 - V: $B_P = 0$ T, $g = 0.22$
- Appears to not be due to subband broadening
- Possible explanation: attractive "negative-U" interaction between e^-
 - Then: B_P is field when pairing interaction (singlets) balances with E_Z (alignment of spins)
 - Critical field/temp higher than plausible global SC state in 2DEG
 - Pre-formed pairs condensing at low T?
 - Pairing locally enhanced at ferroelastic domain walls?
 - Valence-skipping defects?



Pascal Conductance with $2e^2/h$ Steps?



$$\epsilon_{yz} = \epsilon_y \left(n_y + \frac{1}{2} \right) + \epsilon_z \left(n_z + \frac{1}{2} \right) + E_Z(B, s)$$



That's It

- Interesting material and device architecture
- Unconventional quantized conductance
- Strange attractive e-e effects without global superconductivity

Thanks for listening!!

