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Parity readout of spin qubits in silicon quantum dots

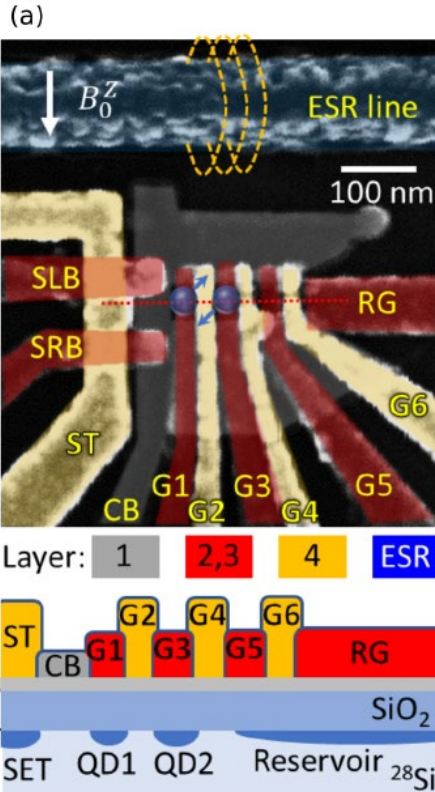
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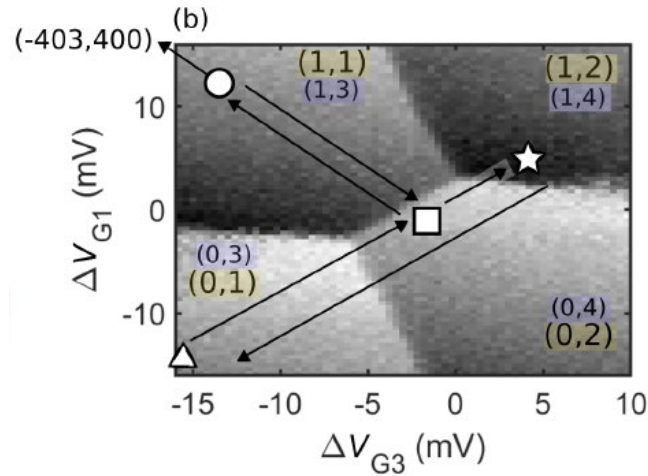


- scalable qubit readout
- Pauli spin blockade breakdown
- parity readout
- quantum error correction

Device and operating regime



- electron accumulation
- 800 ppm ²⁹Si

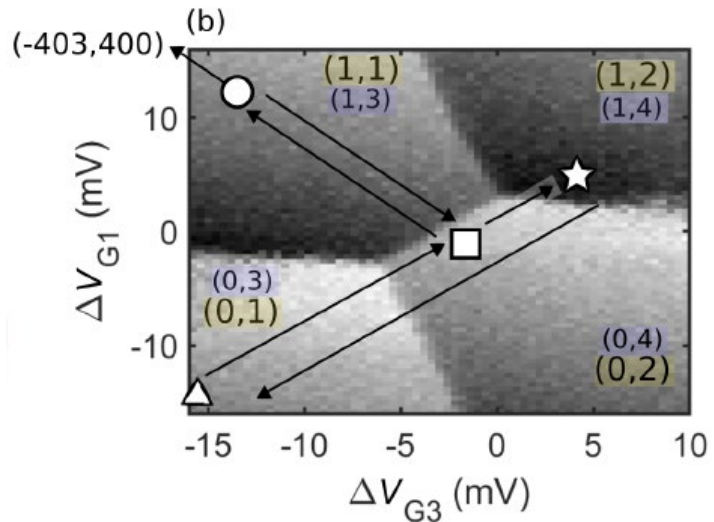
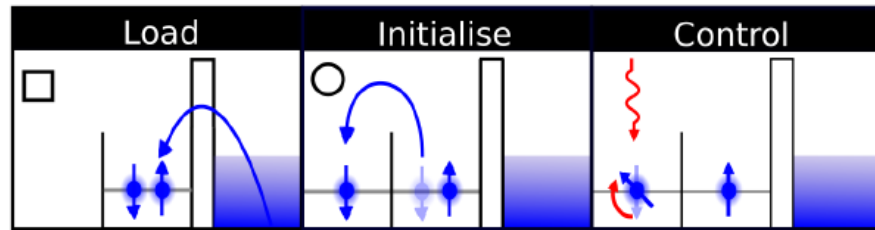


- right dot, two electrons in the lower valley
- the upper valley is >0.1 meV above
- closed shell, no effect on two extra electrons $(0,4) \rightarrow (0,2)$

Measurement scheme

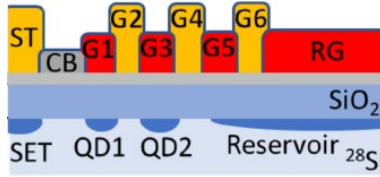
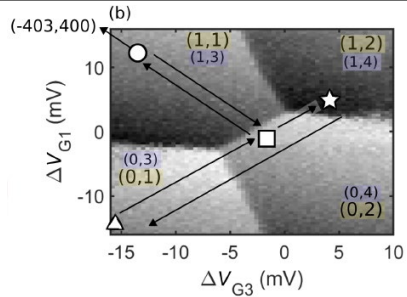


(c)



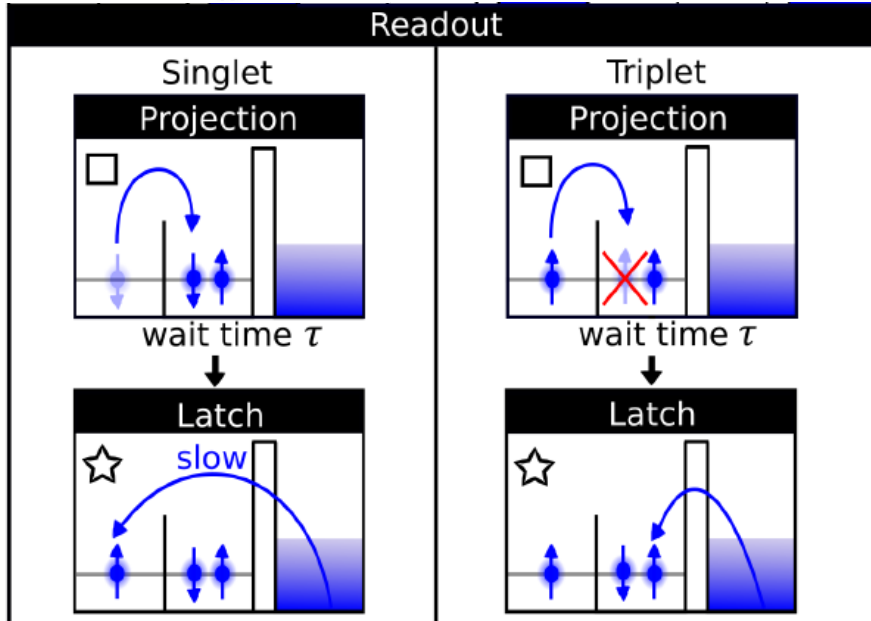
- start at (0,1)
- electron loading into $|S(0,2)\rangle$
- change detuning (0,2) \rightarrow (1,1)
- spin state depends on the ramp rate
anticrossings with $|T(1,1)\rangle$ and $|S(1,1)\rangle$
- manipulation
 - ESR
 - detuning control

Latched single-tripled readout



standard readout

- return to (0,2), □
- only singlet would tunnel
- charge reflects the spin state immediately after reaching □

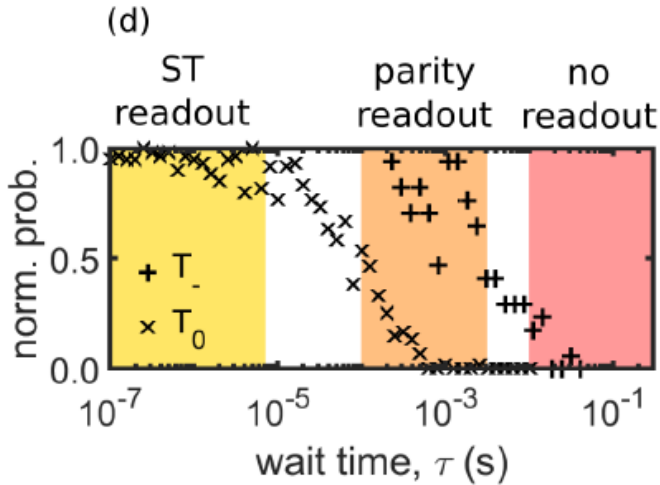


slow charge sensing compromises the spin state

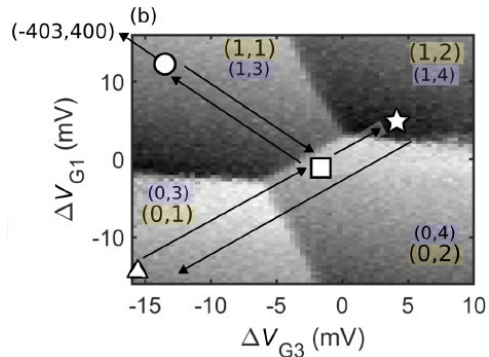
To study the lifetime of the Pauli blockade:

- wait time τ in the (0,2)
- move to the latched region (1,2)
 - singlet will stay in (0,2)
 - slow tunneling into the left dot
 - triplet will go to (1,2)
- charge state maps the spin state at the moment of latching

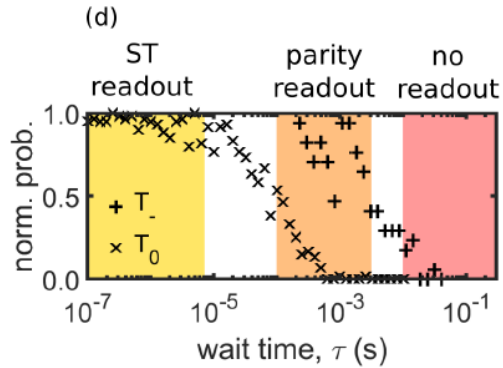
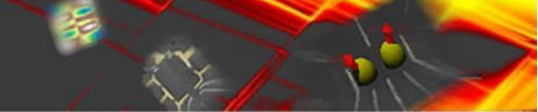
Parity readout



- Triplet states will eventually decay to $|S(0,2)\rangle$
- odd-parity $|T_0\rangle$ decays in 200 μ s
- $|T_-\rangle$ decays in 5 ms
- $|T_+\rangle$ is also long
- if the time for the charge sensing at \square is in orange region, odd-parity states ($|S\rangle$, $|T_0\rangle$) are indistinguishable
- even-parity states ($|T_+\rangle$, $|T_-\rangle$) can be distinguished

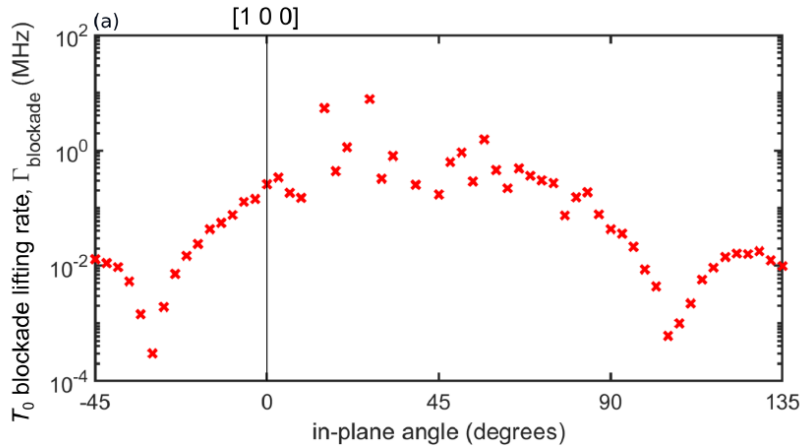


Measure $\Gamma_{blockade}$



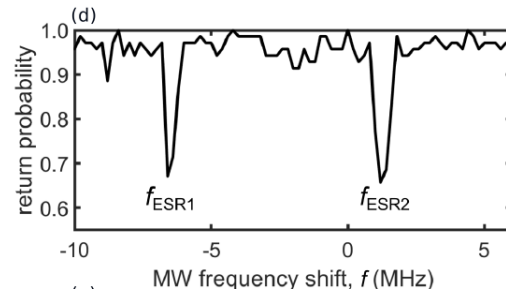
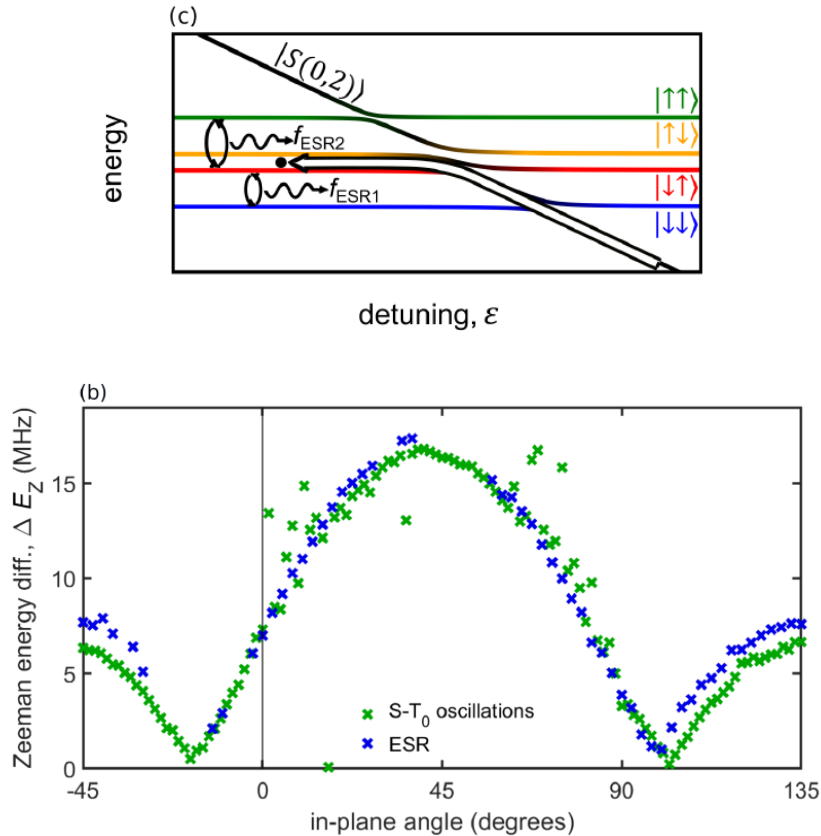
- choose the direction of B
- fit the blockade lifting for the state $|T_0\rangle$
can change by 3-4 orders of magnitude

Look for the correlation between $\Gamma_{blockade}$ and $\Delta E_Z = \Delta g \mu_B B$



- Δg – due to spin-orbit coupling induced by the interface
- large Dresselhaus component, controlled by the direction of B

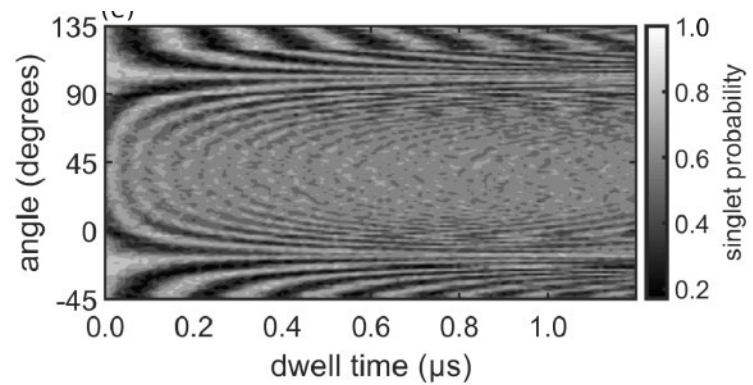
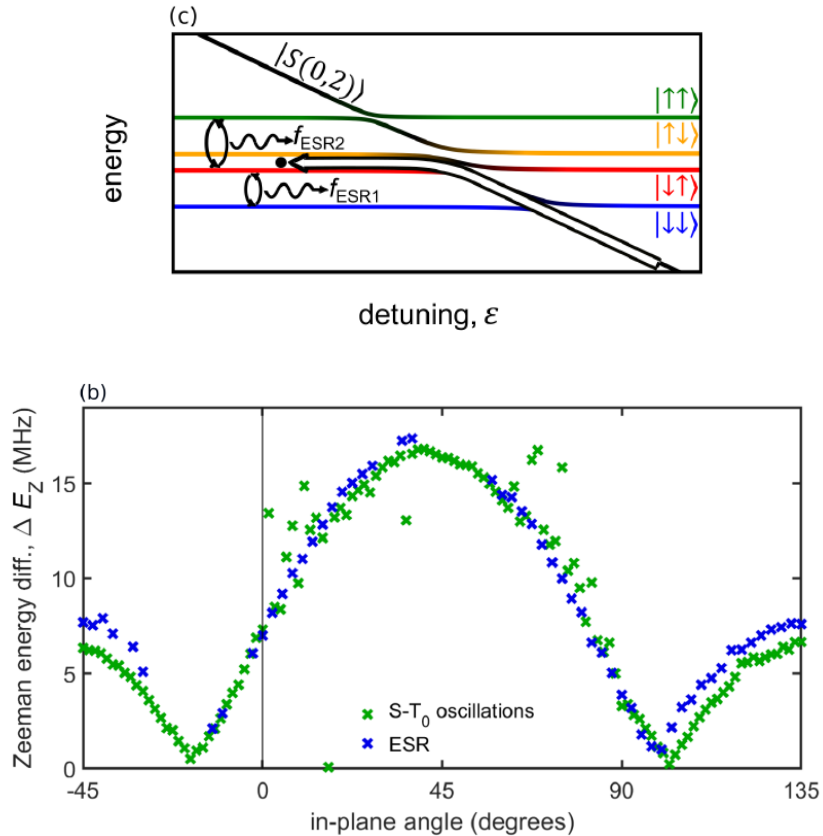
Measure ΔE_Z using ESR



electron spin resonance, rotate one spin

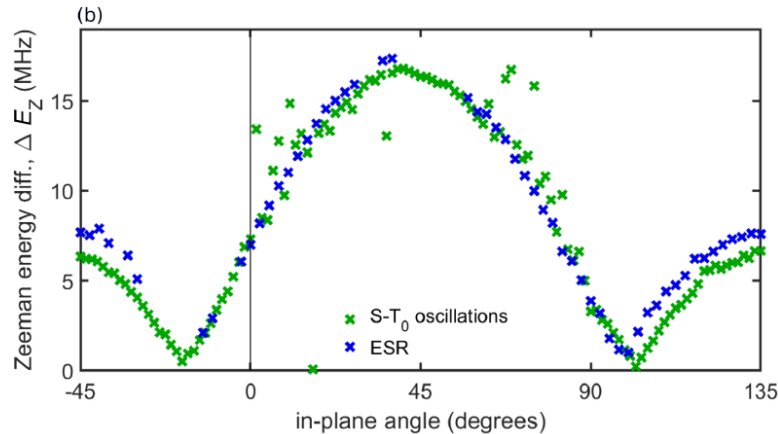
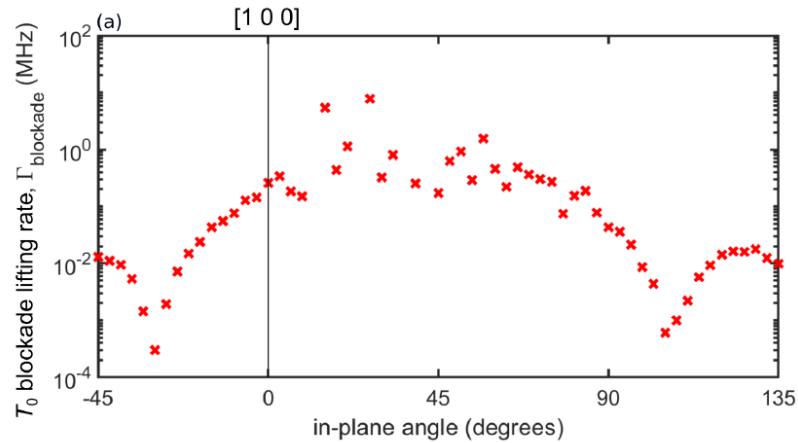
- initialize in $|\downarrow\uparrow\rangle$
 - go adiabatically through $|\downarrow\uparrow\rangle - |\uparrow\downarrow\rangle$ anticrossing
- apply microwave pulse
 - f_{ESR1} to rotate $|\downarrow\uparrow\rangle$ to $|\uparrow\uparrow\rangle$
 - f_{ESR2} to rotate $|\downarrow\uparrow\rangle$ to $|\downarrow\downarrow\rangle$
- return to (0,2), dip in the return probability at the resonance frequency
- difference between f_{ESR1} and f_{ESR2} gives ΔE_Z if $\Delta E_Z \geq J$, $|T_0(1,1)\rangle - |S(1,1)\rangle$ splitting

Measure ΔE_Z using S-T oscillations

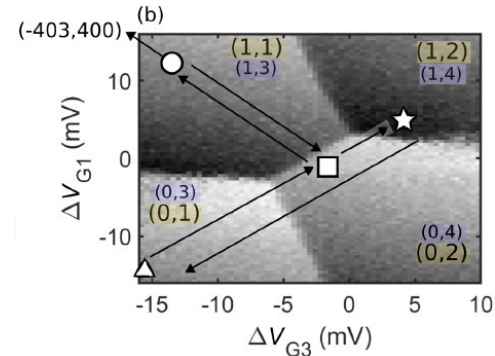


- oscillations between $|S(1,1)\rangle$ and $|T_0(1,1)\rangle$ with frequency ΔE_Z
- initialize in $|S(1,1)\rangle$
ramp quickly through (0,2)-(1,1)
- dwell time – time spent at (1,1) configuration

Measure ΔE_Z using S-T oscillations



- excellent agreement between these two ΔE_Z extraction methods
- horizontal shift due to Stark shift
 - ΔE_Z was measured deep in (1,1)
 - $\Gamma_{blockade}$ was measured at (0,2)
- $\Gamma_{blockade}$ increases as ΔE_Z increases



$$\hat{H} = \begin{pmatrix} -\varepsilon & t & 0 \\ t & 0 & \Delta E_Z \\ 0 & \Delta E_Z & 0 \end{pmatrix} \begin{pmatrix} |S(0,2)\rangle \\ |S(1,1)\rangle \\ |T_0(1,1)\rangle \end{pmatrix}$$

$$\frac{d\hat{\rho}}{dt} = -i[\hat{H}, \hat{\rho}] + \hat{\mathcal{L}}[\hat{a}](\hat{\rho}).$$

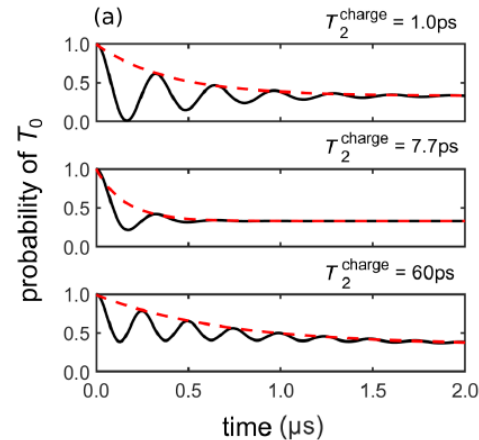
master equation for the density matrix

$$\hat{\mathcal{L}}[\hat{a}](\hat{\rho}) = \hat{a}\hat{\rho}\hat{a}^\dagger - \frac{1}{2}(\hat{a}\hat{a}^\dagger\hat{\rho} + \hat{\rho}\hat{a}\hat{a}^\dagger).$$

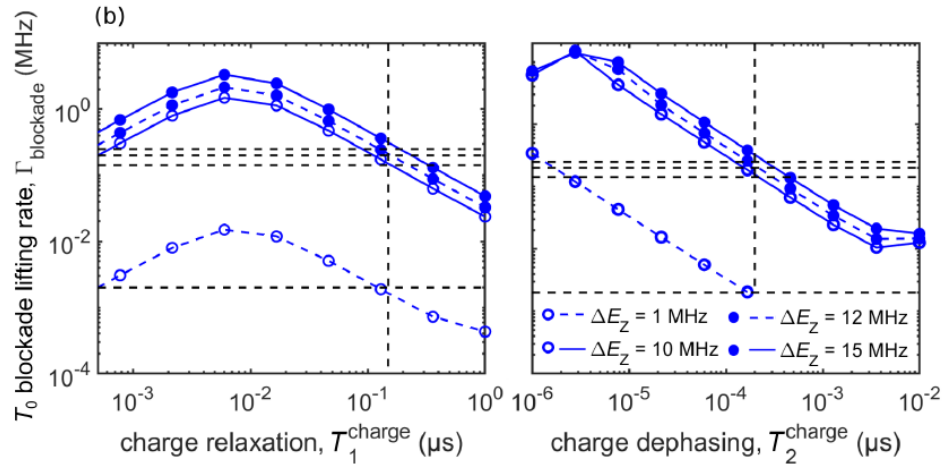
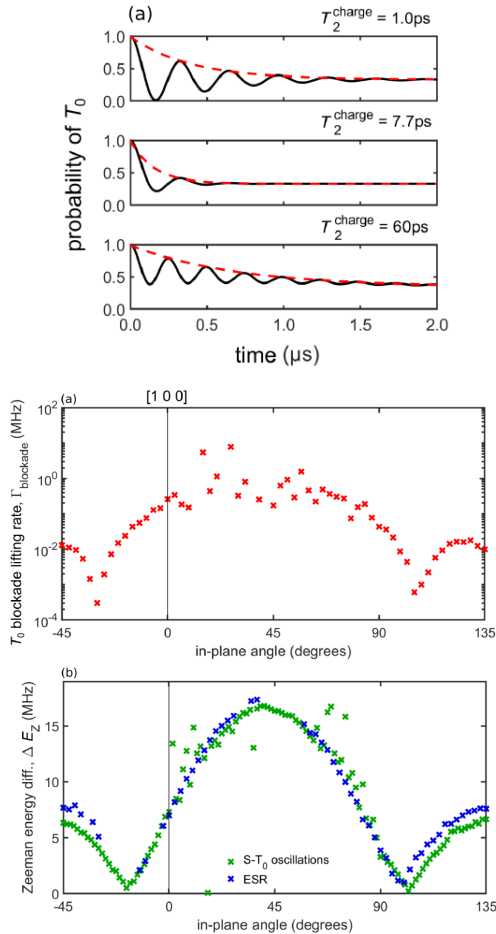
$$\hat{a}_{\text{dephasing}} = \begin{pmatrix} \frac{1}{\sqrt{2T_2}} & 0 & 0 \\ 0 & -\frac{1}{\sqrt{2T_2}} & 0 \\ 0 & 0 & -\frac{1}{\sqrt{2T_2}} \end{pmatrix}$$

$$\hat{a}_{\text{relaxation}} = \begin{pmatrix} 0 & \frac{1}{\sqrt{T_1}} & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

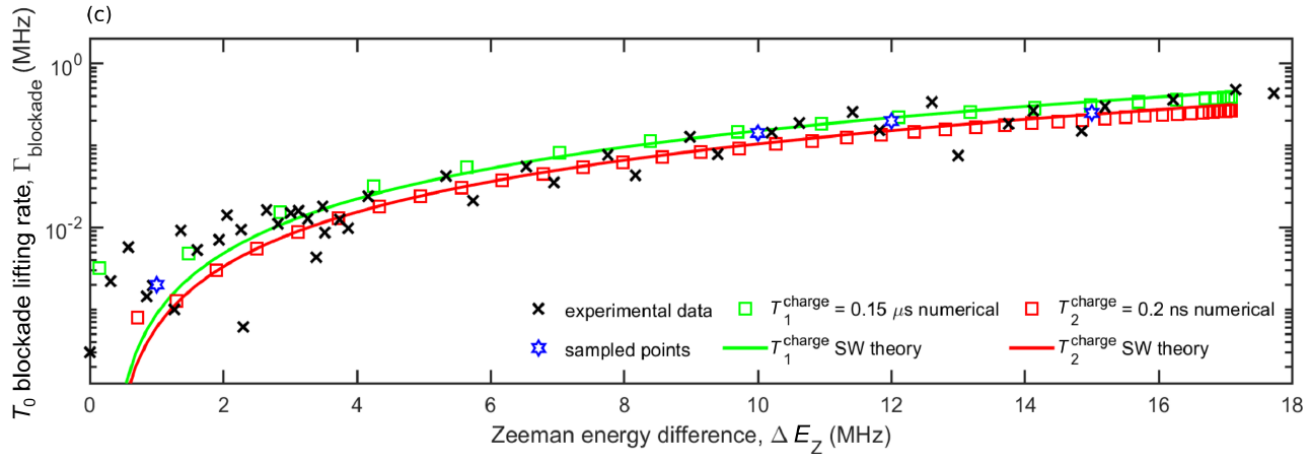
- $|S(0,2)\rangle$ and $|S(1,1)\rangle$ coupled with t
- $|S(1,1)\rangle$ and $|T_0(1,1)\rangle$ coupled through ΔE_Z
- charge dephasing, ε fluctuations, T_2
- charge relaxations, $|S(1,1)\rangle$ into $|S(0,2)\rangle$, T_1
- initialise in pure $|T_0(1,1)\rangle$, oscillations with exponential decay $Ae^{-t\Gamma_{\text{blockade}}} + B$



Charge relaxation and dephasing times



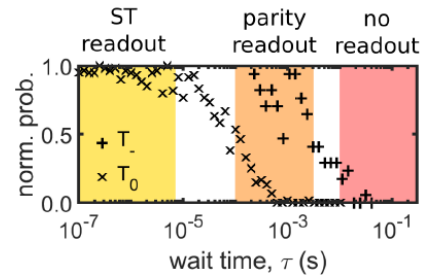
- T_1^{charge} and T_2^{charge} are unknown
- fit: $T_1^{\text{charge}} = 0.2 \mu\text{s}$ and $T_2^{\text{charge}} = 0.2 \text{ ns}$
- T_2^{charge} fits better with the literature – dephasing may be the decay mechanism



$$\Gamma_{\text{blockade}} \approx \frac{2\Delta E_z^2}{T_2^{\text{charge}}} \frac{t^2 + \varepsilon^2}{t^2 \varepsilon^2}$$

$$\Gamma_{\text{blockade}} \approx \frac{2\Delta E_z^2}{T_1^{\text{charge}}} \frac{t^2 + \varepsilon^2}{t^4}$$

- improve parity readout fidelity by reducing t





Conclusion

- origin of the parity readout
- blockade lifting rate as a function of T_1^{charge} and T_2^{charge}
- found indication that charge dephasing causes blockade lifting