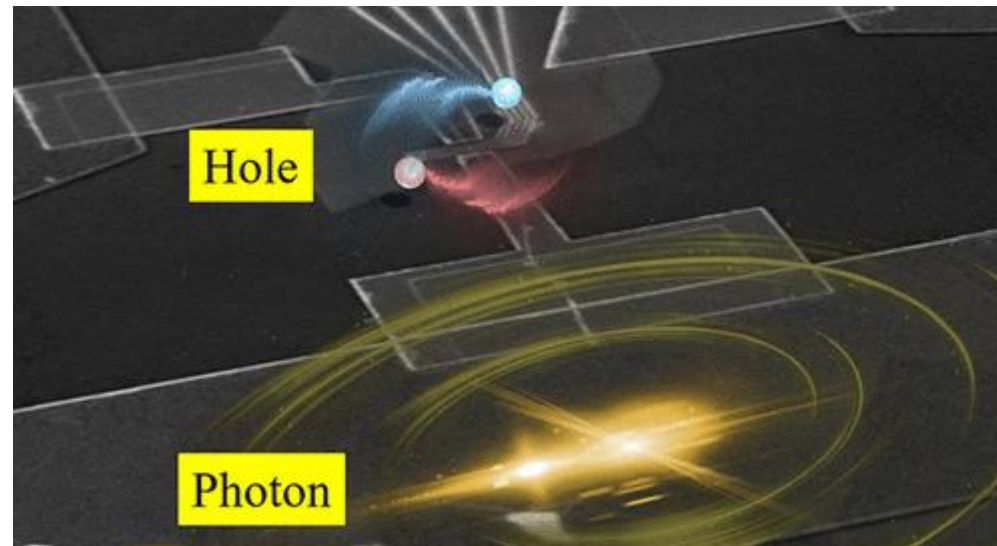


Gate Tunable Hole Charge Qubit Formed in a Ge/Si Nanowire Double Quantum Dot Coupled to Microwave Photons

Rui Wang, Russell S. Deacon, Jian Sun, Jun Yao, Charles M. Lieber, and Koji Ishibashi
Nano Letters **2019** 19 (2), 1052-1060 DOI: 10.1021/acs.nanolett.8b04343



-FAM talk-

Fabian Müller
22.03.2019

- Hole qubit:
 - P-wave nature leads to **enhanced qubit-cavity coupling g_c**
 - Strong spin orbit interaction (SOI) → **Fast manipulation** of hole spin qubit
- Ge/Si NW:
 - Lack of nuclear spin scattering → no hyperfine → **lower decoherence rate γ**

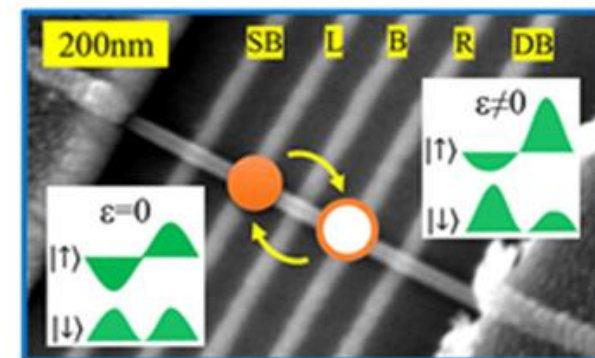
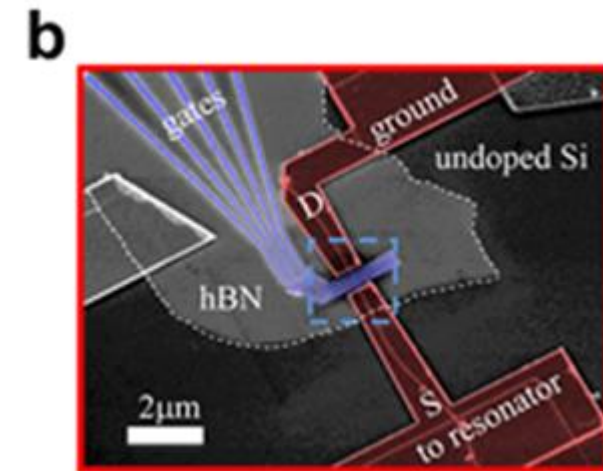
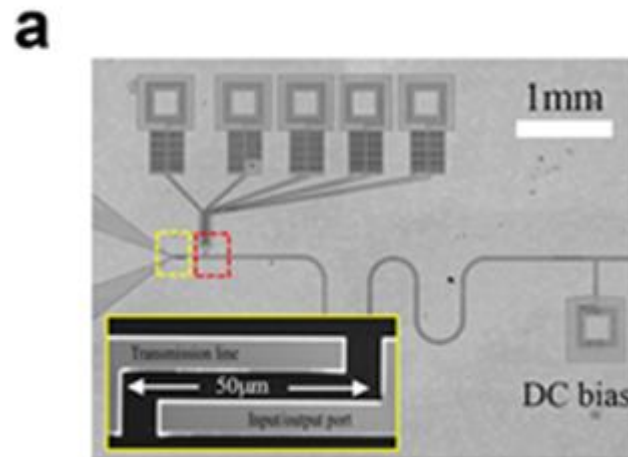
→ long-distance interaction between qubits & fast spin manipulation



1. Device
2. Two-level system (TLS) coupled to cavity
3. Resonator transmission spectroscopy
4. Power dependence of transmission spectroscopy
5. Summary & outlook

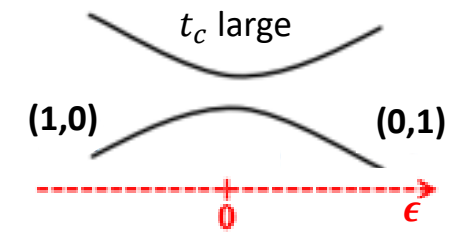
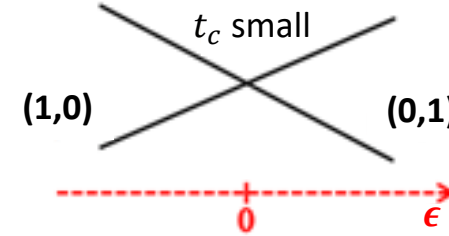
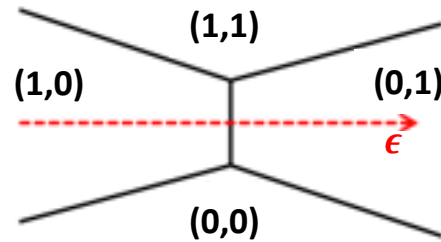


- MoRe superconducting transmission line resonator ($Q \approx 6000$), $f_c = 5.9667\text{GHz}$
- Ge/Si core/shell NW
- hBN as dielectric
- Tens of holes in QD, few holes regime unstable



- Qubit energy (TLS):

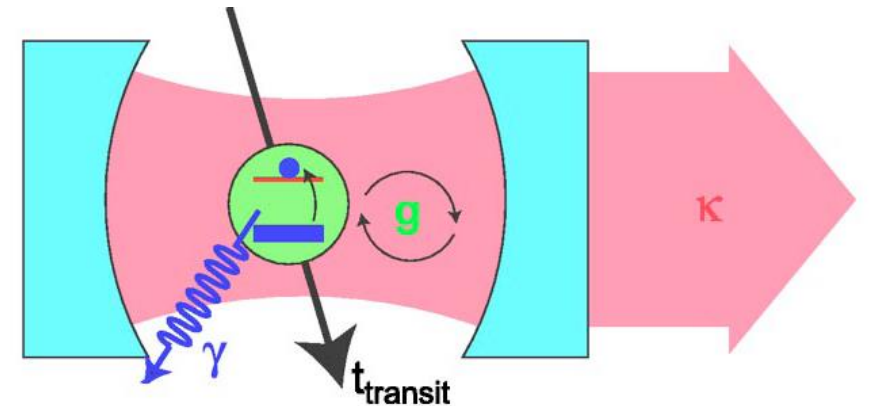
$$E_{qb} = \hbar\omega_{qb} = \sqrt{\epsilon^2 + (2t_c)^2}$$



- Jaynes-Cumming:

$$H_{JC} = \hbar\omega_c \left(a^\dagger a + \frac{1}{2} \right) + \frac{1}{2} \hbar\omega_{qb} \sigma_z + \hbar g_{eff} (a \sigma^+ + a^\dagger \sigma^-)$$

- Dressed States: $\frac{|\downarrow,1\rangle + |\uparrow,0\rangle}{\sqrt{2}}$ and $\frac{|\downarrow,1\rangle - |\uparrow,0\rangle}{\sqrt{2}}$



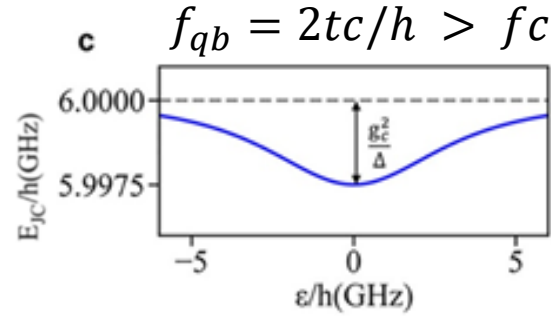
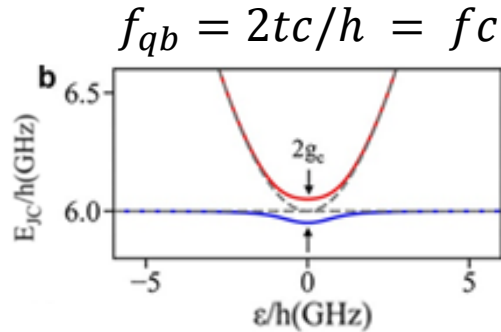
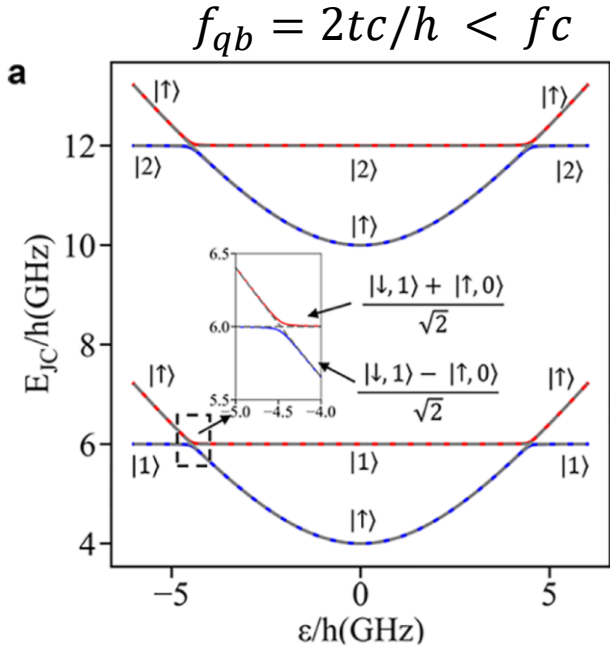
[1] A. Blais *et al.*, Phys. Rev. A **69**, 062320

Simulation of cQED System

- Qubit-cavity coupling: $g_{eff} = g_c \times 2t_c/E_{qb}$ with $E_{qb} = \hbar\omega_{qb} = \sqrt{\epsilon^2 + (2t_c)^2}$

Resonant regime: $f_{qb} \leq f_{resonator}$

Dispersive regime: $f_{qb} > f_{resonator}$



Cavity mode shift of $\Delta\theta = \frac{g_{eff}^2}{\Delta}$

$$\Delta = h(f_{qb} - f_c)$$

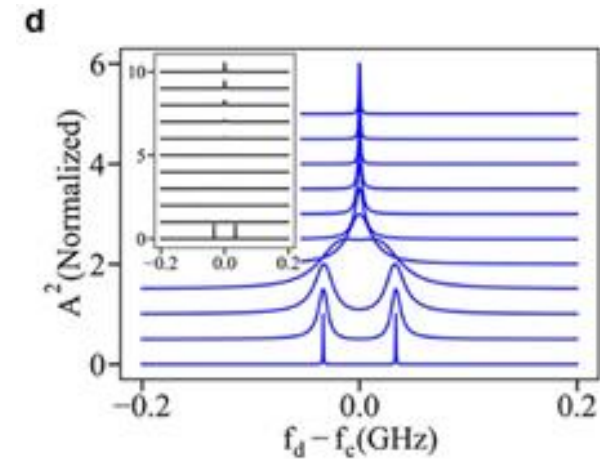
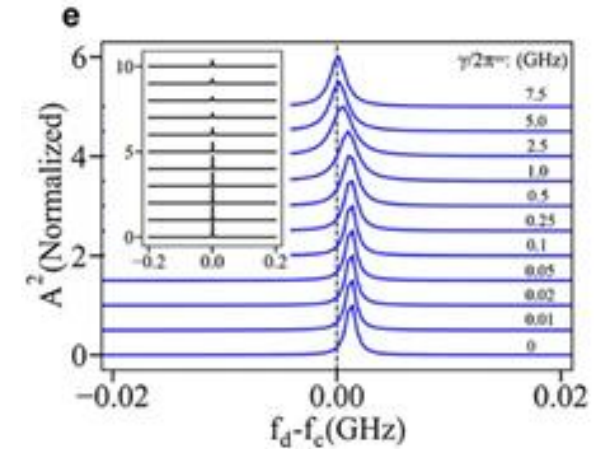


Simulation of Transmission Spectrum

- $g_c = 50\text{MHz}$
- $\Delta > 0$ **dispersive**, $f_{qb} = \frac{2t_c}{h}$
- $\forall \gamma \rightarrow g_{eff} / \Delta$ Shift due to dispersive pull

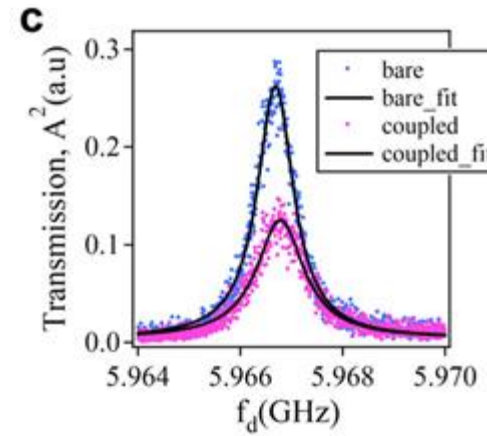
- $\Delta = h(f_{qb} - f_c) = 0$ **resonant**
- $\gamma < 2g_{eff}$ \rightarrow two separated peaks
- $\gamma \gg g_{eff}$ \rightarrow narrow peak

- Strong coupling regime: $g_c > \gamma, \kappa$

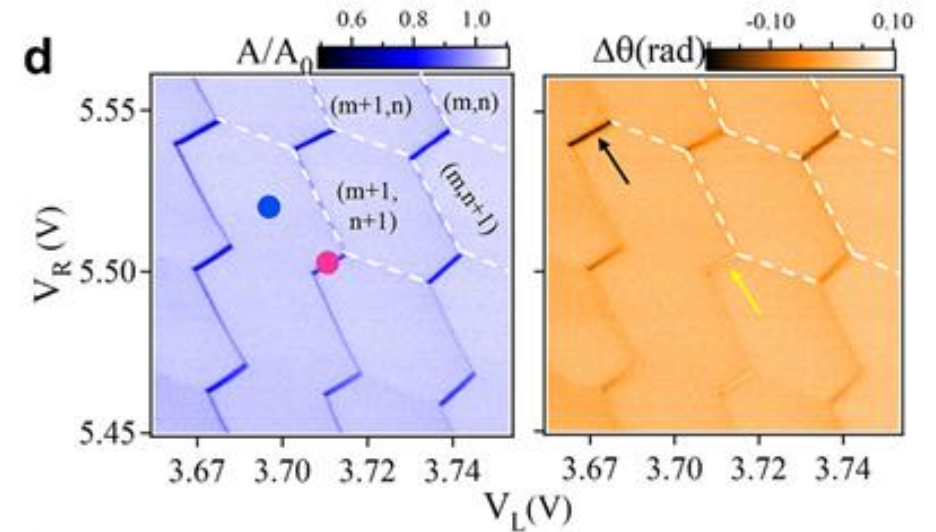


Transmission Spectroscopy

- $\epsilon \gg 0$: qubit is only capacitively coupled to the resonator
 - $\epsilon = 0$: if t_c is close to photon energy, there is qubit-resonator dipole coupling
- altered amplitude (dissipation) and shift (dispersive pull).



- Positive and negative phase shift implies that the qubit is close to resonance with the resonator frequency



Transmission Spectroscopy

- V_B in 10mV steps
 → change of tunneling rate $2t_c/h$
 → lowers E_{qb}
- Change from $f_{qb} > f_c$ ($+\Delta$) to $f_{qb} < f_c$ ($-\Delta$)
- $\Delta\theta = -\arctan\left(\frac{2g_c^2}{\kappa\Delta}\right)$

$$\gamma = 2\pi \times 5.25\text{GHz}$$

$$g_c = 2\pi \times 0.038\text{GHz}$$

$$\gamma = 2\pi \times 4.5\text{GHz}$$

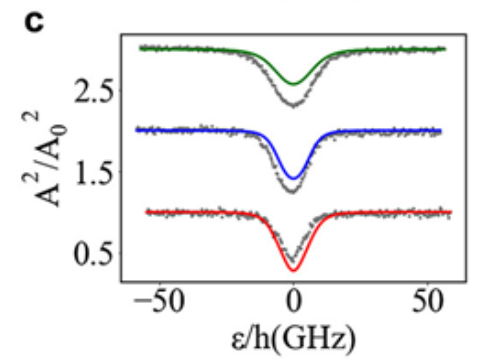
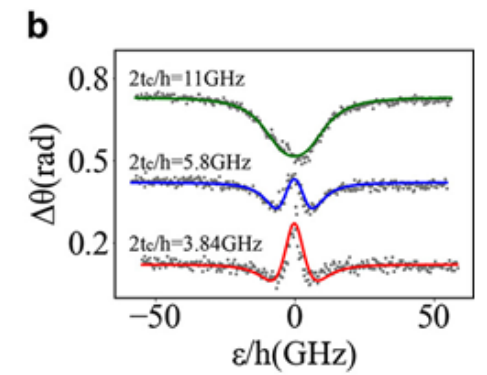
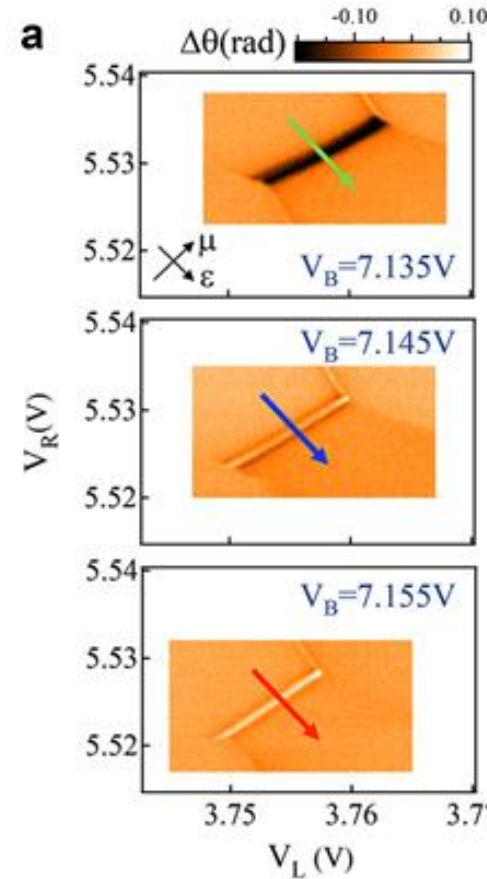
$$g_c = 2\pi \times 0.035\text{GHz}$$

$$\gamma = 2\pi \times 6.49\text{GHz}$$

$$g_c = 2\pi \times 0.055\text{GHz}$$

$$f_c = 5,9667\text{GHz}$$

$$\kappa = 6\text{MHz}$$



[1] K. Petersson *et al.*, *Nature* **490**, 380–383

Transmission Spectroscopy

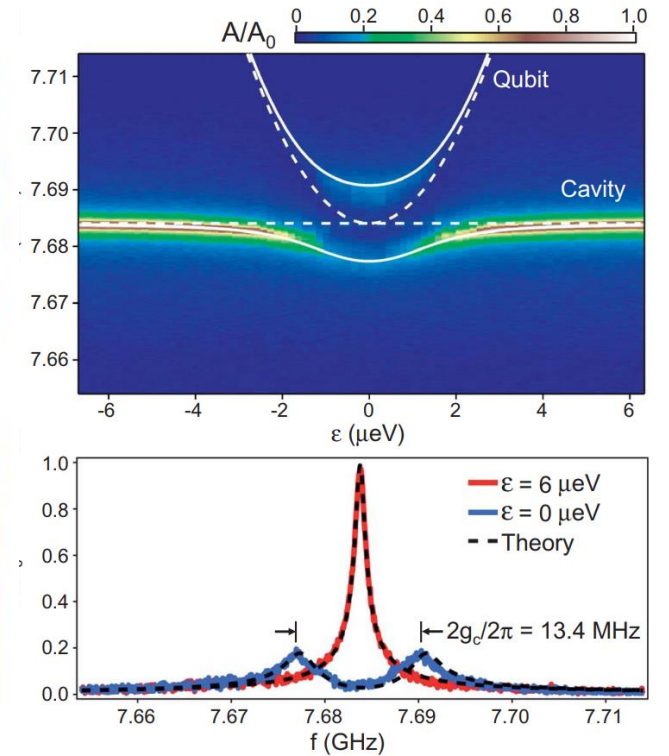
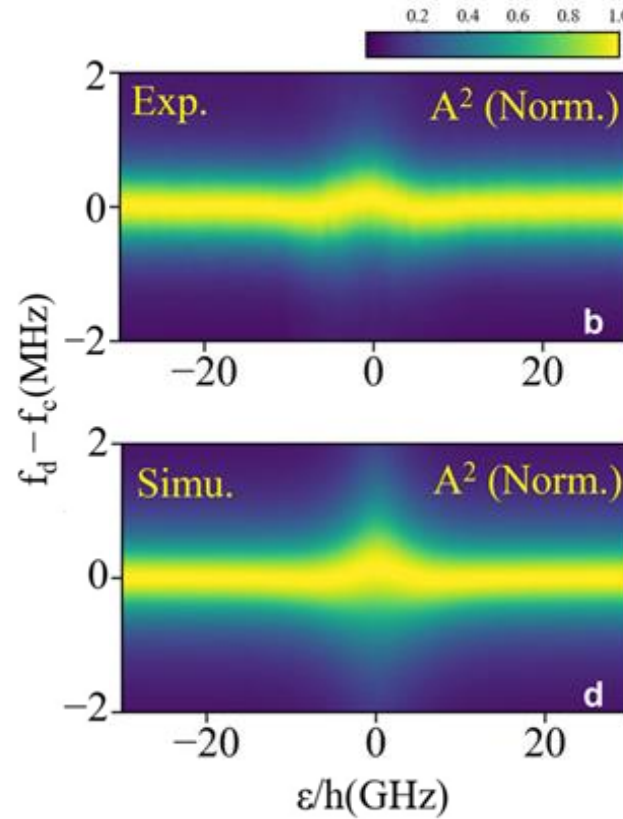
- No splitting (dressed state) resolved!

$$\rightarrow \gamma \gg g_c$$

$$(\gamma = 2\pi \times 4 - 6\text{GHz}, g_c = 90\text{MHz})$$

- Lower charge noise sensitivity of real device

$$2t_c/h = 3.84\text{GHz} < f_c$$



Compared to strong coupling regime [1]

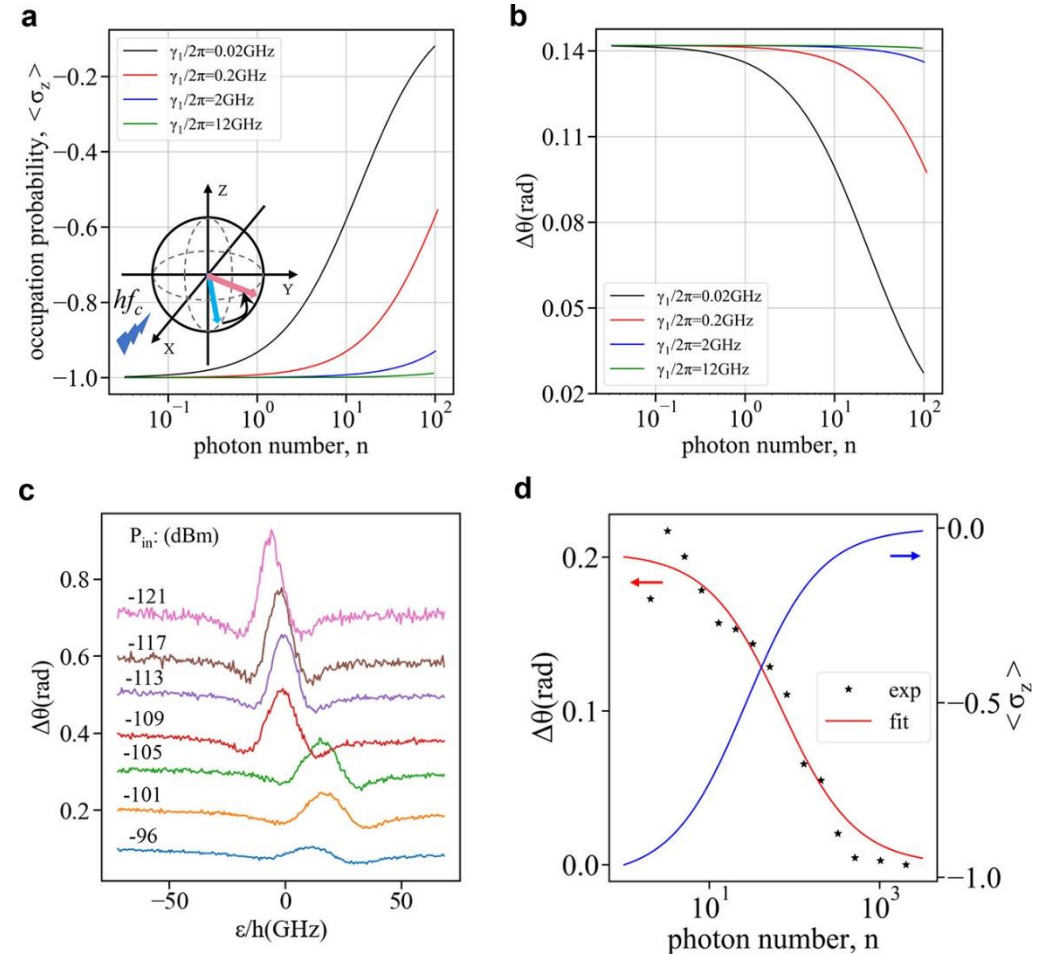
[1] X. Mi *et al.*, Science **355**, 156-158

Power Dependence of Transmission

Separate influence of qubit energy relaxation γ_1 and pure dephasing γ_ϕ

- Phase-shift goes to zero for large Power (n)
 $\gamma = 2\pi \times 5\text{GHz}, \gamma_1 = 2\pi \times 70\text{MHz}$
 $\rightarrow \gamma_\phi = 2\pi \times 5\text{GHz}$ (charge noise)
- Large drive power induces broadening, merging, splitting of Rabi peaks when qubit and resonator resonantly coupled

\rightarrow Pure dephasing dominates the decoherence



Conclusion and Outlook

- Expect stronger coupling for hole qubit than for electron qubit, but they could not reach strong coupling regime
 - Power dependence of the cavity mode dispersive shift revealed pure dephasing dominates decoherence of the qubit
- Improve qubit coherence times for strong coupling by **increase cleanness of hBN/NW interface**
- Zajac et al claim **accumulation type QD** will have reduced environmental noise
- **High-impedance resonator** to elevate coupling strength (several fold)

