

High-impedance surface acoustic wave resonators

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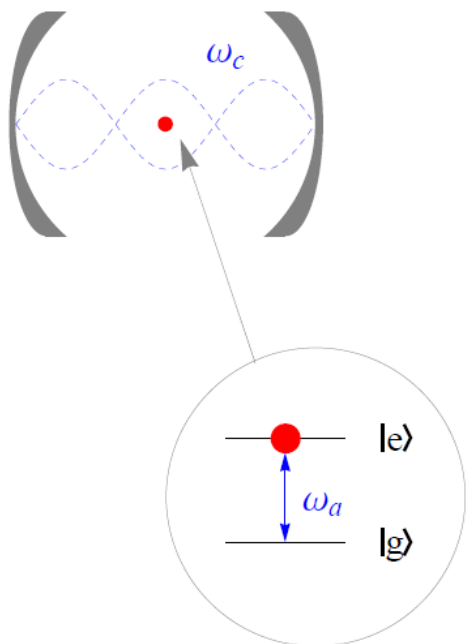
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Because of their small size, low loss, and compatibility with magnetic fields and elevated temperatures, surface acoustic wave resonators hold significant potential as future quantum interconnects. Here, we design, fabricate, and characterize GHz-frequency surface acoustic wave resonators with the potential for strong capacitive coupling to nanoscale solid-state quantum systems, **including semiconductor quantum dots.** Strong capacitive coupling to such systems requires a large characteristic impedance, and the **resonators we fabricate have impedance values above 100 Ω .** We achieve such high impedance values by tightly confining a Gaussian acoustic mode. At the same time, the resonators also have low loss, with **quality factors of several thousand at millikelvin temperatures.** These high-impedance resonators are expected to exhibit large vacuum electric-field fluctuations and have the potential for strong coupling to a variety of solid-state quantum systems.

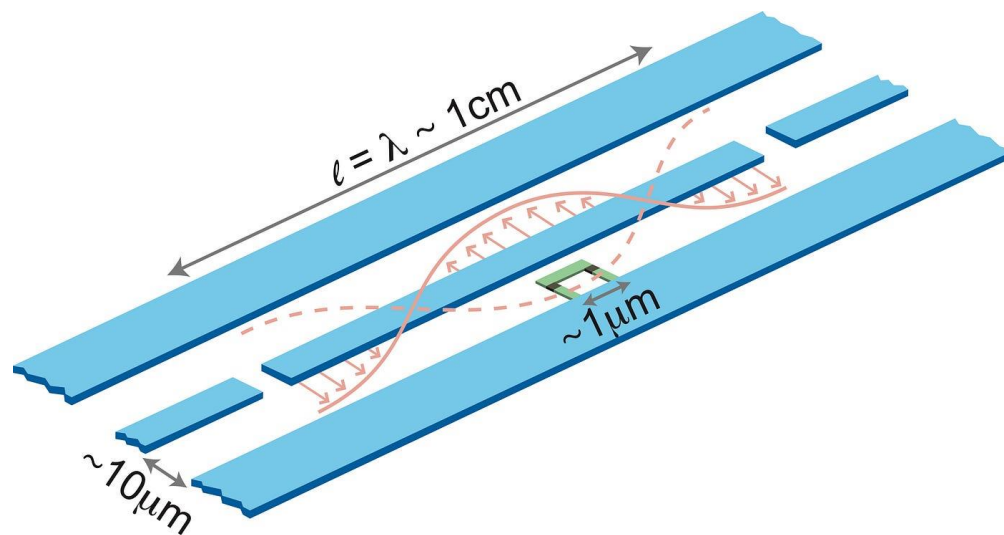
Cavity QED

$$H_{JC} = \hbar\omega_{cavity}a^+a + \hbar\omega_{atom}\frac{\sigma_z}{2} + \hbar g(a^+\sigma_- + a\sigma_+)$$



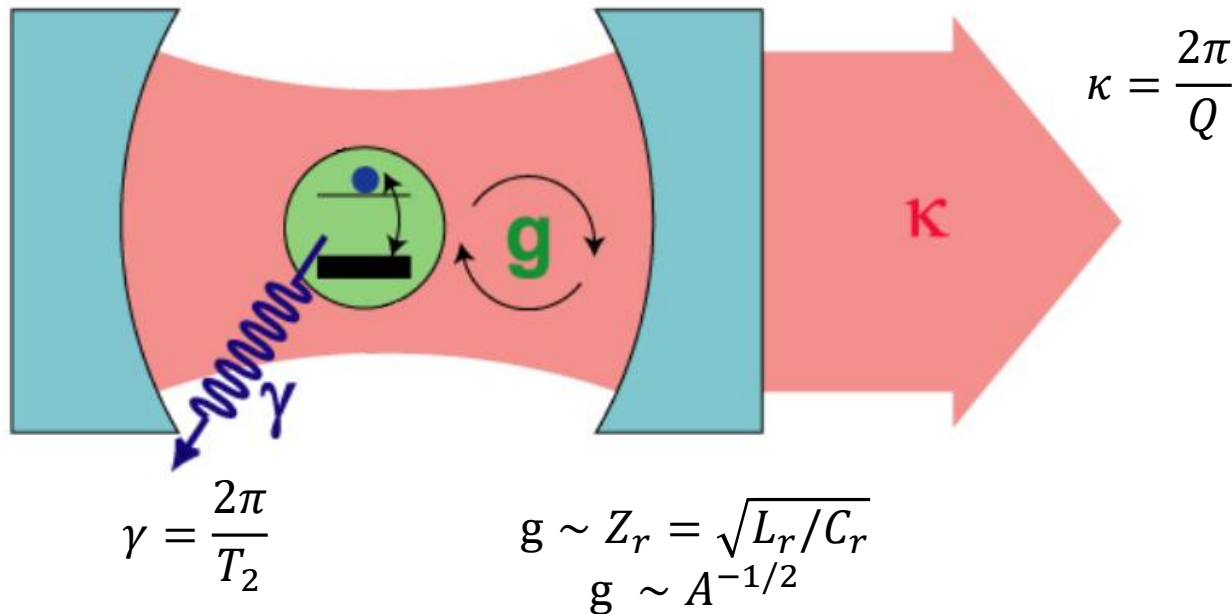
circuit QED

$$H_{JC} = \hbar\omega_{resonator}a^+a + \hbar\omega_{qubit}\frac{\sigma_z}{2} + \hbar g(a^+\sigma_- + a\sigma_+)$$



A. Blais, R.-S. Huang, A. Wallraff, S. M. Girvin, and R. J. Schoelkopf, Physical Review A **69**, 062320 (2004).

What are the relevant parameters?



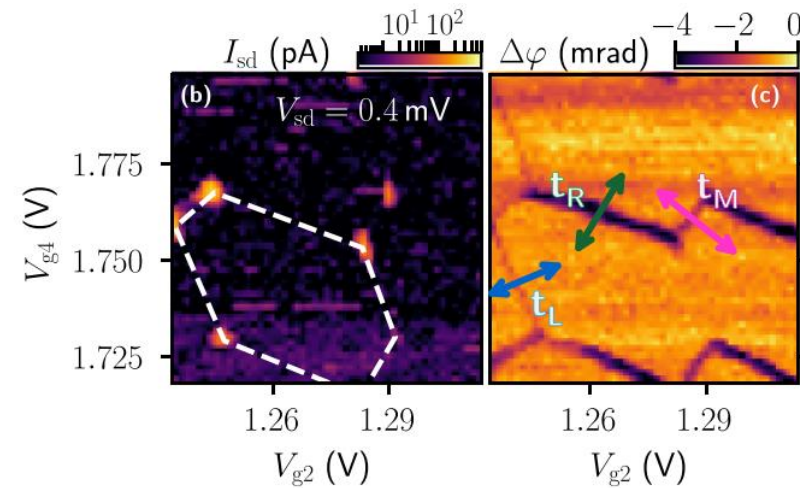
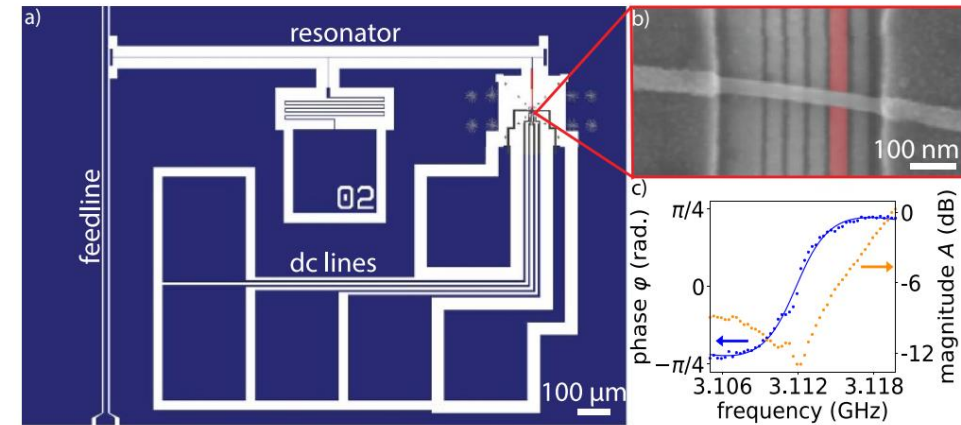
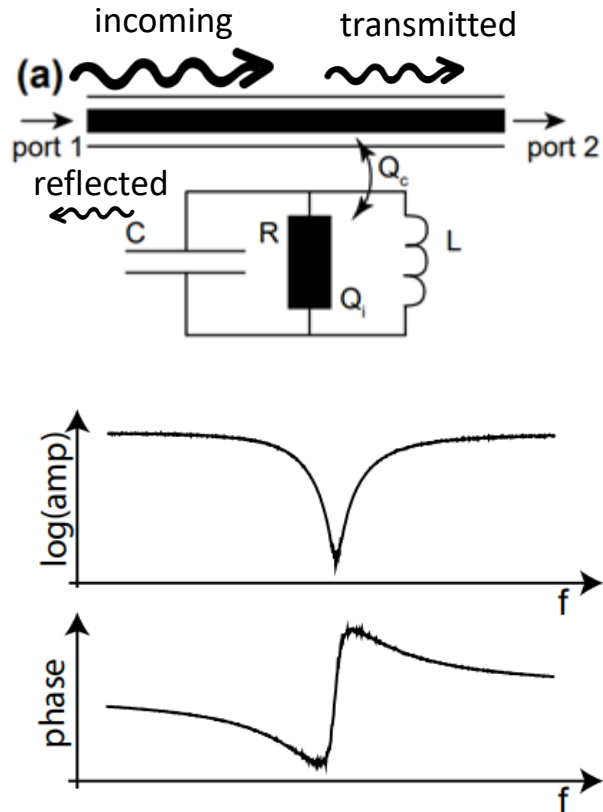
γ is determined by the coherence of your 2 level system

κ is determined by the confinement of your cavity

g is determined by the nature of the interaction and the resonator circuit parameters

When $g \gg \kappa, \gamma$ then you are in the strong coupling regime and the JC Hamiltonian describes the model

How do we measure a resonator?

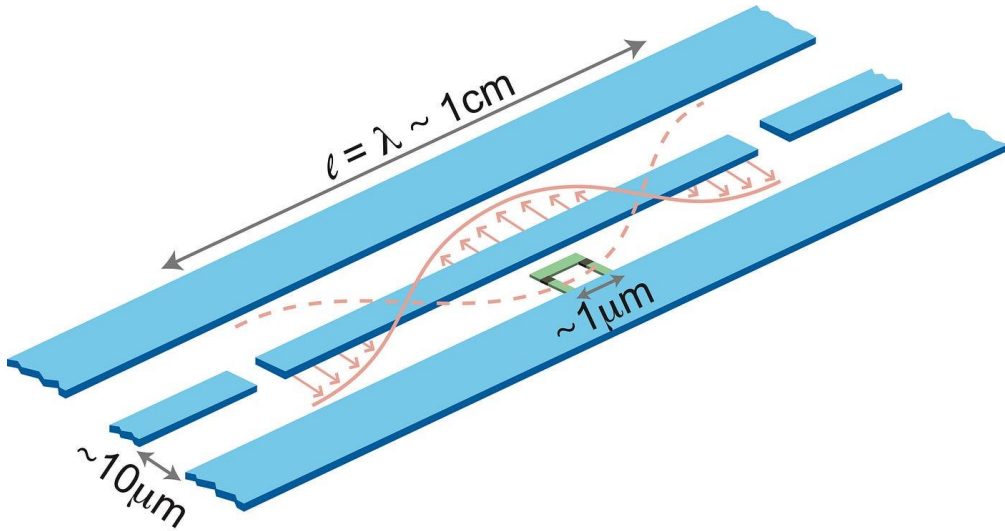


Efficient and robust analysis of complex scattering data under noise in microwave resonators
 S. Probst et al Rev. Sci. Instrum. 86, 024706 (2015)

Charge-sensing of a Ge/Si core/shell nanowire double quantum dot using a high-impedance superconducting resonator
 J. H. Ungerer, P. Chevalier Kwon et al Mater. Quantum. Technol. 3 031001 (2023)

Circuit QED

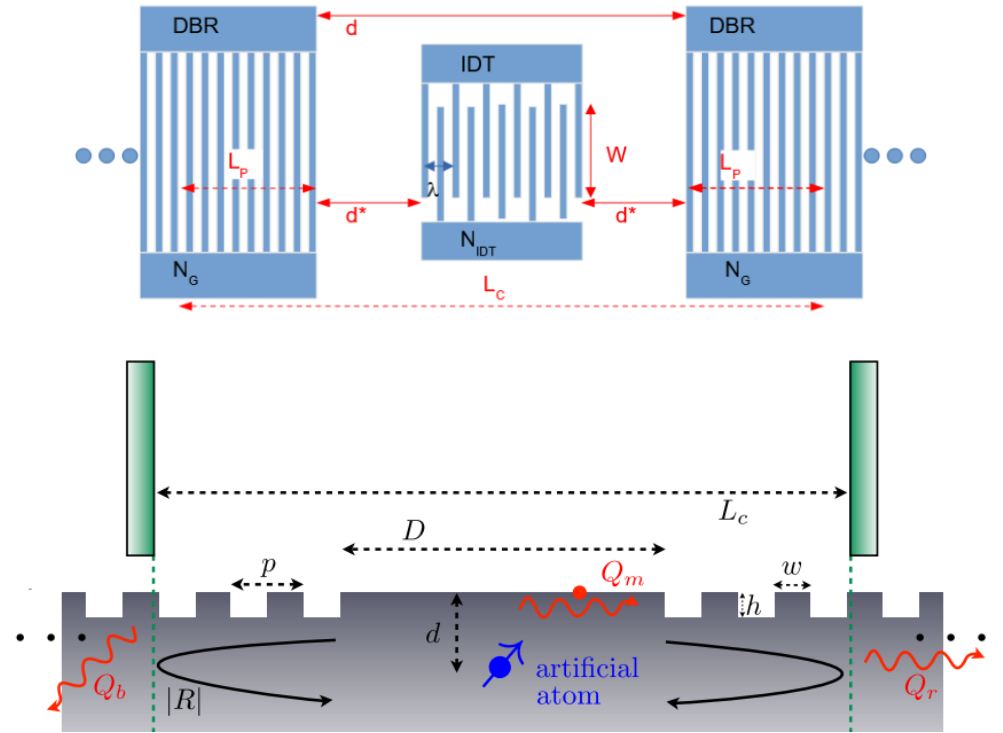
$$H_{JC} = \hbar\omega_r a^\dagger a + \hbar\omega_q \frac{\sigma_z}{2} + \hbar g(a^\dagger \sigma_- + a \sigma_+)$$



Cavity quantum electrodynamics for superconducting electrical circuits: An architecture for quantum computation
 A. Blais et al Physical Review A **69**, 062320 (2004).

Circuit QED with phonons (QAD?)

$$H_{JC} = \hbar\omega_r a^\dagger a + \hbar\omega_q \frac{\sigma_z}{2} + \hbar g(a^\dagger \sigma_- + a \sigma_+)$$



Universal Quantum Transducers Based on Surface Acoustic Waves
 M. J. A. Schuetz et al Phys. Rev. X **5**, 031031

The 1 port Saw resonator

$$Q_0 = \frac{\pi(d + 2L_p)}{\lambda_0(1 - \tanh(|r_S| N_G))}$$

$$Q_{\text{Diffraction}} = \frac{5\pi}{|1 + \gamma|} \left(\frac{W}{\lambda_0}\right)^2$$

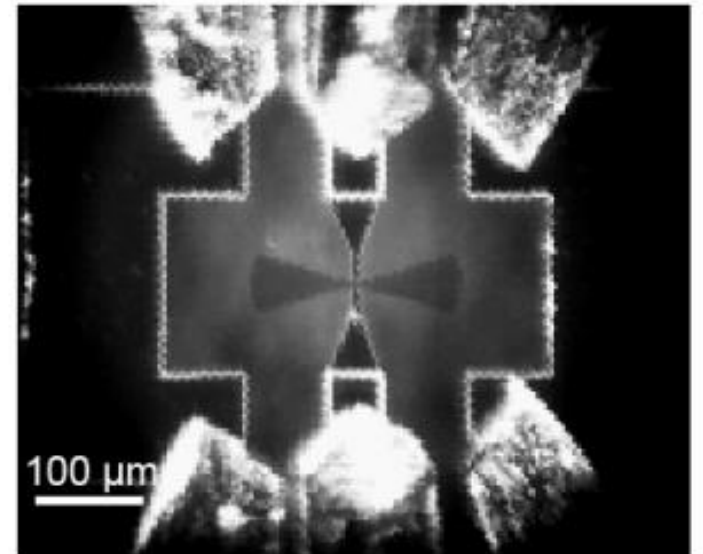
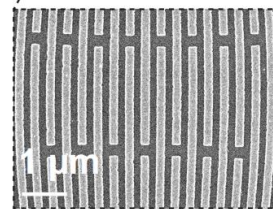
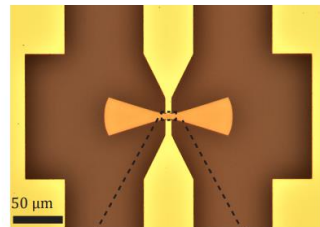
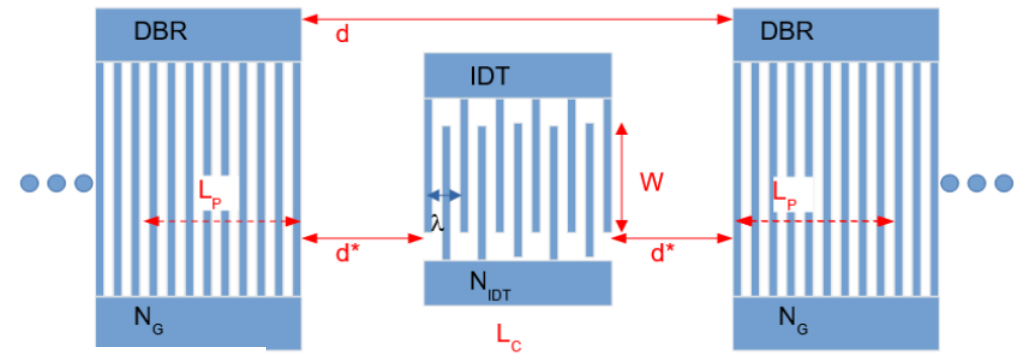
$$Q_{\text{ext}} = \frac{1}{5.74 v_{\text{SAW}} Z_C C_S W K^2 (0.5 N_{\text{IDT}})^2} \cdot L_c$$

$$Q_{\text{bk}} = 2\pi N_{\text{eff}} / [C_b (h/\lambda_c)^2]$$

	GaAs	LiNbO3
V_saw	2864 m/s	3488 m/s
C_s	1.2 pF/cm	4.6 pF/cm
K ²	0.07	4.8
gamma	-0.537	-1.08

Fabrication Parameters

Material Parameters



Modeling the resonator

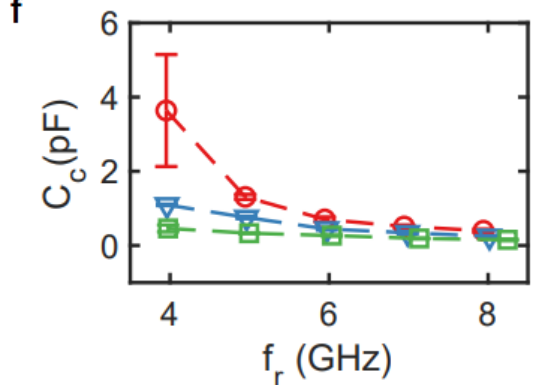
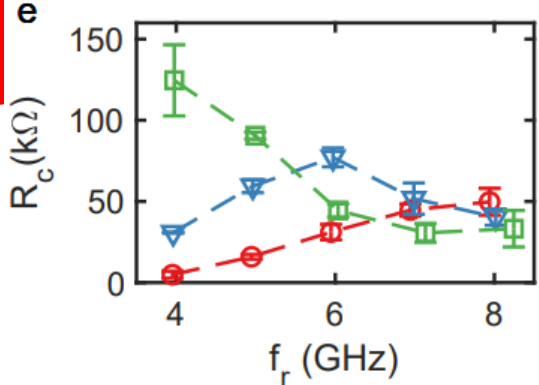
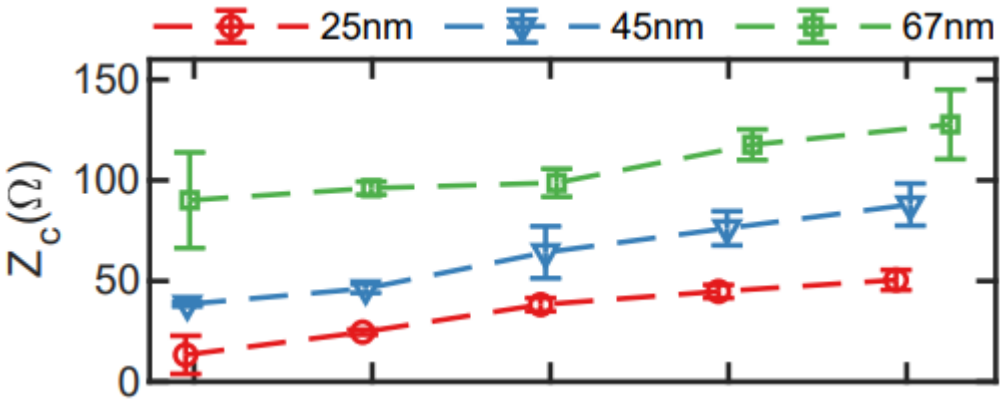
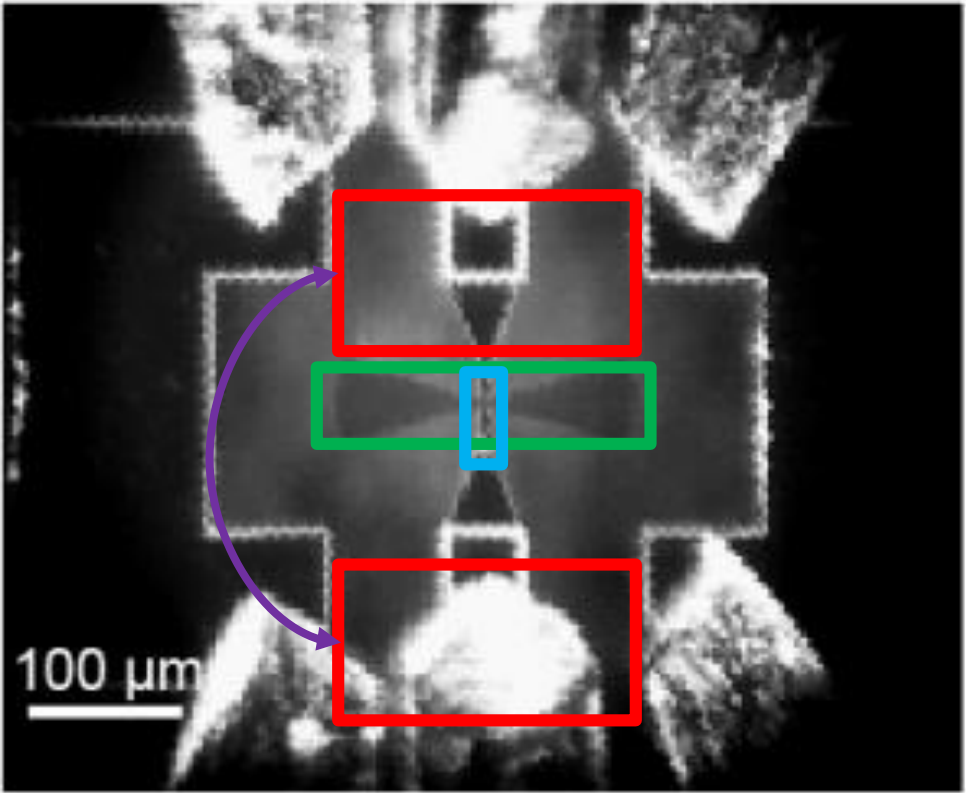
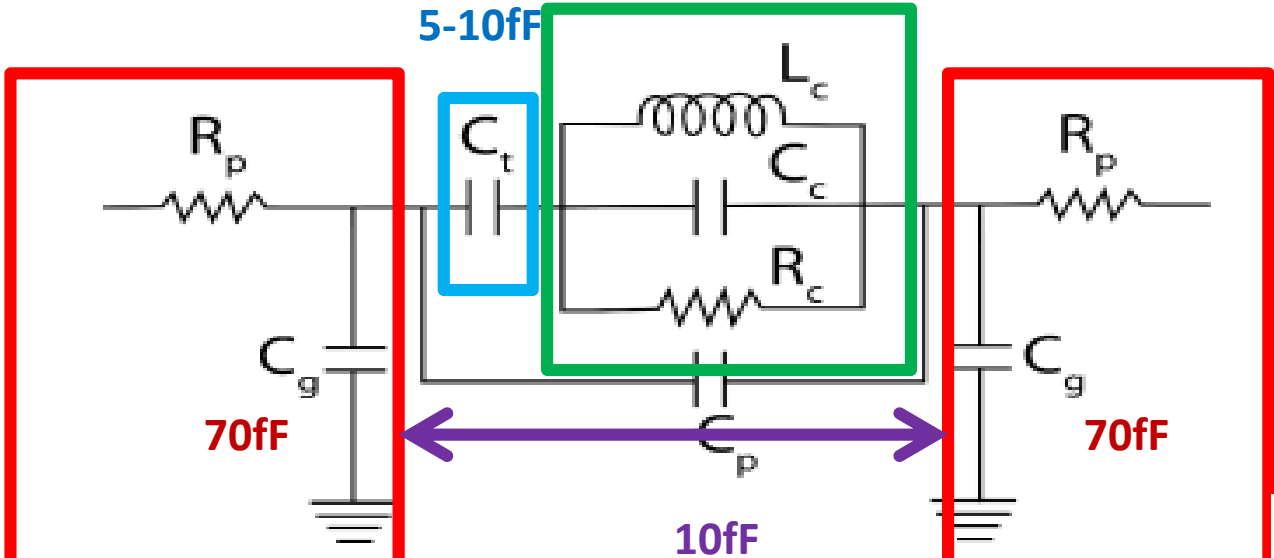


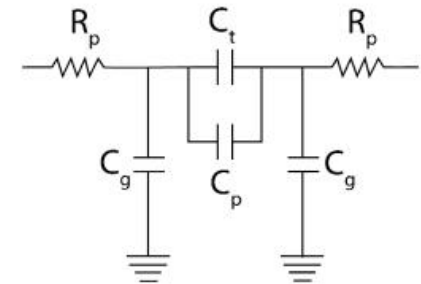
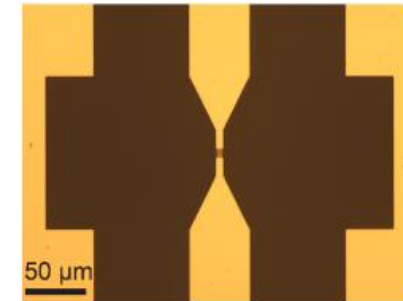
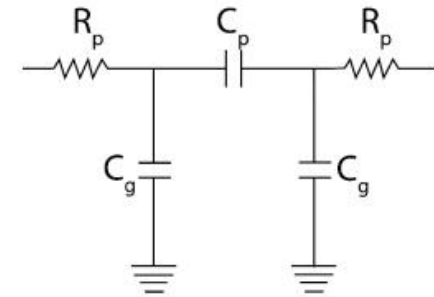
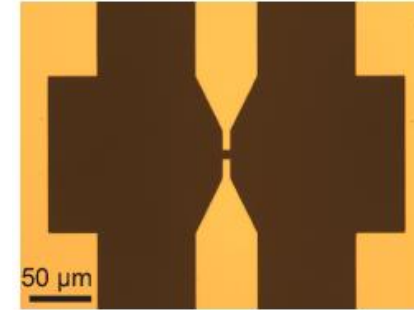
FIG. S4.

Fitting the parameters

Calibrate the VNA

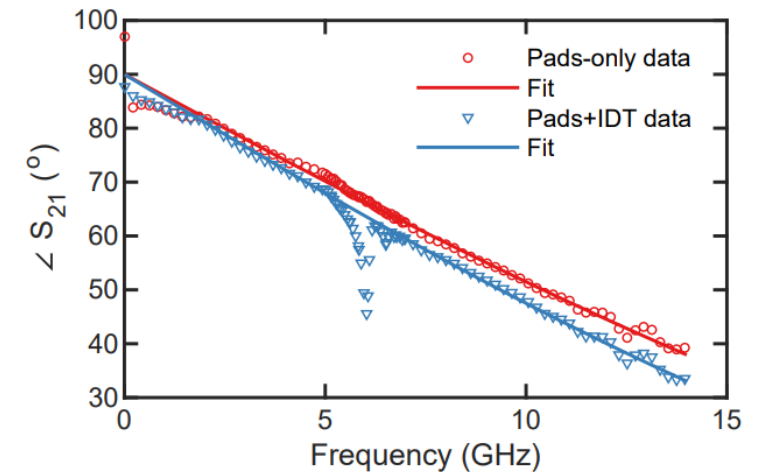
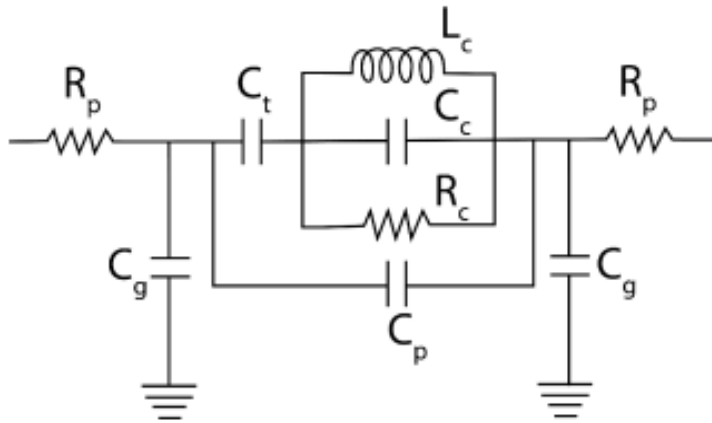
Use the sim-Cg and fit R_p and C_p from a pad only device

Simulate C_p C_g with a pad only device

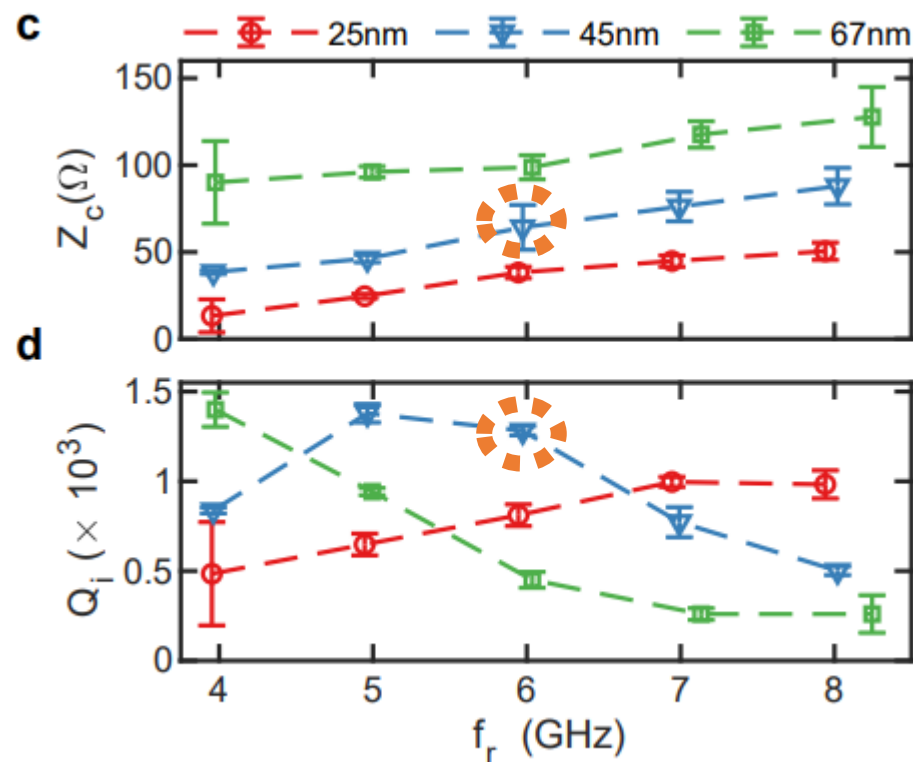
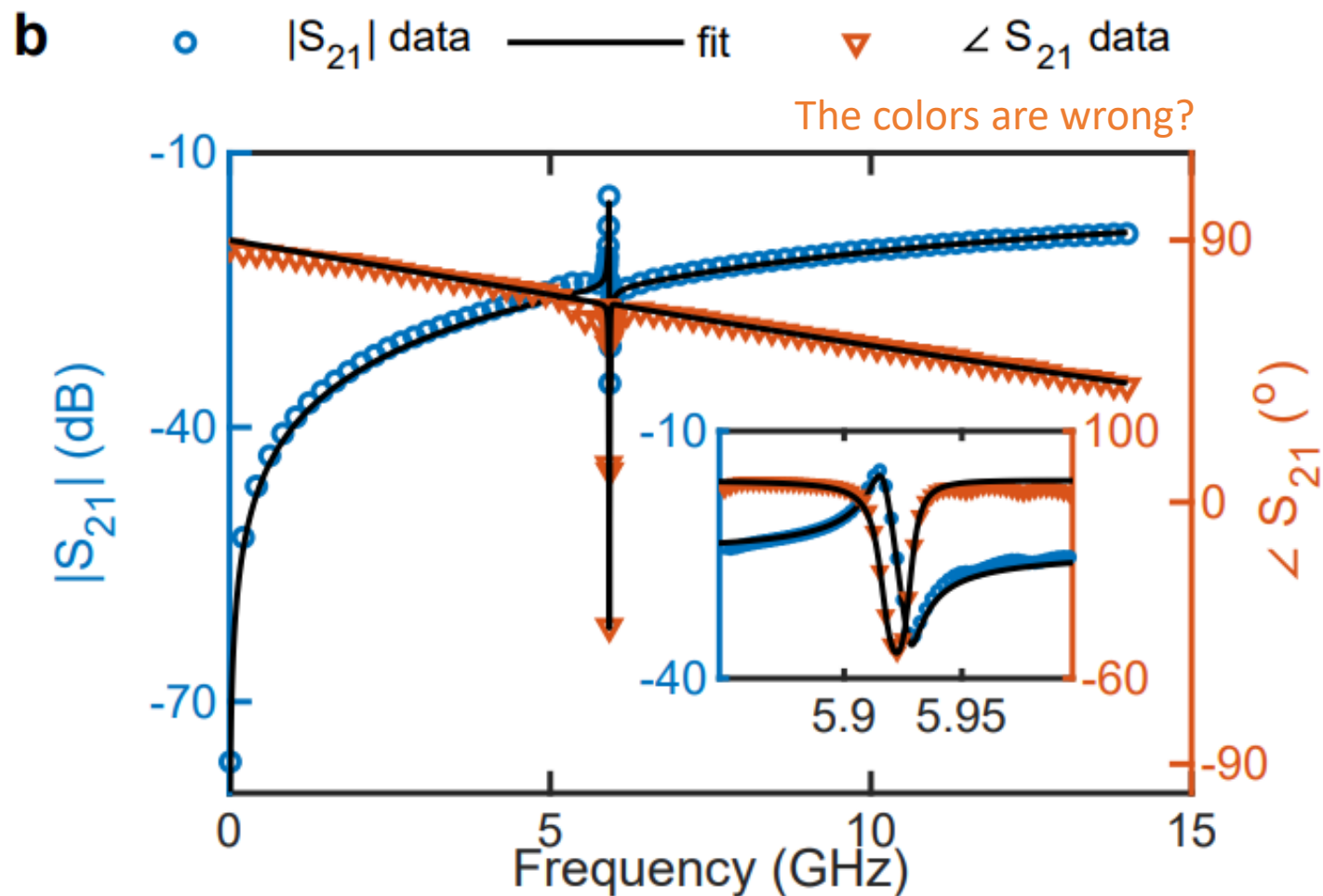


fit C_t and update R_p in a IDT device, ignoring the resonance

Fit C_c R_c fitting near the resonance



Characterization at RT



QRAP resonators

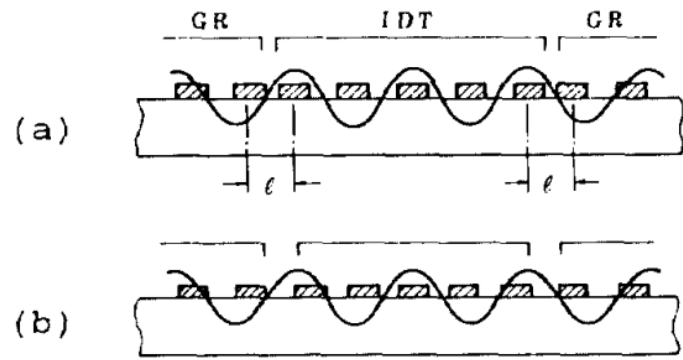


Figure-1 Schematic of SAW Resonators
 (a) Conventional (b) QRAP Structure

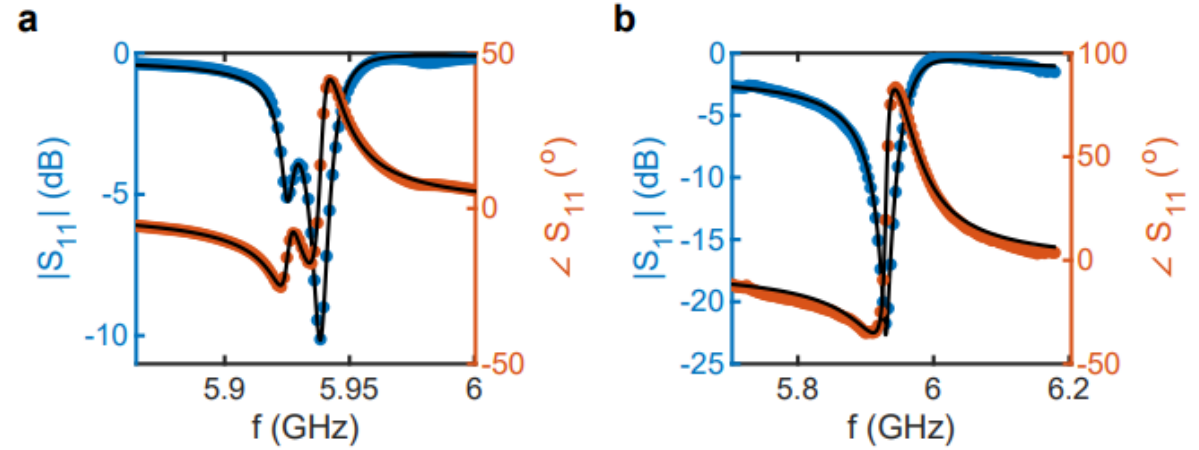


Fig 3 a

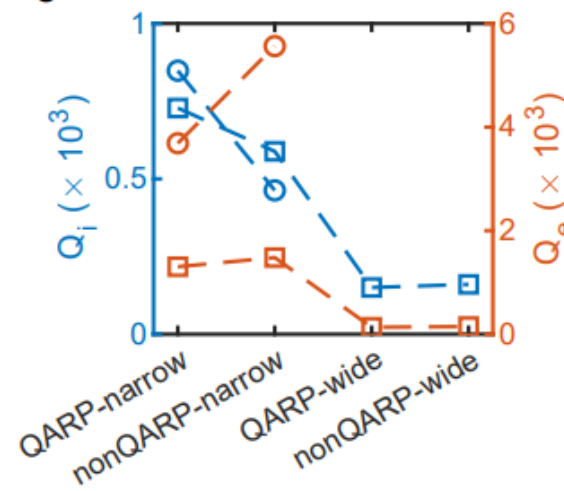
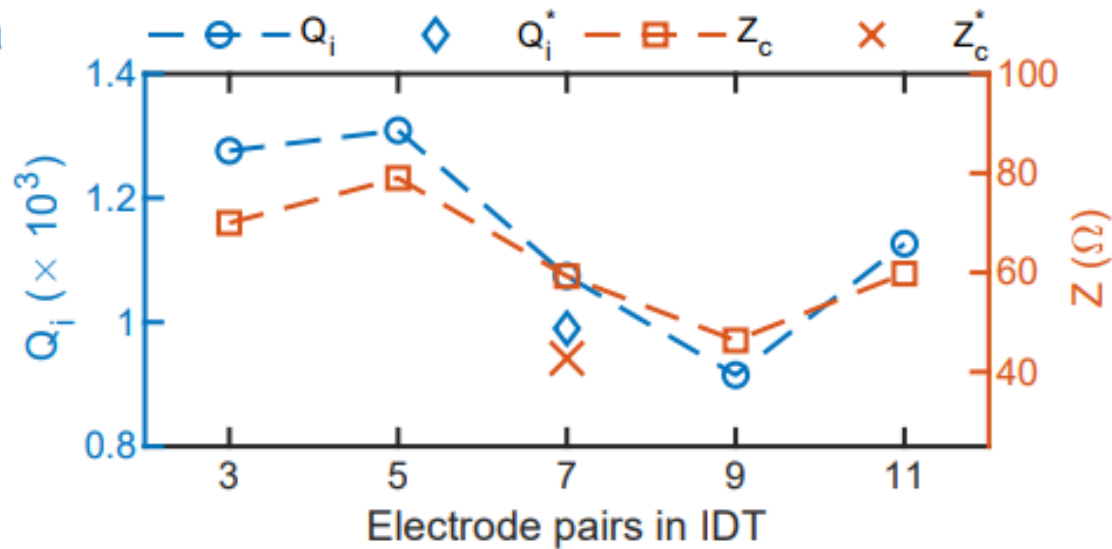
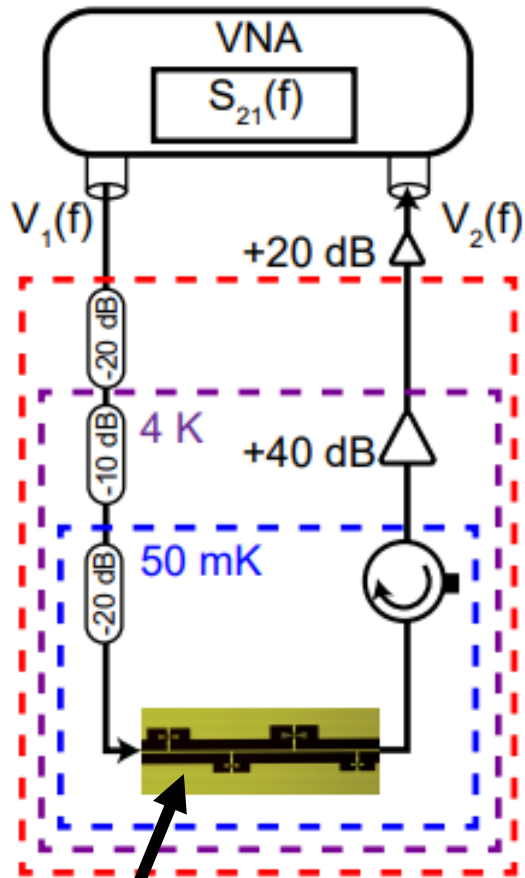
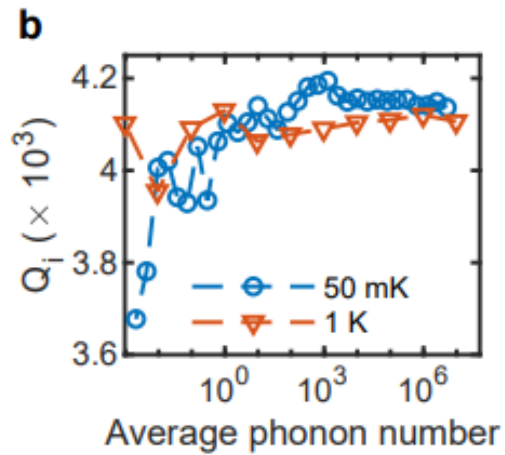
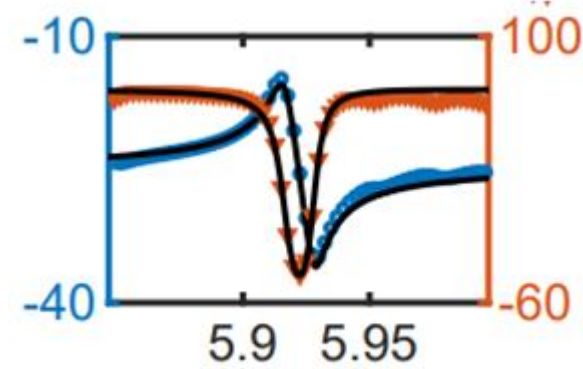
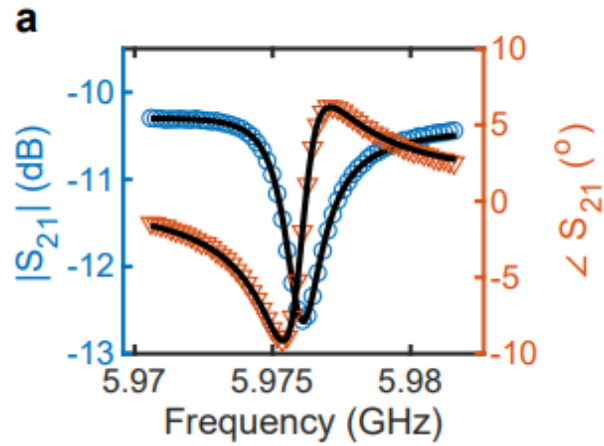


Fig S8

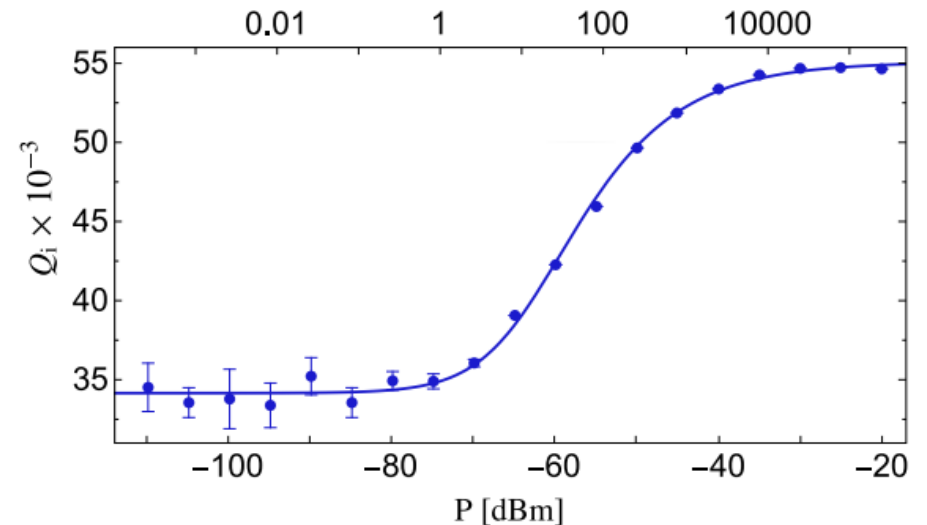
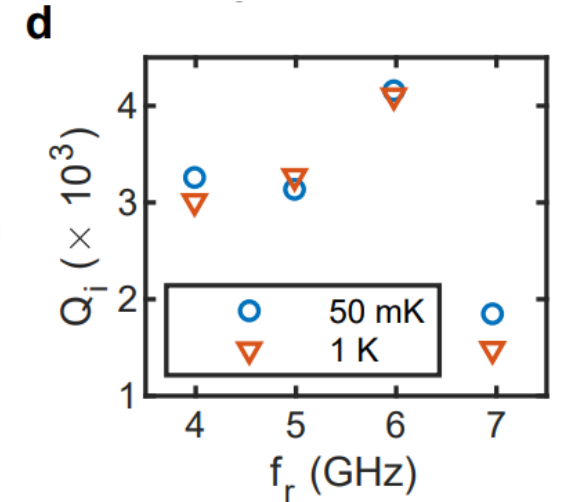
Cryogenic characterization



Nb CPW, T_c 9.3K



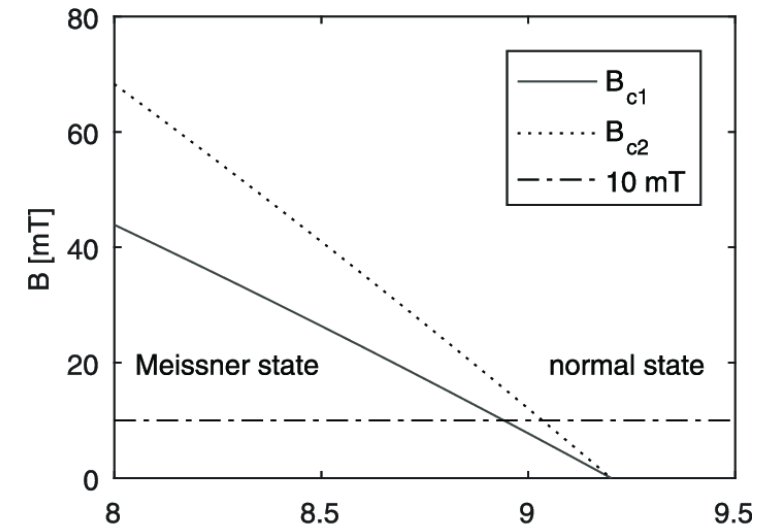
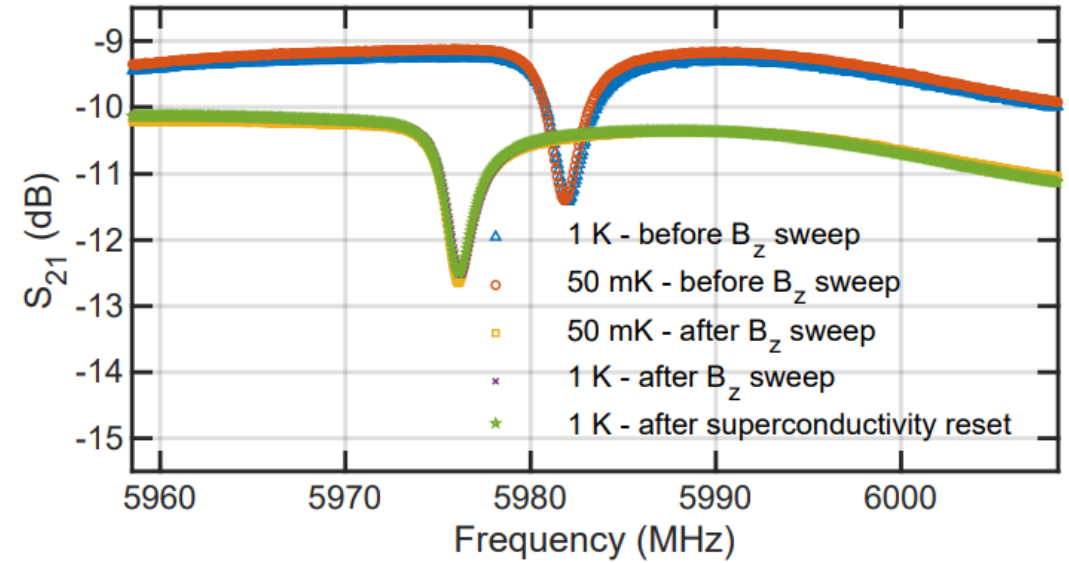
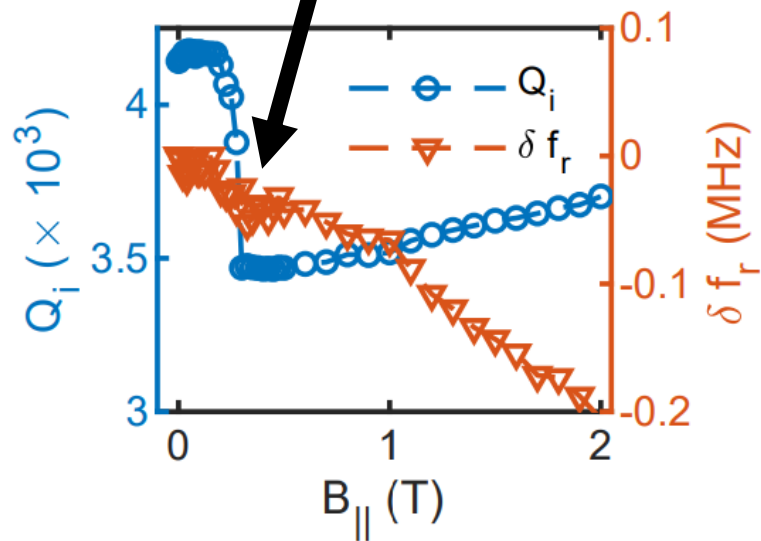
$$\langle n \rangle = \frac{Q_c}{\omega_r} \left(\frac{Q_i}{Q_i + Q_c} \right)^2 \frac{P_{in}}{\hbar \omega_r}$$



Surface acoustic wave resonators in the quantum regime PRB 93, 041411(R) R. Manenti et al

Magnetic field dependence

Loss of SC of the electrodes



In conclusion



- It's possible to have high Z at RT
- Loss of SC is not changing Q in a large way
- Do they really show high Z at low temperature?
- The substrate used does not allow for semiconductor dots

Because of their small size, low loss, and compatibility with magnetic fields and elevated temperatures, surface acoustic wave resonators hold significant potential as future quantum interconnects. Here, we design, fabricate, and characterize GHz-frequency surface acoustic wave resonators with the potential for strong capacitive coupling to nanoscale solid-state quantum systems, **including semiconductor quantum dots**. Strong capacitive coupling to such systems requires a large characteristic impedance, and the **resonators we fabricate have impedance values above 100 Ω** . We achieve such high impedance values by tightly confining a Gaussian acoustic mode. At the same time, the resonators also have low loss, with **quality factors of several thousand at millikelvin temperatures**. These high-impedance resonators are expected to exhibit large vacuum electric-field fluctuations and have the potential for strong coupling to a variety of solid-state quantum systems.