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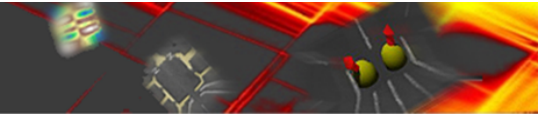
Coherent transfer of quantum information in a silicon double quantum dot using resonant SWAP gates

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2019

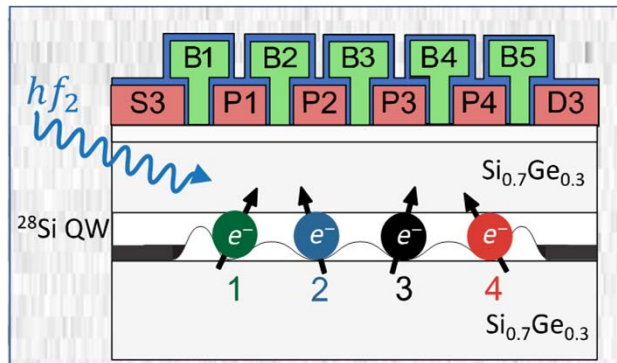
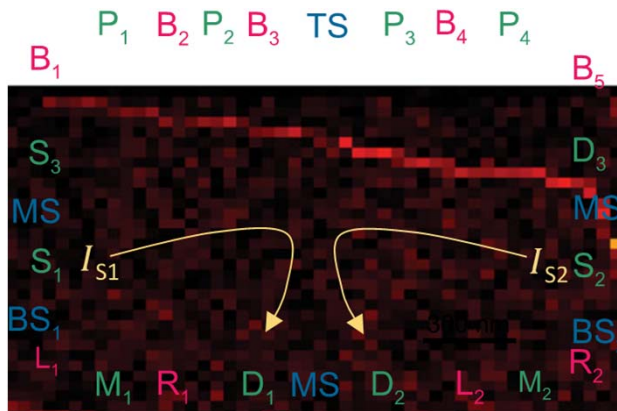
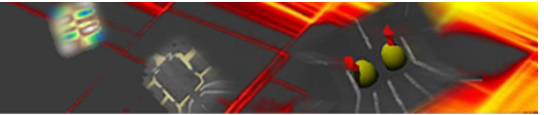
Taras Patlatiuk

20.03. 2020



- electron spin qubits in Si
- high-fidelity single- and two-qubit (exchange) gates
 - error correction
 - random access memory
 - multiqubit algorithms
- quantum SWAP gate → phase coherent SWAP
 - move spin eigenstates in 100 ns, $\bar{F}_{SWAP}^{(p)} = 98\%$
 - transfer product states in 300 ns, $\bar{F}_{SWAP}^{(p)} = 84\%$
- coupling of non-adjacent qubits

Device architecture



Phys. Rev. Appl. 11, 061006 (2019)

- isotopically enriched $^{28}\text{Si}/\text{SiGe}$
- micromagnet
- electron dipole spin resonance (single spin control)
- qubits Q_3 and Q_4 (under plunger gates P_3 and P_4)
- $(N_3, N_4) = (1, 1)$
- charge sensing with I_{S2} (readout 3 ms)
- measure/initialize Q_4 (via spin-selective tunneling to reservoir under D_3)
- loading fidelity 95% limited by 110 mK

Two-qubit interactions

modulate exchange interaction $J_{3,4}(V_{B4})$

$$V_{B4}(t) = V_{B4}^{(dc)} + V_{B4}^{(ac)} \cos(2\pi f_{\text{SWAP}} + \phi)$$

$J_{3,4} \gg \gamma_e |B_3^{\text{tot}} - B_4^{\text{tot}}|$ - magnetic field gradient
 γ_e - gyromagnetic ratio

No

Yes

CPHASE-like evolution

SWAP oscillations

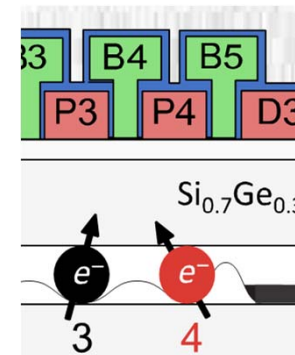
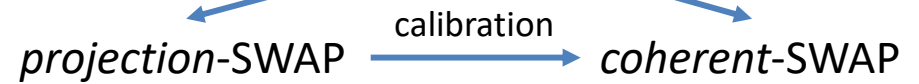
$$2\pi f_{\text{SWAP}} = \gamma_e |B_3^{\text{tot}} - B_4^{\text{tot}}|$$

$|\uparrow\uparrow\rangle, |\downarrow\downarrow\rangle$ - unaffected

$|\phi_3, \phi_4\rangle \in \{|\uparrow\downarrow\rangle, |\downarrow\uparrow\rangle\}$

π -pulse = SWAP + single-qubit phases

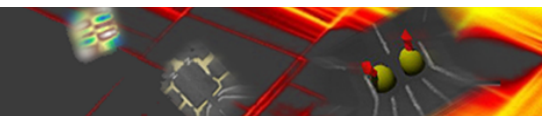
resonant SWAP gate



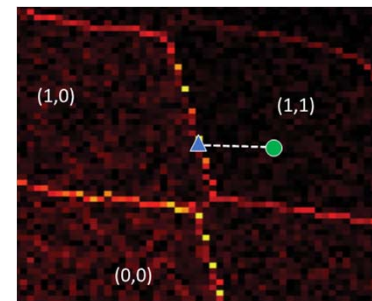
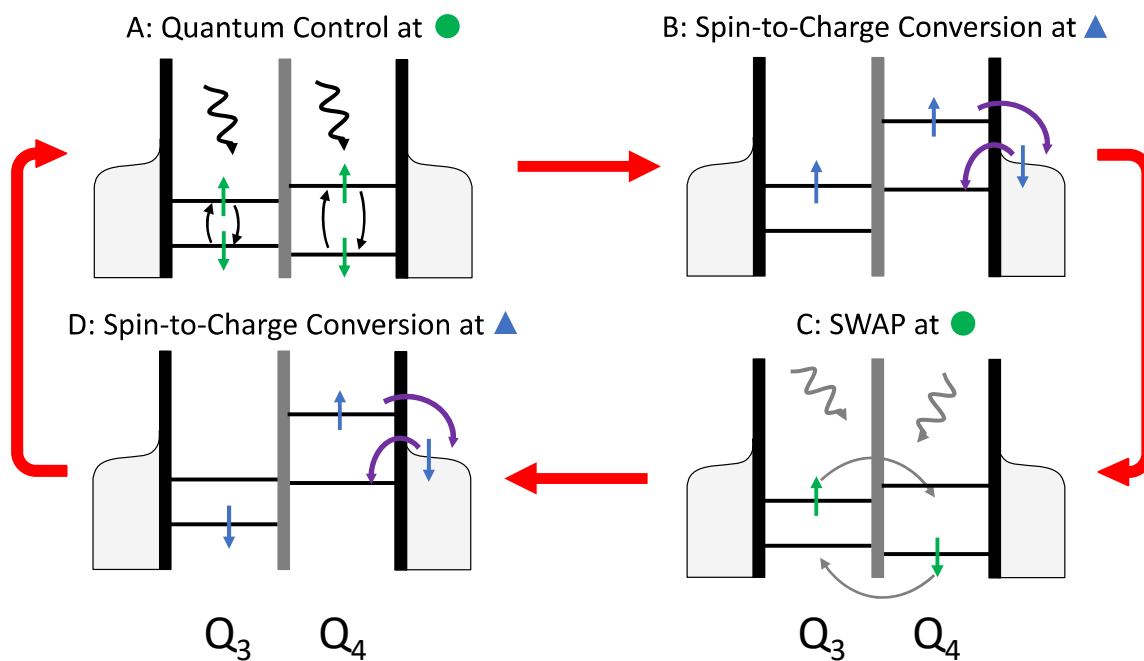
$$\gamma_e B_3^{\text{tot}} = 16.949 \text{ GHz}$$

$$\gamma_e B_4^{\text{tot}} = 17.089 \text{ GHz}$$

$$B_{\text{ext}} = 410 \text{ mT}$$



Measurement Cycle

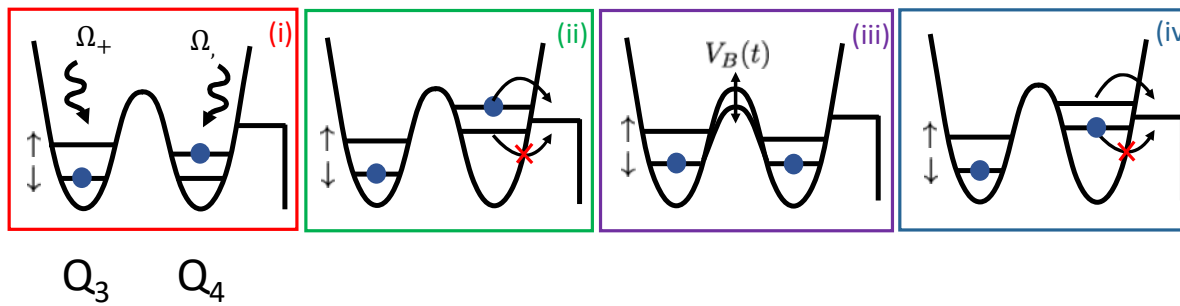
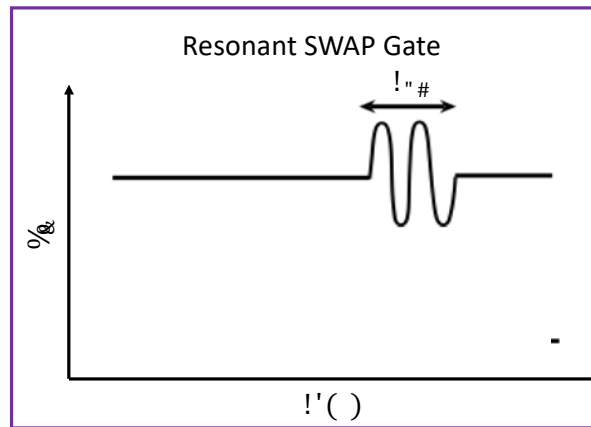
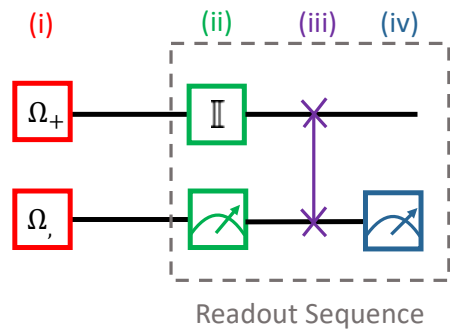
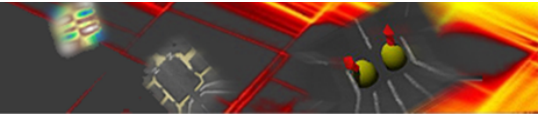


A: individual manipulation

B: spin-selective tunneling
leaves Q_4 in the state $|\downarrow\rangle$

C: modulate the exchange
interaction, map Q_3 to Q_4

D: read out " Q_3 "
leaves Q_3 and Q_4 in state $|\downarrow\rangle$



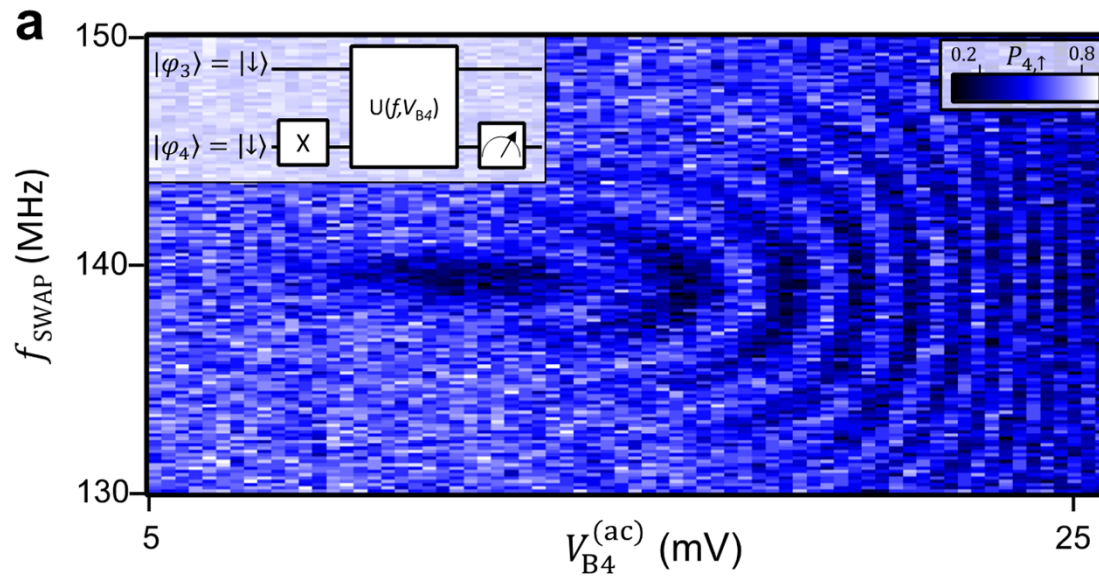
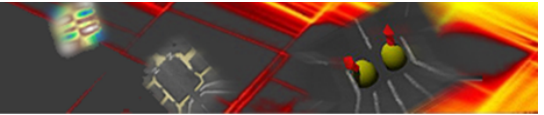
i: state preparation

ii: measurement

iii: resonant SWAP

iv: read out "Q₃"

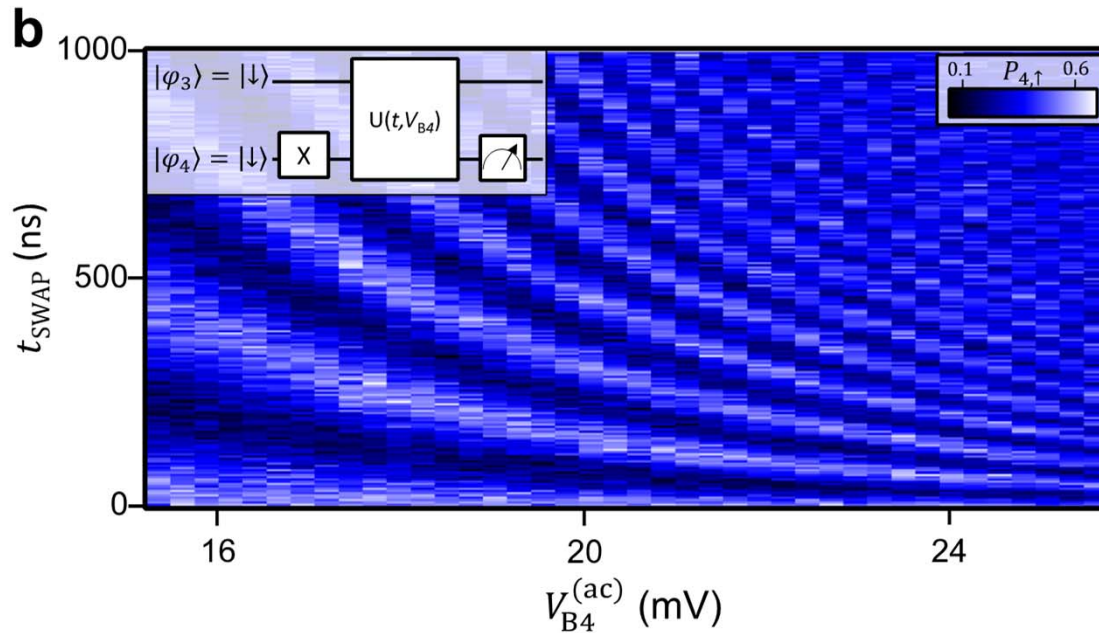
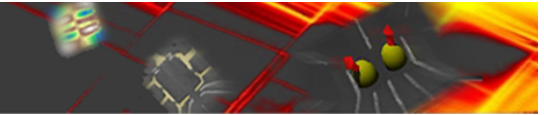
SWAP gate calibration



- initialize $|\phi_3, \phi_4\rangle = |\downarrow\downarrow\rangle$
- flip Q_4 using X gate
- f_{SWAP} burst on B_4 for 600 ns

- no oscillations at small $V_{B_4}^{(\text{ac})}$
- SWAP at $V_{B_4}^{(\text{ac})} = 10$ mV
- pattern is symmetric about $f_{\text{SWAP}} = 140$ MHz

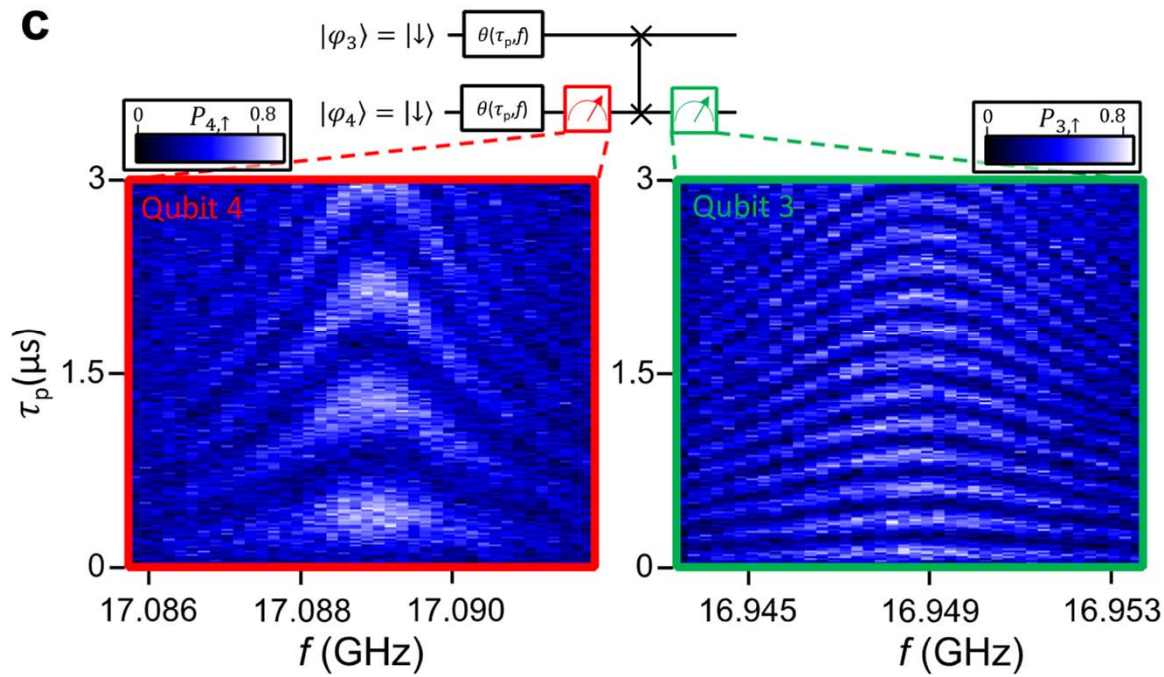
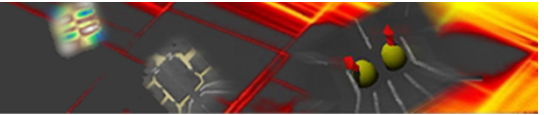
Minimize SWAP time



- fix $f_{\text{SWAP}} = 140$ MHz
- change burst time and ac amplitude

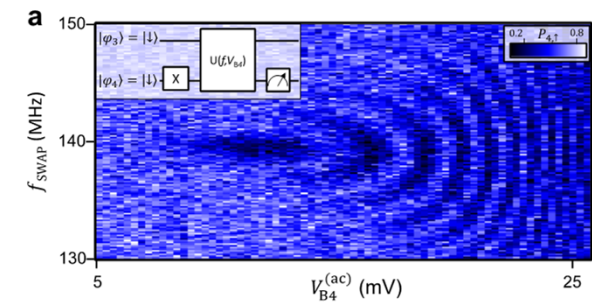
- bright fringe – even number of SWAPs
- Minimum $t_{\text{SWAP}} = 23$ ns (limited by control electronics)

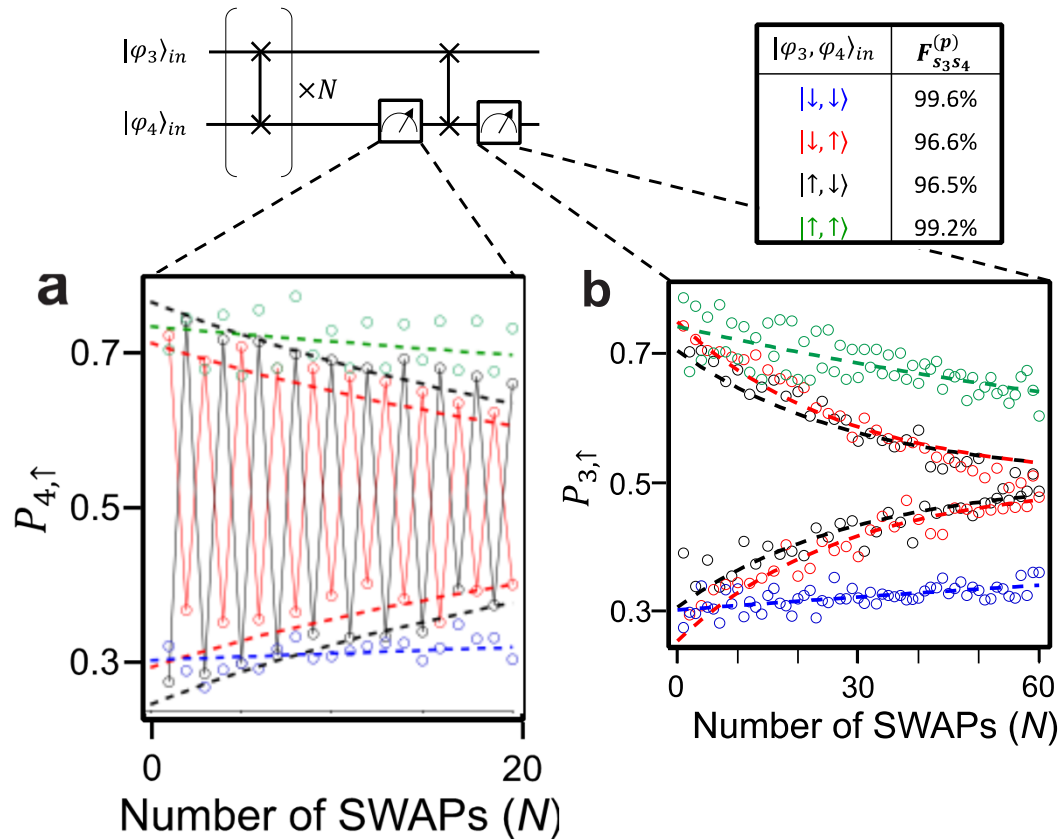
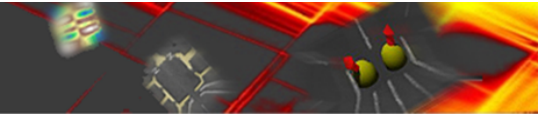
$$T_2^* \approx 10 \mu\text{s for both dots } (T_1 = 134 / 52 \text{ ms})$$



- Rabi oscillations (spacing largest on resonance)
- qubit difference frequency 140 MHz
- Initialize, control, and readout DQD with Q_4

- initialize $|\phi_3, \phi_4\rangle = |\downarrow\downarrow\rangle$
- RF burst
- measure Q_4
- projection-SWAP
- measure “ Q_3 ” (measure Q_4 infer Q_3)





- experiment insensitive to state preparation and measurement (SPAM) errors

- initial states:

$$|\phi_3, \phi_4\rangle_{in} = |\downarrow\downarrow\rangle, |\downarrow\uparrow\rangle, |\uparrow\downarrow\rangle, |\uparrow\uparrow\rangle$$

- execute SWAP gate N times

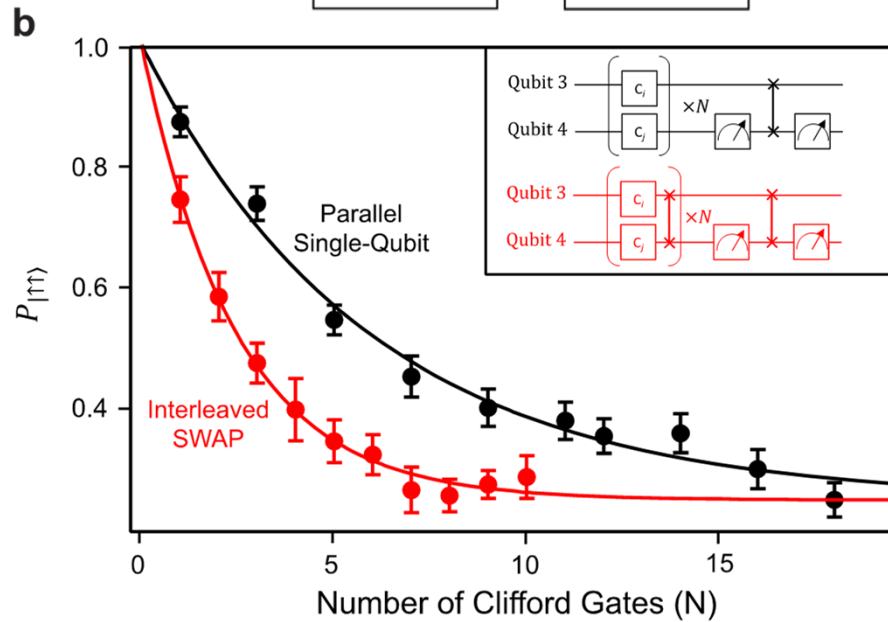
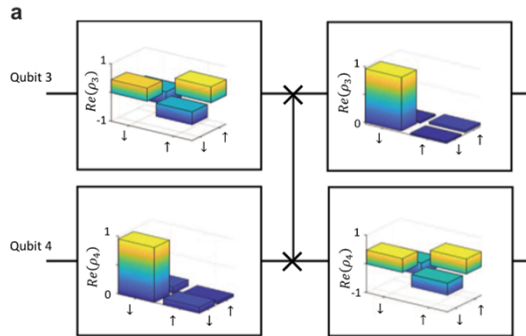
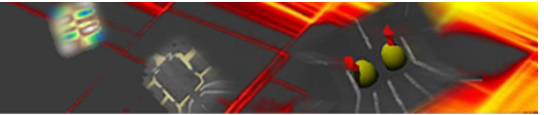
- $|\downarrow\uparrow\rangle$ and $|\uparrow\downarrow\rangle$ flip-flop for each SWAP
- decay envelope is given by fidelity

$$F_{\downarrow\uparrow}^{(p)} = F_{\uparrow\downarrow}^{(p)} = 96.5\%,$$

$$F_{\downarrow\downarrow}^{(p)} = 99.6\%, F_{\uparrow\uparrow}^{(p)} = 99.2\%$$

- $\bar{F}_{SWAP}^{(p)} = 98\%$

coherent-SWAP fidelity



- state tomography before and after SWAP of the superposition state Q_3 and spin down state Q_4 .

- additional calibration
- $(\alpha_1|\uparrow\rangle + \beta_1|\downarrow\rangle) \otimes (\alpha_2|\uparrow\rangle + \beta_2|\downarrow\rangle) \rightarrow (\alpha_2|\uparrow\rangle + \beta_2|\downarrow\rangle) \otimes (\alpha_1|\uparrow\rangle + \beta_1|\downarrow\rangle)$

- $\bar{F}_{SWAP}^{(c)} = 84\%$

Gate calibration

$$U = Z_3(\theta_3)Z_4(\theta_4)\text{SWAP},$$

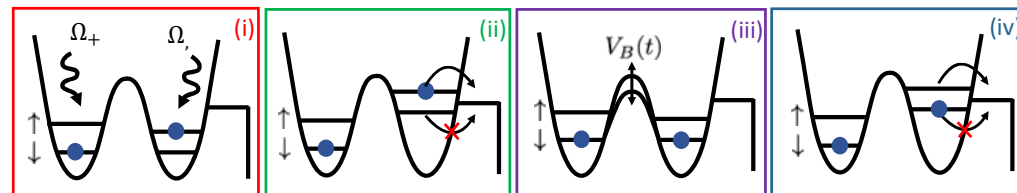
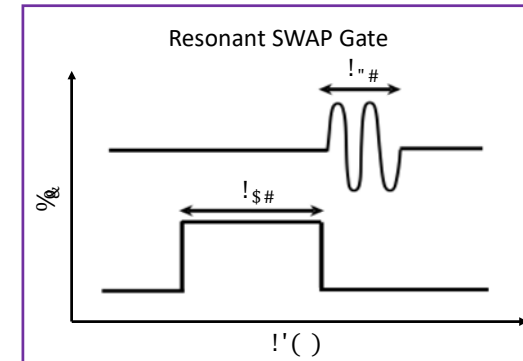
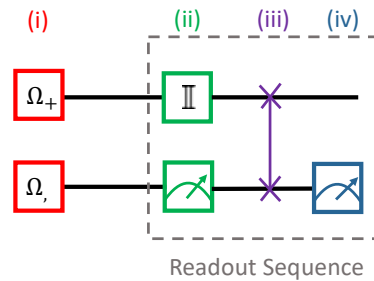
$$\theta_3 = \frac{(\Delta_3^{\text{dc}} + \Delta_4^{\text{dc}})t_{\text{dc}}}{2} - \frac{\Delta_{34}^{\text{dc}}}{2}t_{\text{dc}} + \frac{(\Delta_3^{\text{ac}} + \Delta_4^{\text{ac}})t_{\text{ac}}}{2} + \phi,$$

$$\theta_4 = \frac{(\Delta_3^{\text{dc}} + \Delta_4^{\text{dc}})t_{\text{dc}}}{2} + \frac{\Delta_{34}^{\text{dc}}}{2}t_{\text{dc}} + \frac{(\Delta_3^{\text{ac}} + \Delta_4^{\text{ac}})t_{\text{ac}}}{2} - \phi.$$

$\Delta_i = B_i - 2\pi f_i$ - magnetic field detuning

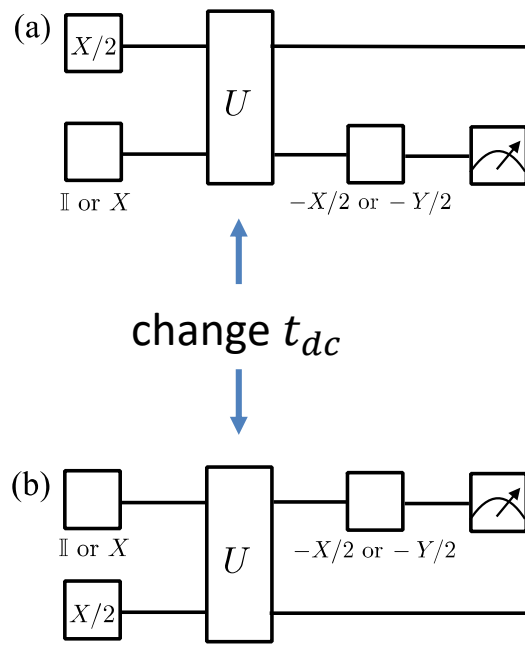
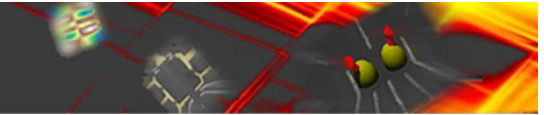
$\Delta_{34} = \Delta_3 - \Delta_4$ - shifted magnetic field gradient

ϕ is the phase of this ac



Calibration of the gate requires a precise measurement of $J_{34}^{\text{dc/ac}}$ and $\Theta_{3/4}$

Gate calibration



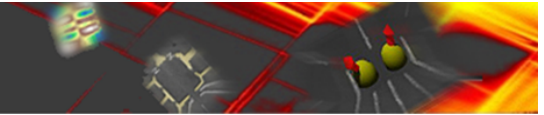
$$p_{\uparrow} = \frac{1}{2} + \frac{1}{2} \sin \left[\left(\Delta_3^{dc} - \frac{J_{34}^{dc}}{2} \right) t_{dc} + \frac{\Delta_3^{ac} + \Delta_4^{ac}}{2} t_{ac} - \frac{J_{34}^{ac}}{2} t_{ac} - \phi \right],$$

$$p_{\uparrow} = \frac{1}{2} - \frac{1}{2} \sin \left[\left(\Delta_3^{dc} + \frac{J_{34}^{dc}}{2} \right) t_{dc} + \frac{\Delta_3^{ac} + \Delta_4^{ac}}{2} t_{ac} + \frac{J_{34}^{ac}}{2} t_{ac} - \phi \right],$$

$$p_{\uparrow} = \frac{1}{2} + \frac{1}{2} \sin \left[\left(\Delta_4^{dc} - \frac{J_{34}^{dc}}{2} \right) t_{dc} + \frac{\Delta_3^{ac} + \Delta_4^{ac}}{2} t_{ac} - \frac{J_{34}^{ac}}{2} t_{ac} + \phi \right],$$

$$p_{\uparrow} = \frac{1}{2} - \frac{1}{2} \sin \left[\left(\Delta_4^{dc} + \frac{J_{34}^{dc}}{2} \right) t_{dc} + \frac{\Delta_3^{ac} + \Delta_4^{ac}}{2} t_{ac} + \frac{J_{34}^{ac}}{2} t_{ac} + \phi \right],$$

$\Delta_{3/4}^{dc}, J_{34}^{dc/ac},$
 $(\Delta_3^{ac} + \Delta_4^{ac})/2, \phi$

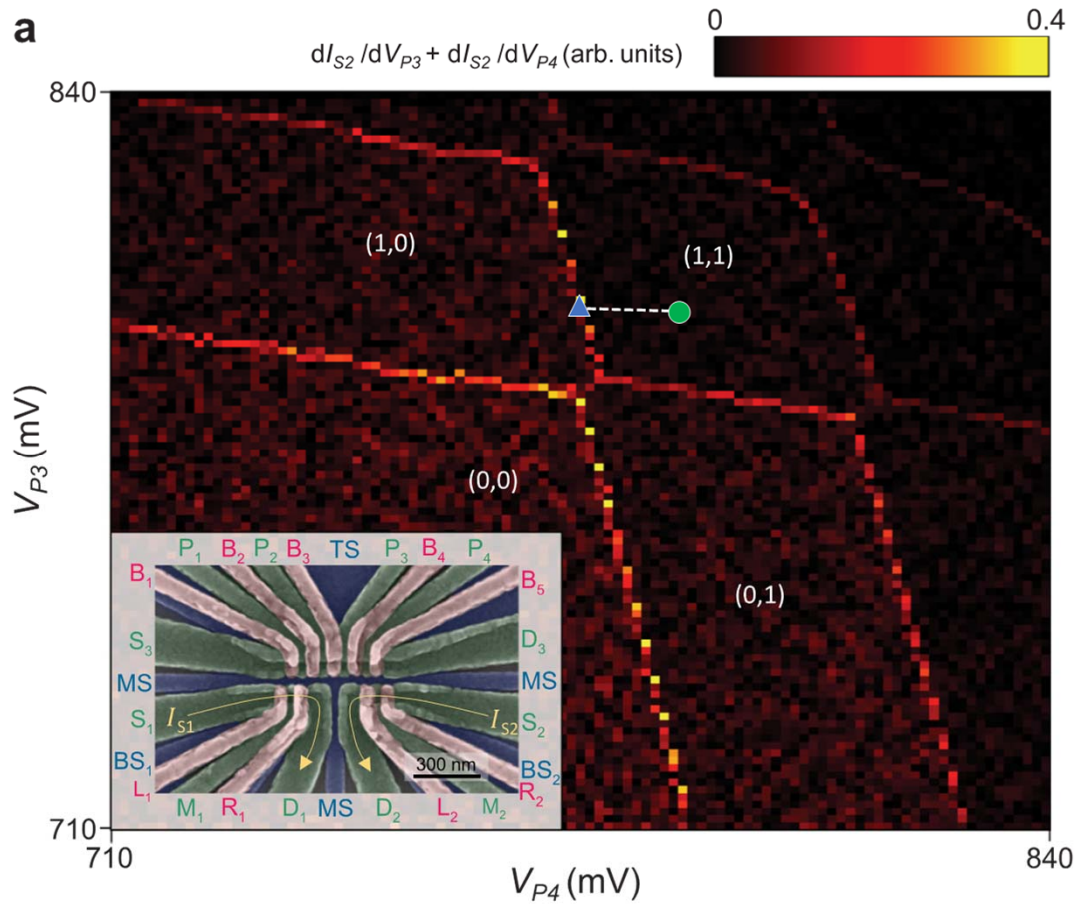
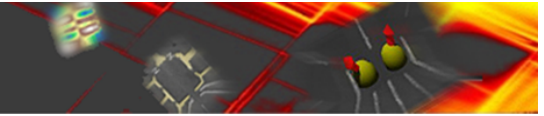


Conclusion and outlook

- coherent spin transport in an array of quantum dots
- transfer arbitrary two-qubit states between spins
- no moving charges

- shuttle a spin projection across nine-dot array with 85% fidelity
- compatible with singlet-triplet and cavity dispersive readouts

Charge stability diagram



Experimental setup

